



# Langevin Model of Low-Energy Fission

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Dedicated to the memory of J. R. Nix, 1938-2008,  
who led the Los Alamos fission theory efforts for 30  
years and inspired the developments presented today.

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## Essential Contributors:

- Peter Moller for years of collaboration and for developing the potential-energy model.
- John Lestone for enlightening discussions and providing an accurate neutron-evaporation code.
- Jorgen Randrup for help with nuclear inertia calculations and discussions of dissipation.
- Morgan White for encouragement and support.

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# Introduction

- Fission was discovered 79 years ago.
- For the first time, a dynamical physical model of fission predicts most of the major features and some of the quantitative details of fission mass and kinetic-energy distributions.

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# Ingredients of a dynamical model

- Potential energy of a nucleus as a function of its shape (nuclear binding energy).
- Kinetic energy of nuclear shape motions.
- Dissipation of nuclear shape motions: damping of motion and fluctuating force (Brownian motion).
- Equation of motion.

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# Potential energy of a deformed nucleus

- Use the macroscopic-microscopic model.
- Parameters of the model are determined by nuclear ground-state properties.
- Fission properties are an extrapolation of the model outside where its parameters are determined.
- Excellence of predictions a tribute to its creators: Nilsson, Nix, and Moller (~1970-2015).

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# Kinetic energy of nuclear shape motions

- Use irrotational fluid flow to define nuclear inertia.
- Allow scaling of the inertia to crudely represent the behavior of real nuclear matter.

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# Dissipation in nuclear shape motions

- Dissipation implies both damping of motion and a random force in the dynamics.
- Dissipation is modeled as a one-body effect with an arbitrary strength parameter.
- The random force is proportional to the square root of (nuclear temperature) x (the dissipation tensor).

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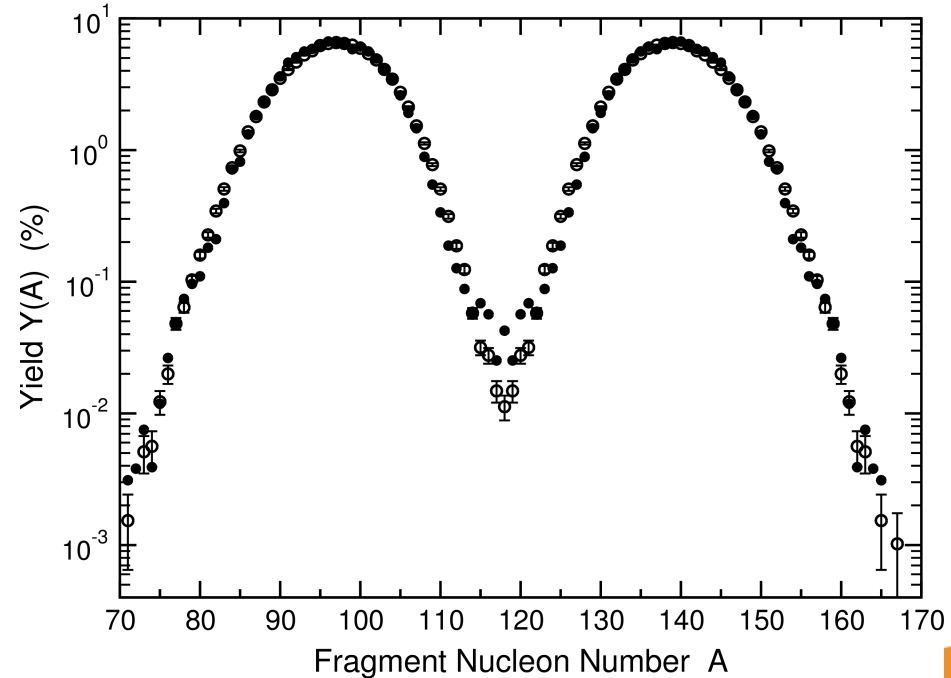
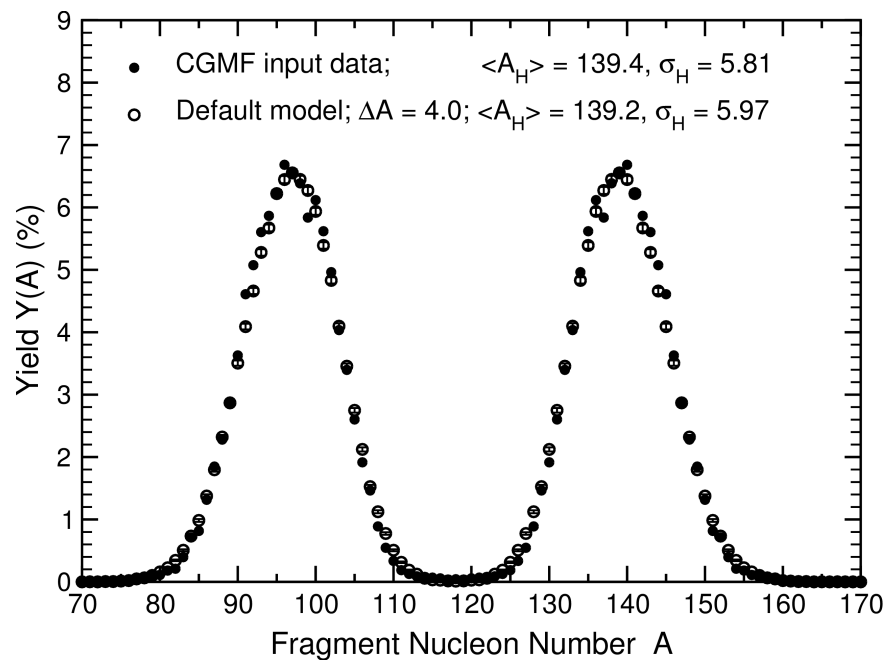


# Equations of motion: Langevin Equation

$$\frac{dq_j}{dt} = \frac{\partial H}{\partial p_j} = \frac{\partial(K + V)}{\partial p_j} = \frac{\partial(\frac{1}{2} M_{ik}^{-1} p_i p_k)}{\partial p_j} = M_{jk}^{-1} p_k \quad (1)$$

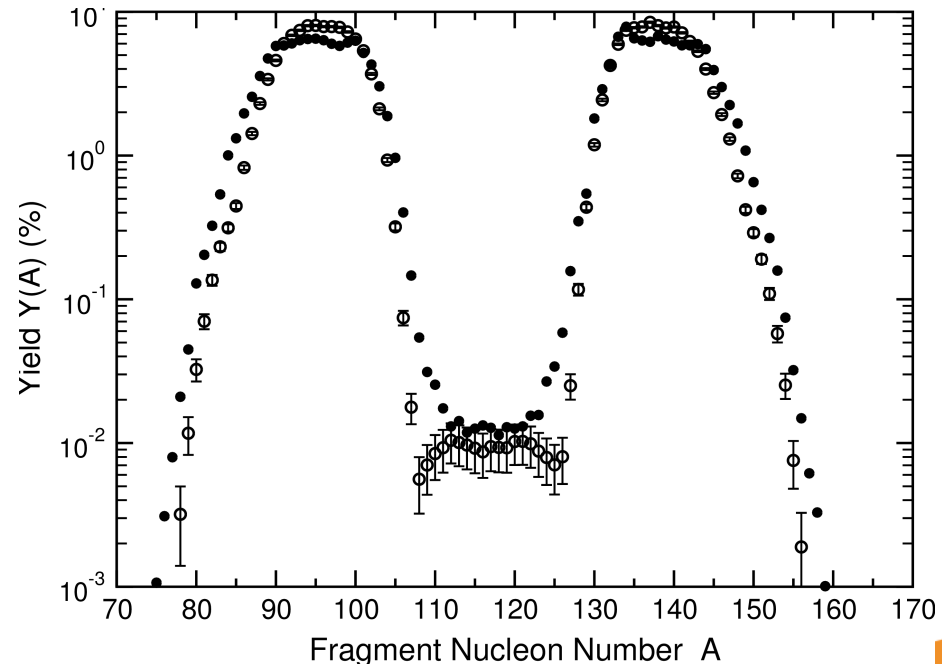
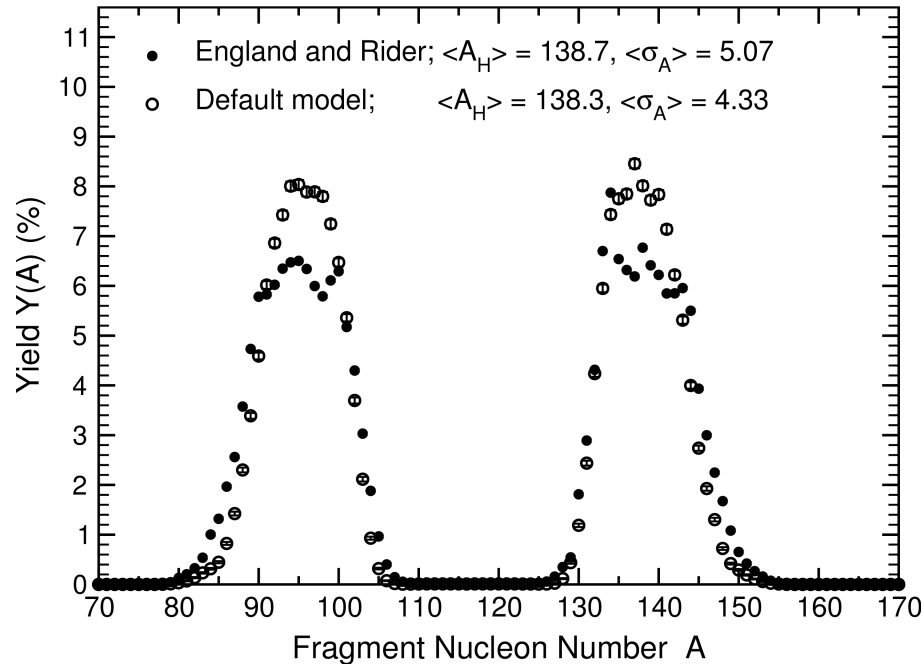
$$\frac{dp_j}{dt} = -\frac{\partial V}{\partial q_j} + \frac{1}{2} \frac{\partial M_{kl}}{\partial q_j} \dot{q}_k \dot{q}_l - \eta_{jk} \dot{q}_k + \sqrt{\frac{2T}{\Delta t}} \gamma_{jk} \Theta_k(t). \quad (2)$$

# Unadjusted baseline model; thermal fragment yields; U235 (n,f)



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# Unadjusted baseline model; thermal prompt yields U235 (n,f)



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# Default model predictions

- Prompt yields poorly predicted.
- Primary yields consistent with measurements.
- Model  $\langle \text{TKE} \rangle = 173.7 \text{ MeV}$ ; Expt. =  $170.9 \text{ MeV}$ .

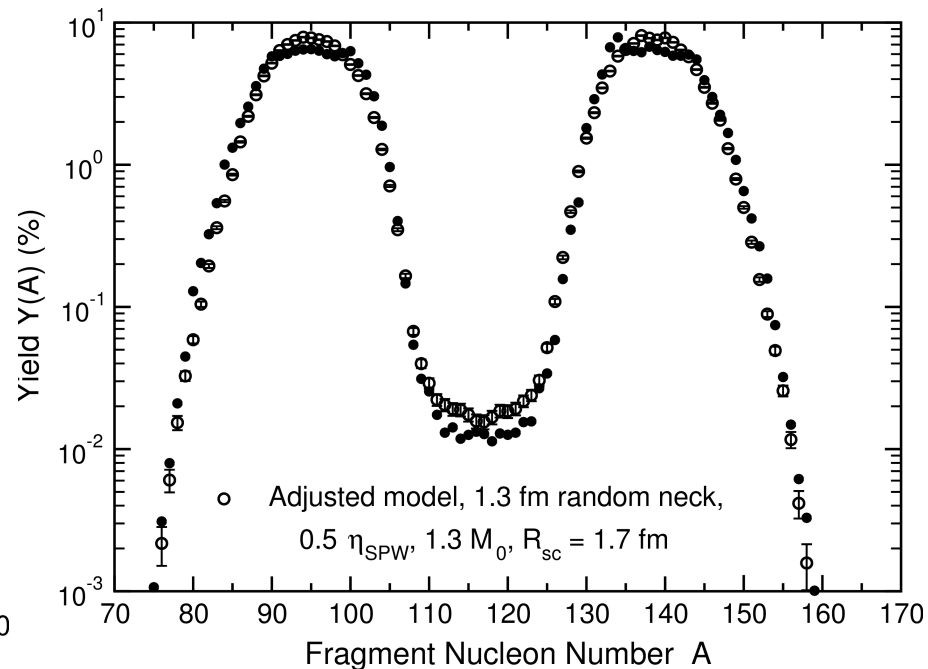
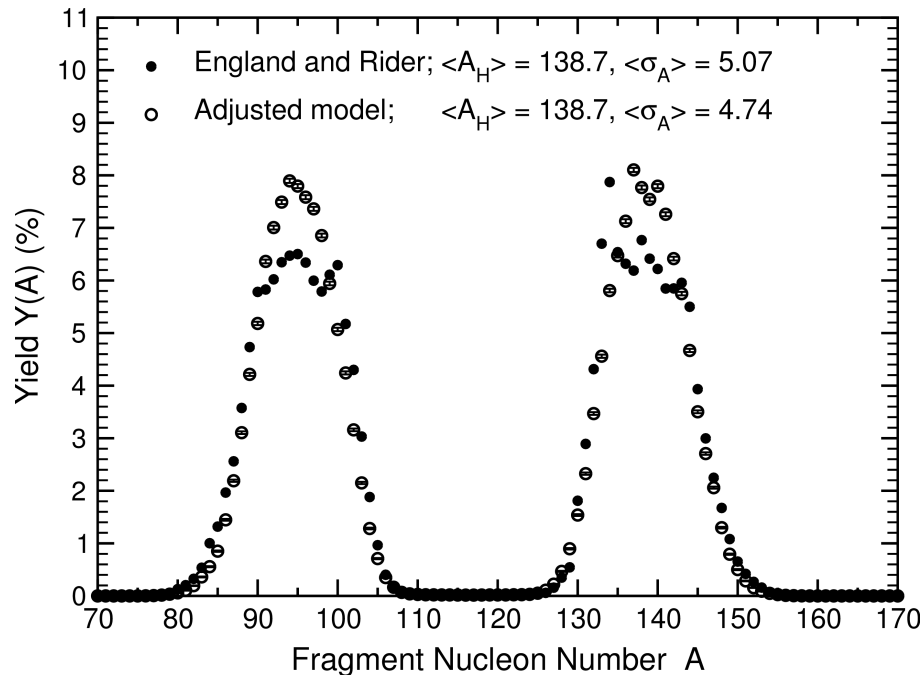
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## Physical parameters in 'adjusted' model

- Scaling of hydrodynamical inertia; value 1.3.
- Dissipation strength; scale = 0.5 of default.
- Neck radius at scission; 1.7 fm vs. 1.0 fm.
- Random neck rupture;  
gaussian with standard deviation of 1.3 fm.
- Probability of starting in symmetric valley; 0.0.

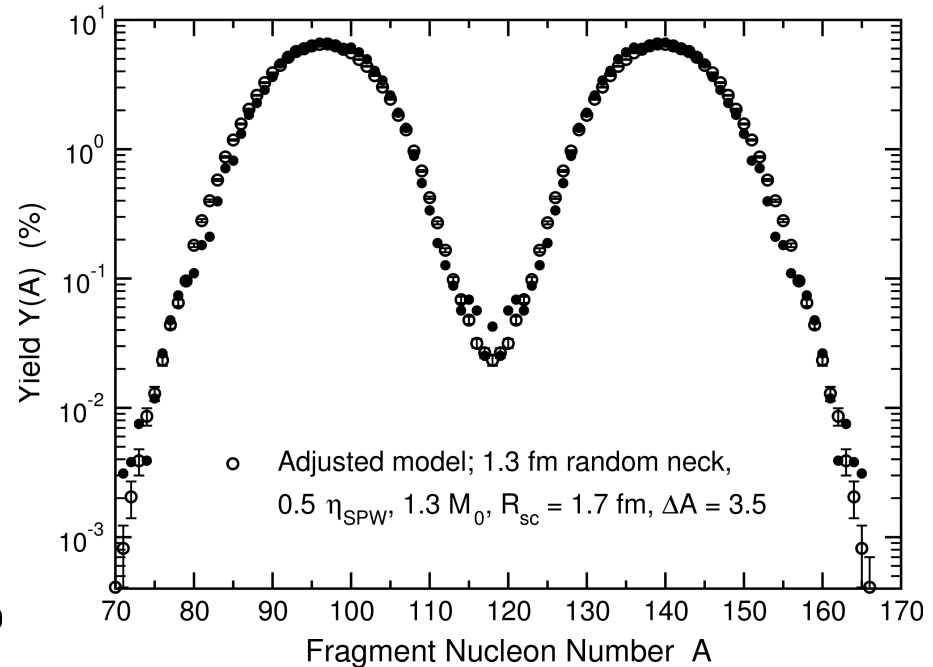
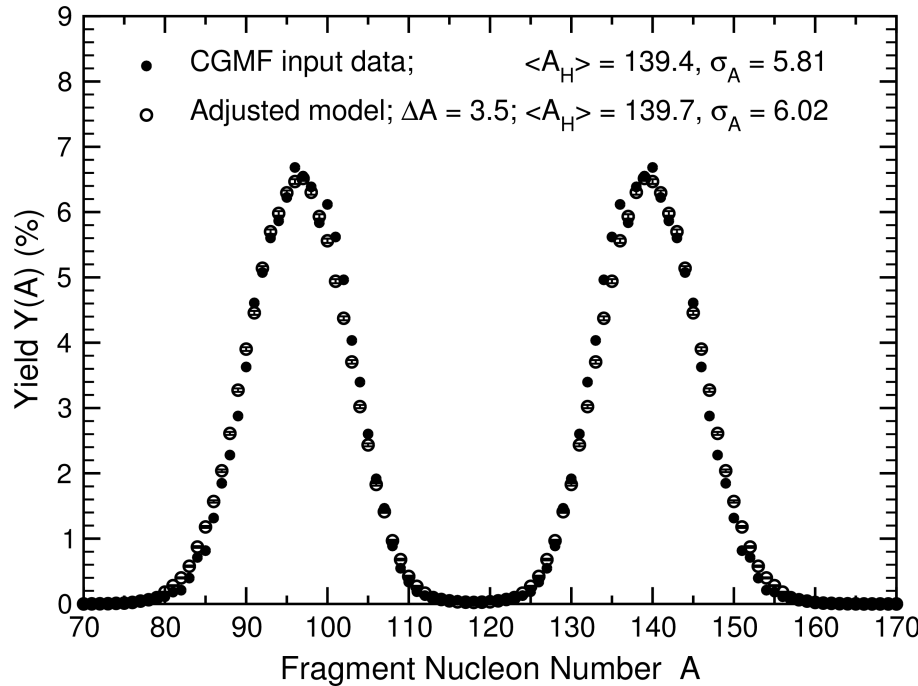
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# 'Adjusted' model; thermal prompt yields; $^{235}\text{U}(n,f)$



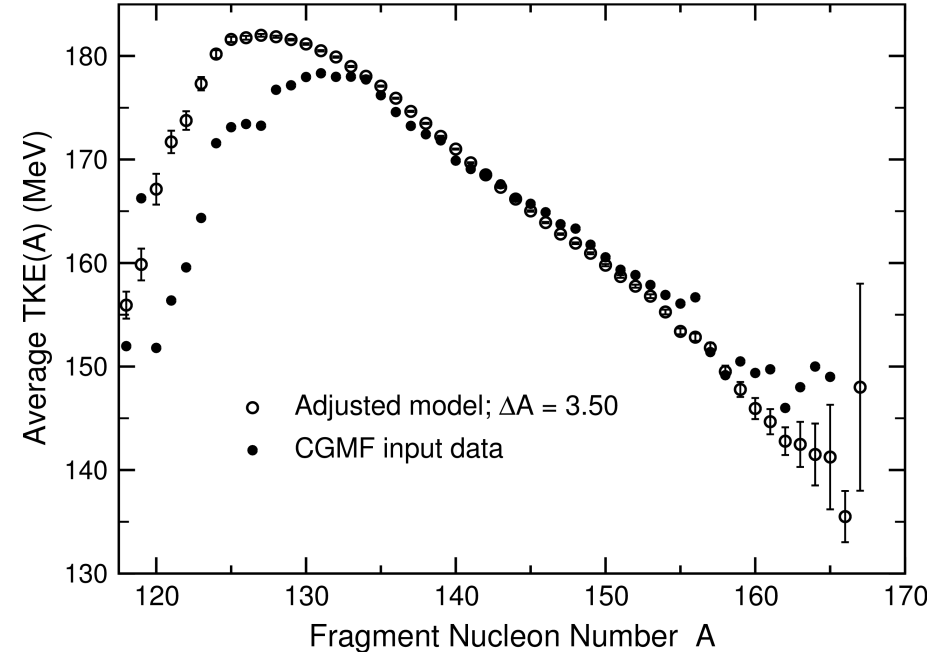
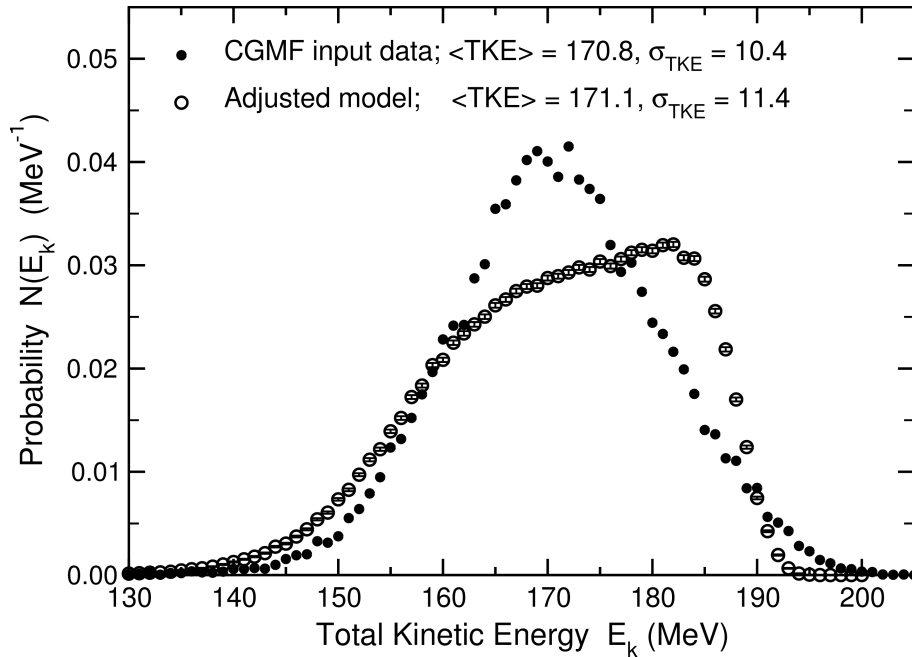
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# 'Adjusted' model; thermal fragment yields; U235 (n,f)



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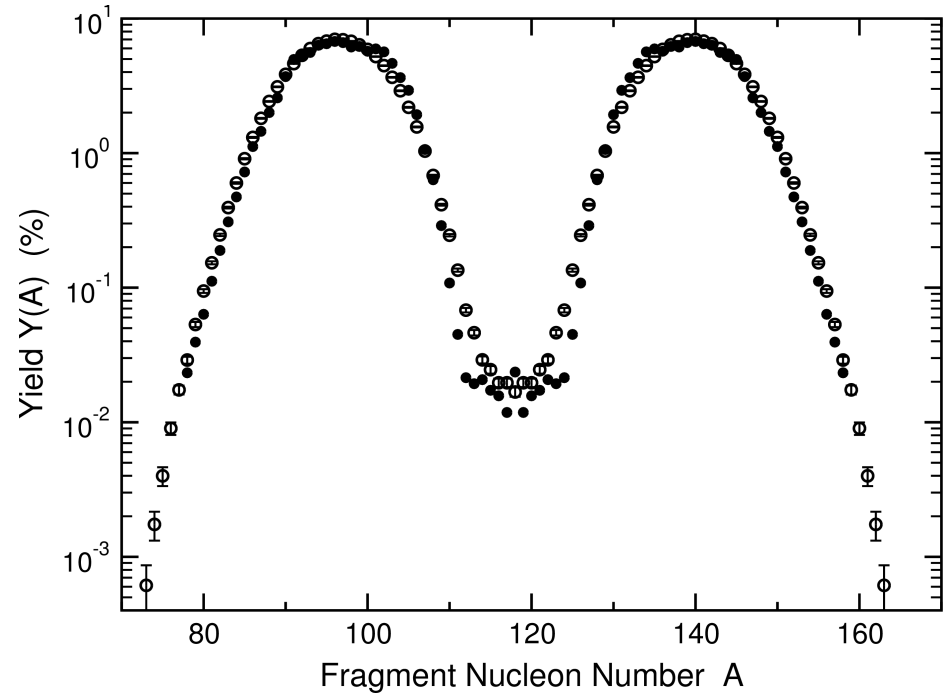
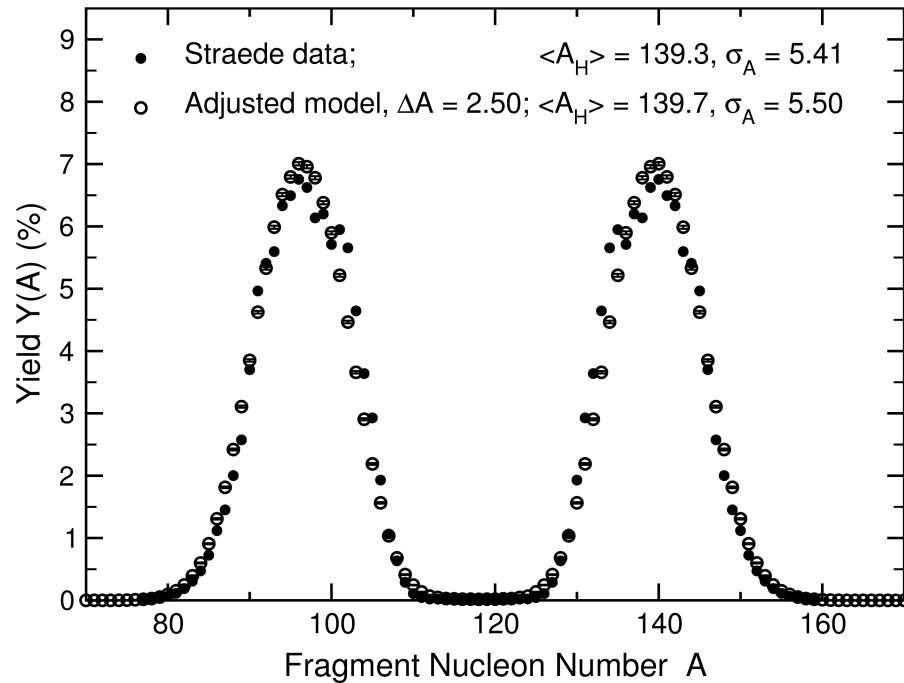
# 'Adjusted' model; TKE yields; U235 (n,f)



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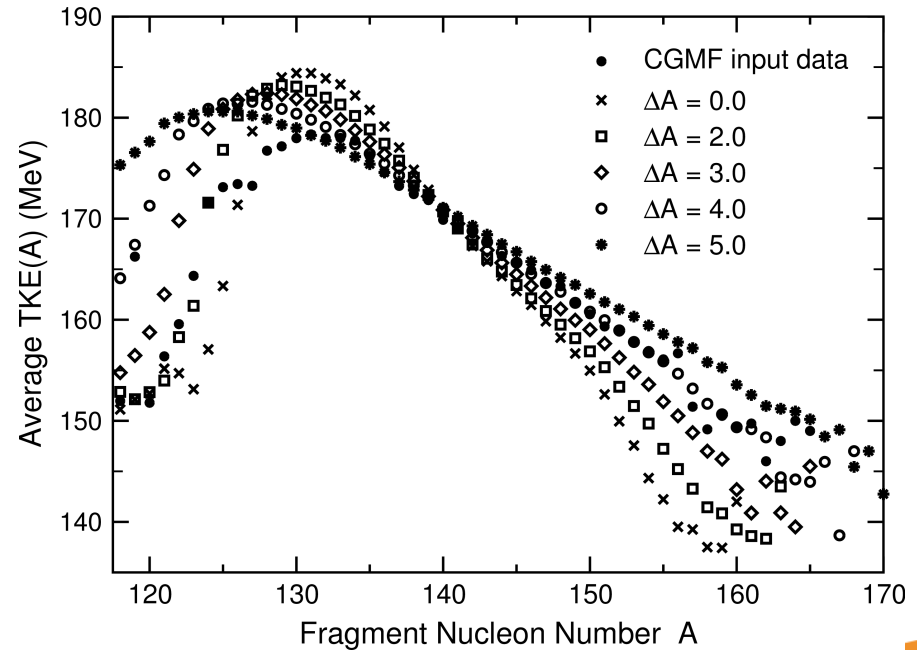
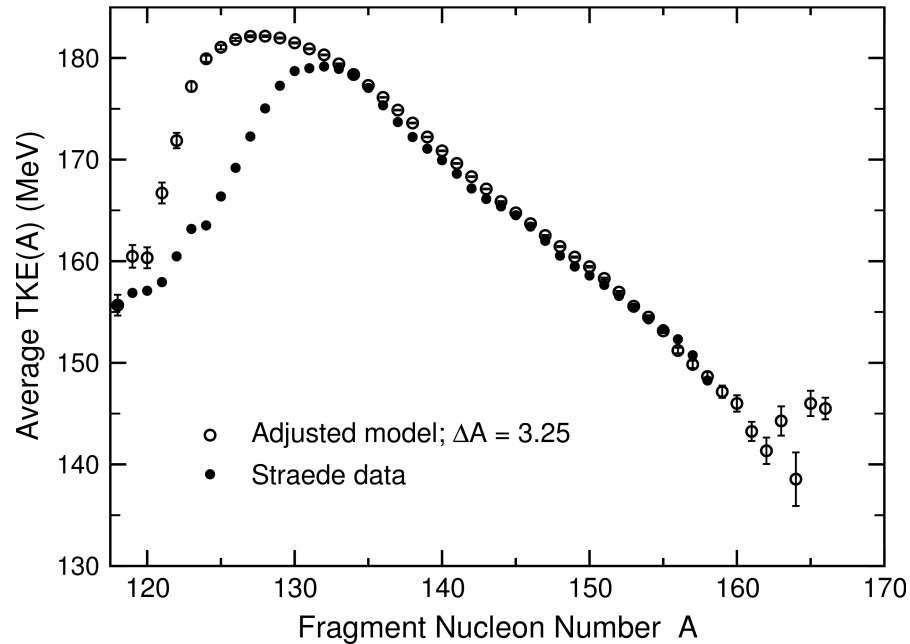


# 'Adjusted' model; thermal fragment yields; U235 (n,f)



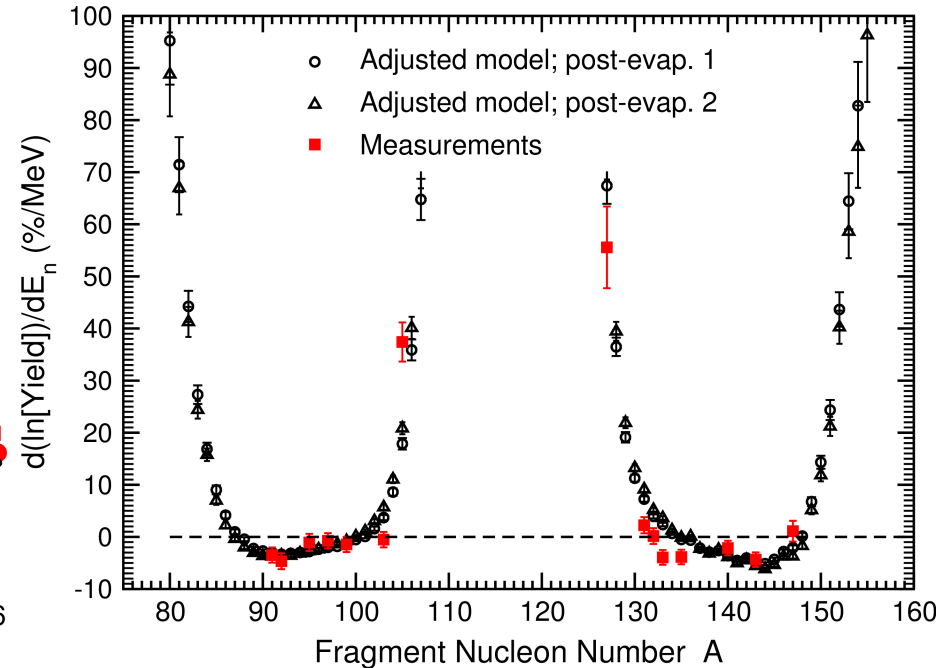
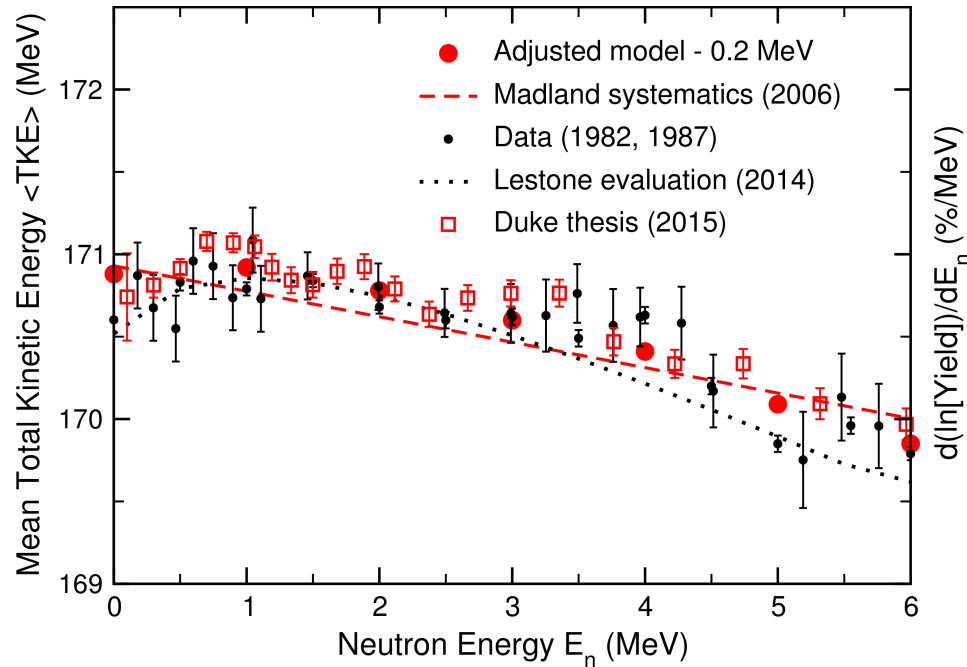
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# 'Adjusted' model; TKE yields; U235 (n,f)



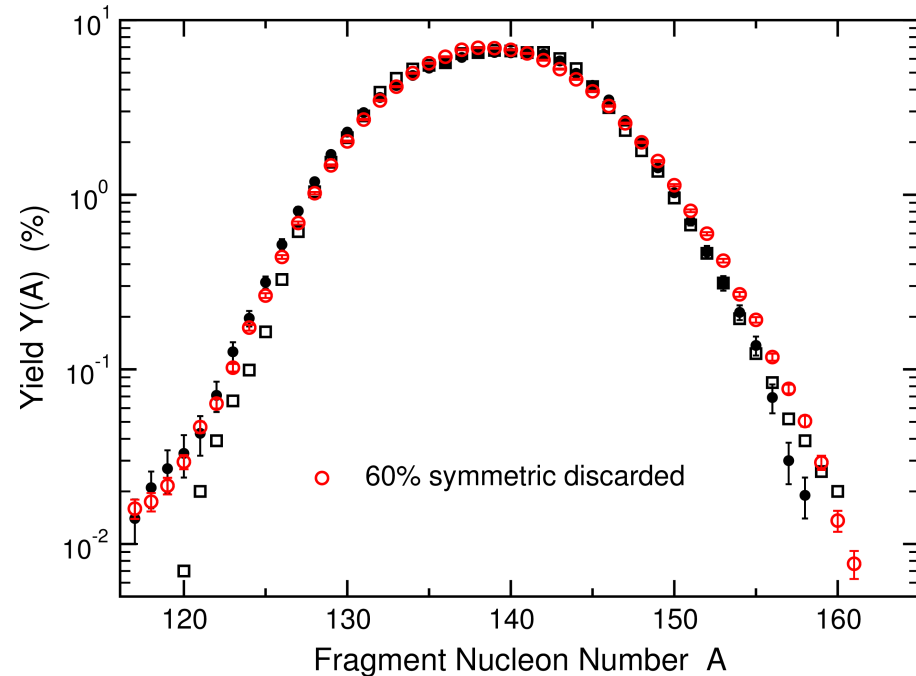
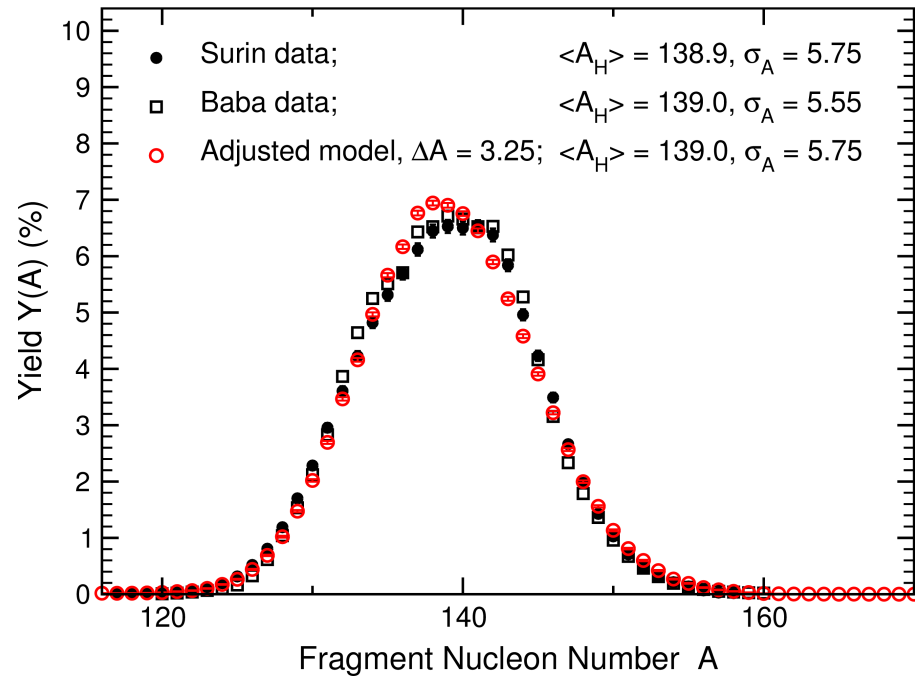
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# 'Adjusted' model; neutron energy dependence



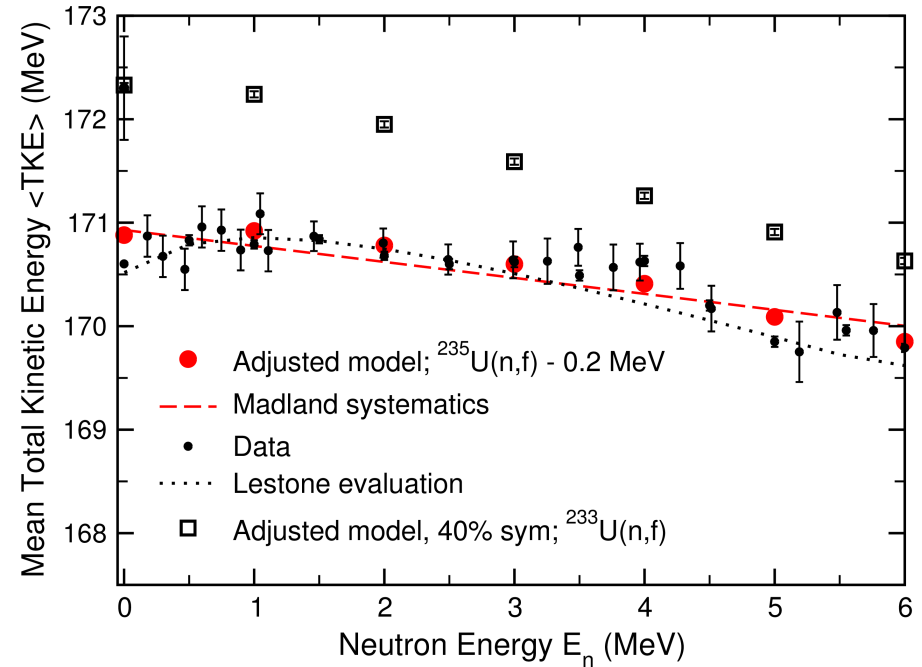
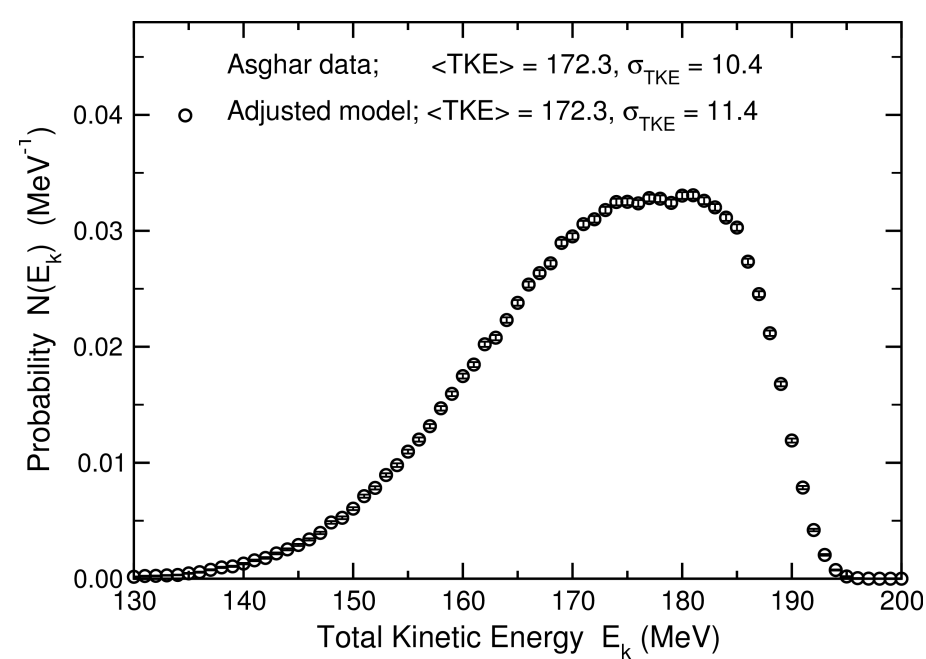
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# 'Adjusted' model; thermal fragment yields; U233 (n,f)



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# 'Adjusted' model; TKE observables; U233 (n,f)



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# Summary

- For the first time, a dynamical physical model of fission predicts most of the major features and some of the quantitative details of fission mass and kinetic-energy distributions.
- The neutron-energy-dependence of average TKE and yields are predicted.

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- This model is already useful for constraining yield and TKE correlations.
- The TKE distributions can be improved by straightforward improvements in the post-scission dynamical model.
- It may be possible to improve detailed yields by relaxing the fixed-isospin assumption.

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# Reference:

- Physical Review C 96, 034603 (2017)

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