

Improved nuclear physics for supernovae

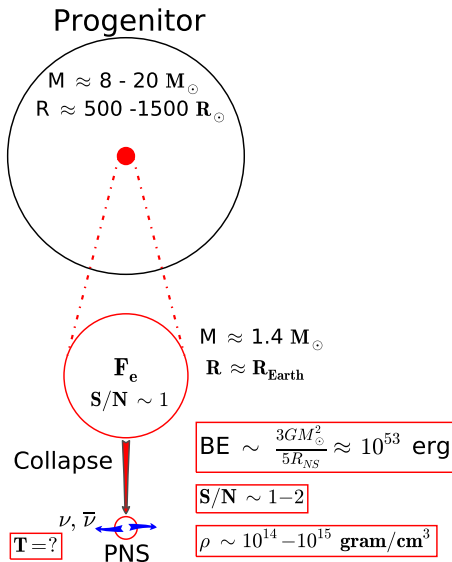
Implications for neutrino spectra, nucleosynthesis and dark matter

Ermal Rrapaj

University of Washington, Seattle

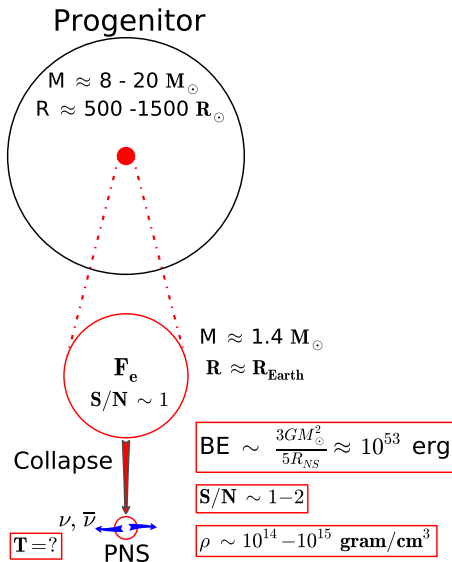
December 16, 2015

Core Collapse Supernovae

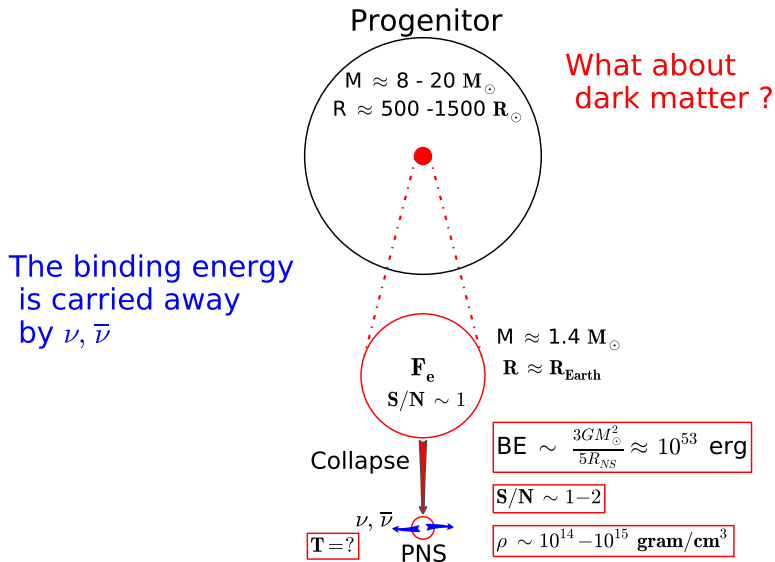


Core Collapse Supernovae

The binding energy
is carried away
by $\nu, \bar{\nu}$

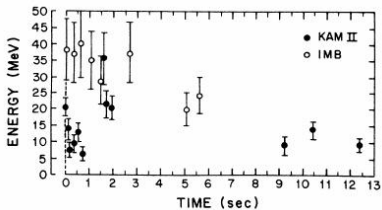


Core Collapse Supernovae



Detection of SN87A

Hirata et al, Phys. Rev. Lett. 58, (1490)



12 events in the burst sample observed in Kamiokande-II, and 8 events in the burst sample observed in the IMB detector

Dark Matter, Why ?

Astron.Astrophys. 571 (2014) A16

Energy composition of our universe: Planck Data

- ▶ Dark Energy: 68.3%
- ▶ Dark Matter (DM): 26.8%
- ▶ Atomic Matter: 4.8%
- ▶ Light: 0.005%
- ▶ Neutrinos: 0.0034%

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So, what is Dark Matter?

Can we 'see' dark matter?

Essig et al, arxiv:1311.0029 (2013)

Portal	Particles	Operators
"Vector"	Dark Photons	$-\epsilon e J_{\mu}^{\text{SM}}$
"Axion"	PseudoScalars	$\frac{a}{f_a} G_{\mu\nu} \tilde{G}^{\mu\nu}, \frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}, \frac{\partial_{\mu} a}{f_a} \bar{\Psi} \gamma^{\mu} \gamma^5 \Psi$
"Higgs"	Dark Scalars	$(\mu S + \lambda S^2) H^{\dagger} H$
"Neutrino"	Sterile Neutrinos	$y_N L H N$

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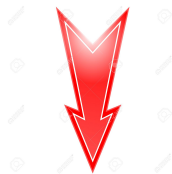
This talk will be focused on Dark Photons!

Dark Photons: How?

Holdom, Phys. Rev. B 166, (1986) 196

High - Energy

$$\mathcal{L} \supset -\frac{1}{4}(B_{\mu\nu})^2 - \frac{1}{4}(F'^{\mu\nu})^2 - \frac{\epsilon_Y}{2 \cos \theta_W} B_{\mu\nu} F'^{\mu\nu} + \frac{1}{2} m_A^2 A'^2 + g_D J_\mu^D A'^\mu$$



Low - Energy

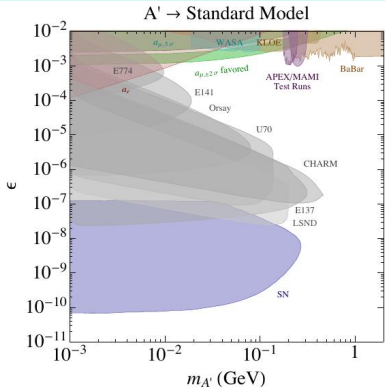
$$\mathcal{L} \supset g_Q A'_\mu J_\mu^{\text{EM}} - \frac{1}{2} m_{\gamma Q}^2 A'_\mu A'^\mu + [g_B V_\mu^B J_\mu^B - \frac{1}{2} m_{\gamma B}^2 (V_\mu^B)^2]$$

Lee, Yang Phys. Rev. 98, 1501 (1955)

Batell, deNerville, McKeen, Pospelov, Ritz Phys. Rev. D 90, 115014 (2014)

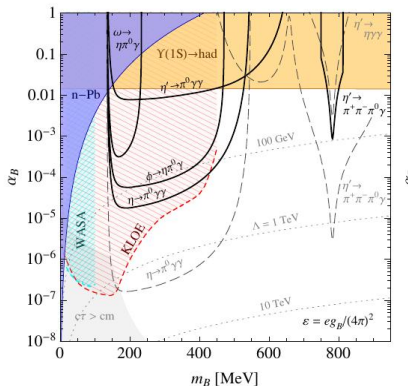
Dark Photons: Why?

γ_Q Parameter Space



Snowmass report (2013) arXiv:1311.0029

γ_B Parameter Space

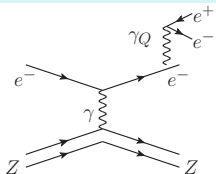


Tulin Phys. Rev. D 89, 11408 (2014)

Dark Photons: Why?

<http://hallaweb.jlab.org/experiment/APEX>

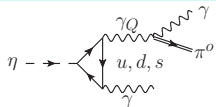
Experimental searches for γ_Q



APEX (spring 2016) at Jefferson Laboratory

- ▶ $10^{-2} \lesssim \epsilon_Q \lesssim 10^{-10}$
- ▶ $65 \text{ MeV} \leq m_{\gamma_Q} \leq 550 \text{ MeV}$

Experimental searches for γ_B



Jlab Eta Factory (JEF) Experiment

- ▶ $10^{-1} \lesssim \alpha_B \lesssim 10^{-7}$
- ▶ $140 \text{ MeV} \leq m_{\gamma_B} \leq 550 \text{ MeV}$

<https://cnidlamp.jlab.org/RareEtaDecay/JDocDB/system/files/biblio/2015/04/jef-gan-aps-pdf.pdf>

And many more other experiments

Implications for supernovae

Dark matter cooling of PNS

1. produced in the hot core
2. light in mass \rightarrow copious amounts!
3. sap energy from the core
4. reduce neutrino energy and burst duration!

Raffelt's Criteria

Conditions for fiducial calculations:

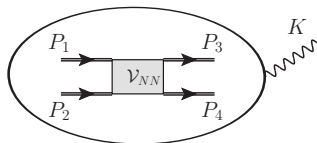
- ▶ $T = 30 \text{ MeV}$
- ▶ $\rho = 3 \times 10^{14} \text{ gram/cm}^3$

$$\dot{E} \approx 10^{19} \frac{\text{erg}}{\text{gram s}} \implies \text{Neutrino burst duration is halved!}$$

"Stars as laboratories for fundamental physics: The astrophysics of neutrinos, axions, and other weakly interacting particles" (University of Chicago Press, 1996)

How is light dark matter produced ?

Nucleon Bremsstrahlung



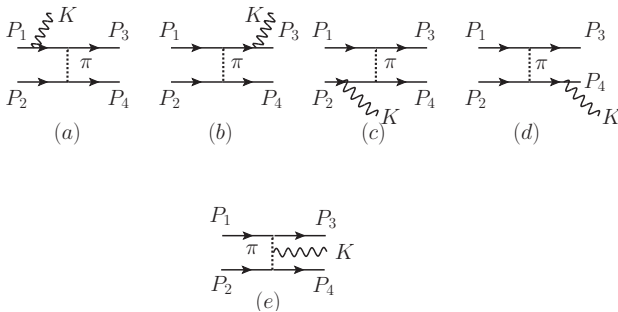
What has been studied so far ?!

- ▶ Axions, Burrows, Turner, Brinkman, *Phys. Rev. D* 39, 1020,(1989)
- ▶ Kaluza–Klein gravitons and dilatons, Hannart, Phillips,Reddy, Savage, *Nuc. Phys. B* 595(2001)
- ▶ Neutralinos Dreiner, Hanhart, Langenfeld, Phillips *Phys. Rev. D* 68, 055004 (2003)
- ▶ Dark Photons Dent, Ferrer, Kraus *arxiv: 1201.2683* (2012)
Kazanas, Mohapatra, Nussinov Teplitz, Zhang *Nuc. Phys. B.* 90, 17, (2014)

Current Dark Photon Constraint: Supernova Cooling

Dent, Ferrer, Kraus arxiv: 1201.2683 (2012)

Production by Bremsstrahlung

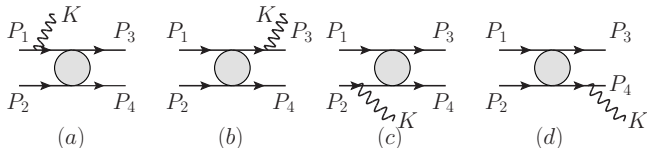


Emissivity

$$\mathcal{L}_{\tilde{\gamma}} \leq 10^{53} \text{ erg/s} \longleftrightarrow \dot{E} \leq 8 \times 10^{22} \text{ erg/g/s}$$

Bremsstrahlung: Soft Radiation Approximation (SRA)

Low Phys. Rev. 110, 974 (1958)



$$d\sigma_{pp \rightarrow pp\tilde{\gamma}} \approx -4\pi\alpha_{\text{em}}\epsilon^2 \frac{d^3k}{2\omega} (\epsilon^\mu J_\mu)^2 d\sigma_{NN \rightarrow NN}$$

$$j_\mu^{(2)} = \left(\frac{P_1}{P_1 \cdot K} - \frac{P_3}{P_3 \cdot K} \right)_\mu, \quad j_\mu^{(4)} = \left(\frac{P_1}{P_1 \cdot K} + \frac{P_2}{P_2 \cdot K} - \frac{P_3}{P_3 \cdot K} - \frac{P_4}{P_4 \cdot K} \right)_\mu$$

Rrapaj, Reddy arxiv:1511.09136

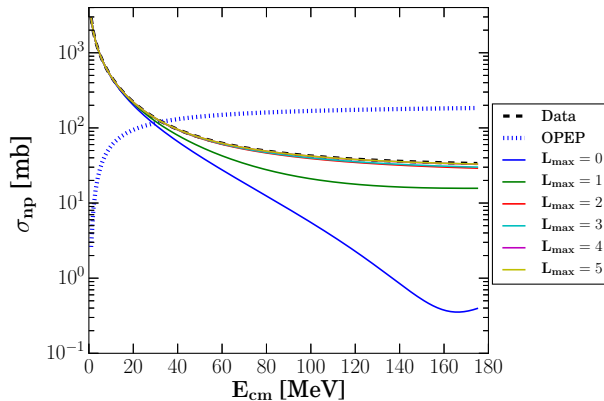
$$\dot{\epsilon}_{np \rightarrow np\gamma Q} = \frac{2\alpha_{\text{em}}\epsilon_Q^2}{\sqrt{\pi}} \frac{n_n n_p}{(MT)^{3/2}} T^4 \int_{m\gamma_Q/T}^{\infty} dx e^{-x} x^3 \mathcal{I}^{(2)}\left(\frac{m\gamma_Q/T}{x}\right) \sigma_{np}^{(2)}(xT)$$

$$\dot{\epsilon}_{ij \rightarrow ij\gamma B} = \frac{2\alpha_{\text{em}}\epsilon_B^2}{\sqrt{\pi}} \frac{n_i n_j}{(MT)^{3/2}} T^5 \int_{m\gamma_B/T}^{\infty} dx e^{-x} \frac{x^4}{M} \mathcal{I}^{(4)}\left(\frac{m\gamma_B/T}{x}\right) \sigma_{ij}^{(4)}(xT)$$

$$\sigma_{ij}^{(2)} = \int d\cos\theta_{\text{cm}} \frac{d\sigma_{n_i n_j \rightarrow n_i n_j}}{d\theta_{\text{cm}}} (1 - \cos\theta_{\text{cm}}), \quad \sigma_{ij}^{(4)} = \int d\cos\theta_{\text{cm}} \frac{d\sigma_{n_i n_j \rightarrow n_i n_j}}{d\theta_{\text{cm}}} (1 - \cos^2\theta_{\text{cm}})$$

Partial Wave Expansion vs OPEP

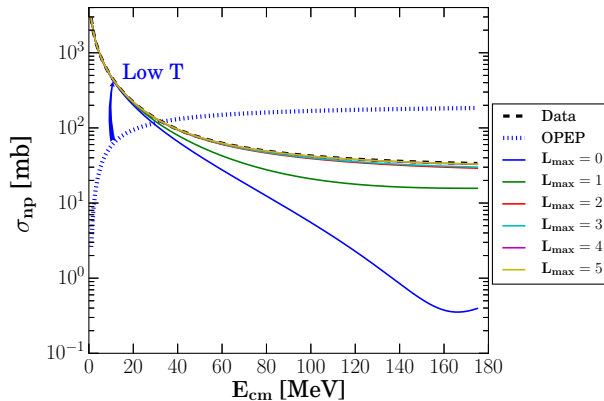
Rrapaj, Reddy arxiv:1511.09136



Data from Nijmegen University database

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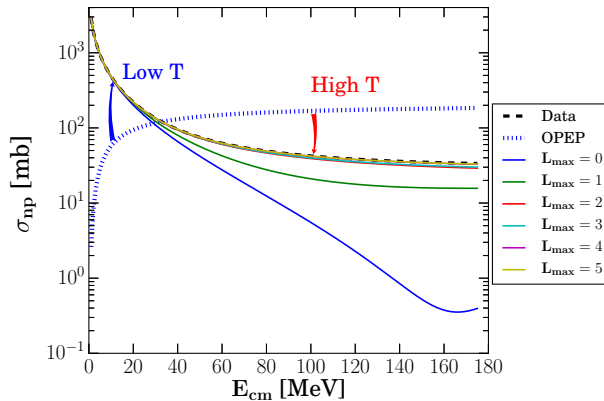
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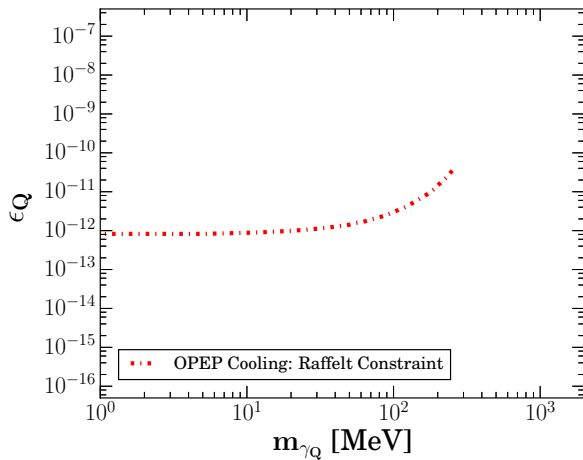
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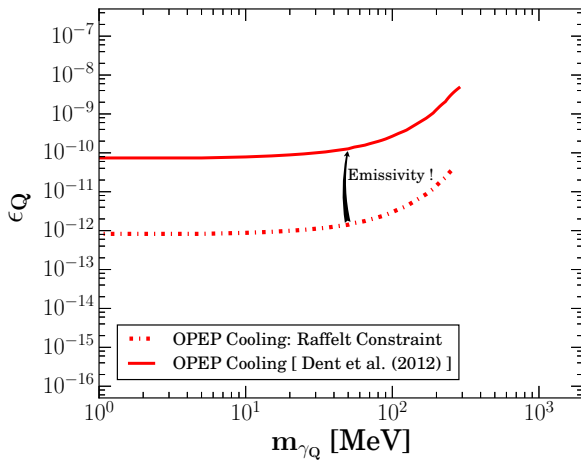
Dark Photon Constraint

Rrapaj, Reddy [arxiv:1511.09136](https://arxiv.org/abs/1511.09136)



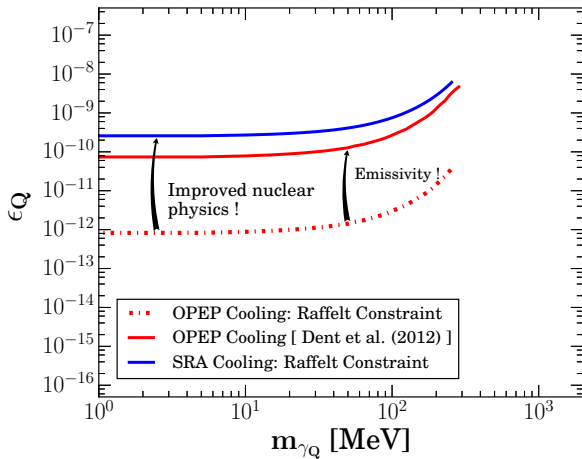
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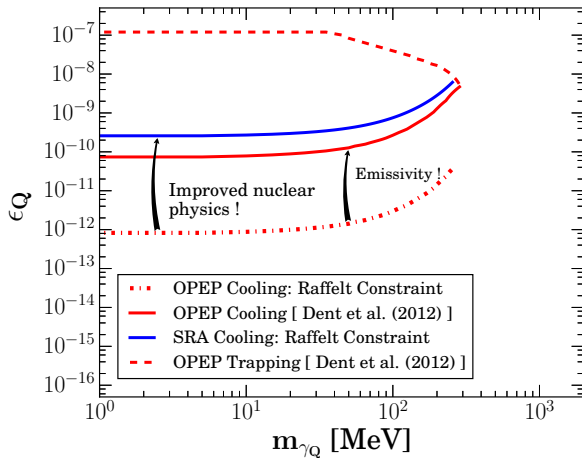
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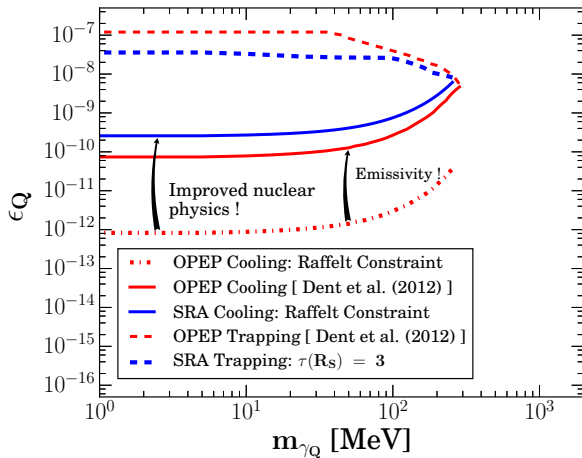
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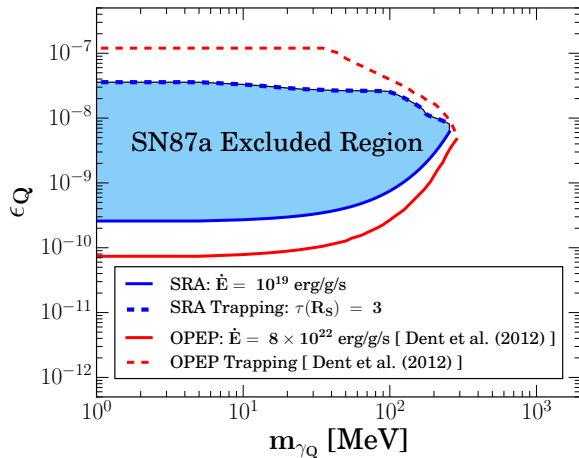
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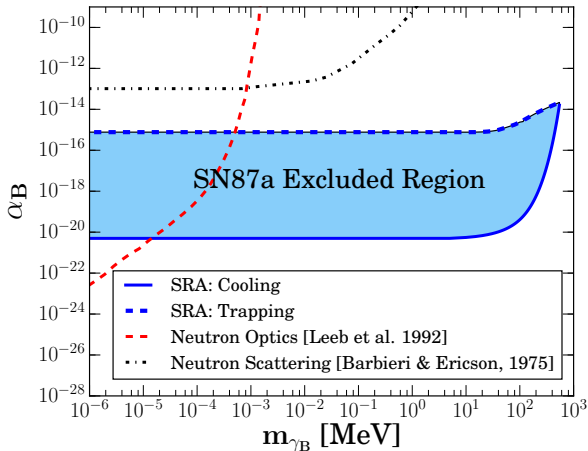


What about the dark leptophobic photon ?!

No Current Supernova constraint!

Dark photon coupled only to baryonic current

Rrapaj, Reddy arxiv:1511.09136



So, did we 'solve' this issue?

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NOPE

Scattering Rate: Bremsstrahlung

- ▶ SRA valid only for $\omega/E_{CM} \ll 1$
(perhaps effective field theories and two body currents)
- ▶ medium effects not included
(rescattering, nucleon excitation lifetimes)
- ▶ ...

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- ▶ ...

Supernova environment

- ▶ $S/N \sim 1 - 2$
- ▶ $\rho_0 = 0.16 \text{ fm}^{-3} \leq \rho_C \leq ? \rightarrow$ Equation of State (EoS)
- ▶ $T = ? \rightarrow$ Specific heat \rightarrow EoS $\rightarrow ?$

Constraining mean field models

- ▶ Saturation density properties Dutra, Lourenco, Martins,Delfino, *Phys.Rev.C* 85, 035201 (2012)
- ▶ Low density neutron matter, ab-initio methods Brown,Schwenk *Phys.Rev.C* 91,049902 (2015)
Compare many-body perturbation and Monte Carlo using χ EFT
Rrapaj, Roggero, Holt *arxiv:1510.00444* (2015)

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Found mean field models consistent with all constraints!

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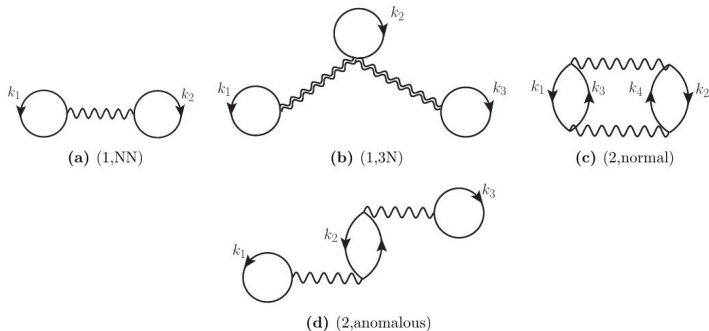
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Not currently used in supernovae simulations!

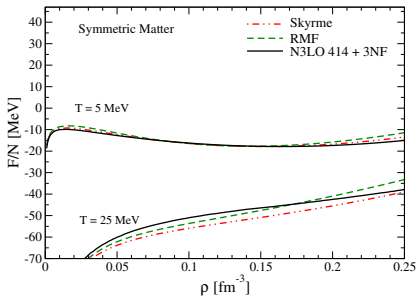
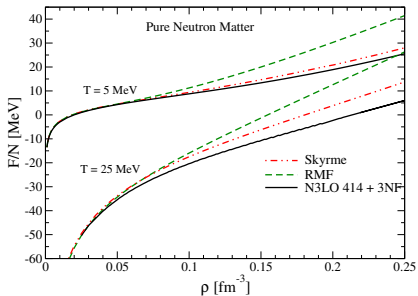
Many-body perturbation theory , $T > 0$

Wellenhofer, Holt, Kaiser, Weise Phys. Rev. C 89, 064009 (2014)



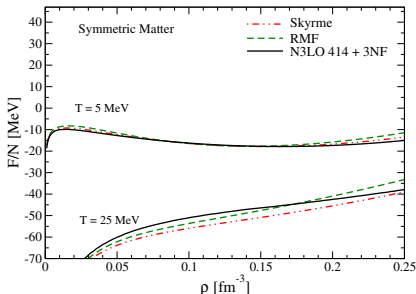
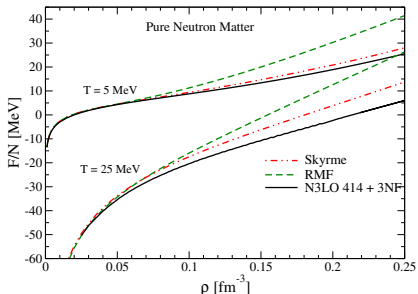
$T > 0$ EoS: Skyrme interactions and relativistic models

Rrapaj, Roggero, Holt [arxiv:1510.00444](https://arxiv.org/abs/1510.00444) (2015)



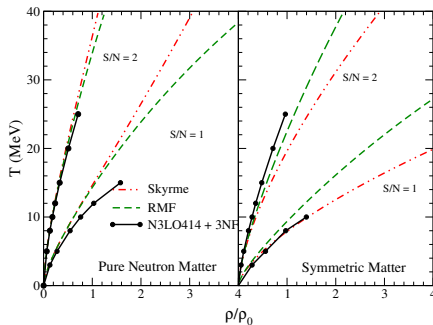
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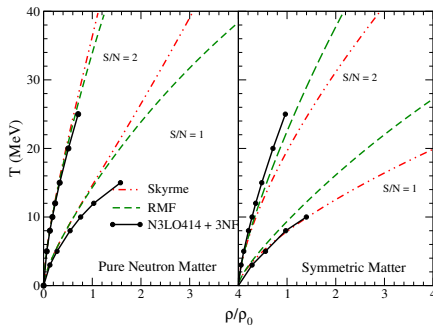
'Qualitatively' comparable with many-body calculations

S = const trajectories



Core temperature uncertain!

S = const trajectories



Beyond mean field signature !

Thank You

- ▶ Sanjay Reddy
- ▶ Jeremy Holt
- ▶ Alessandro Roggero
- ▶ Alexander Bartl
- ▶ Achim Schwenk