

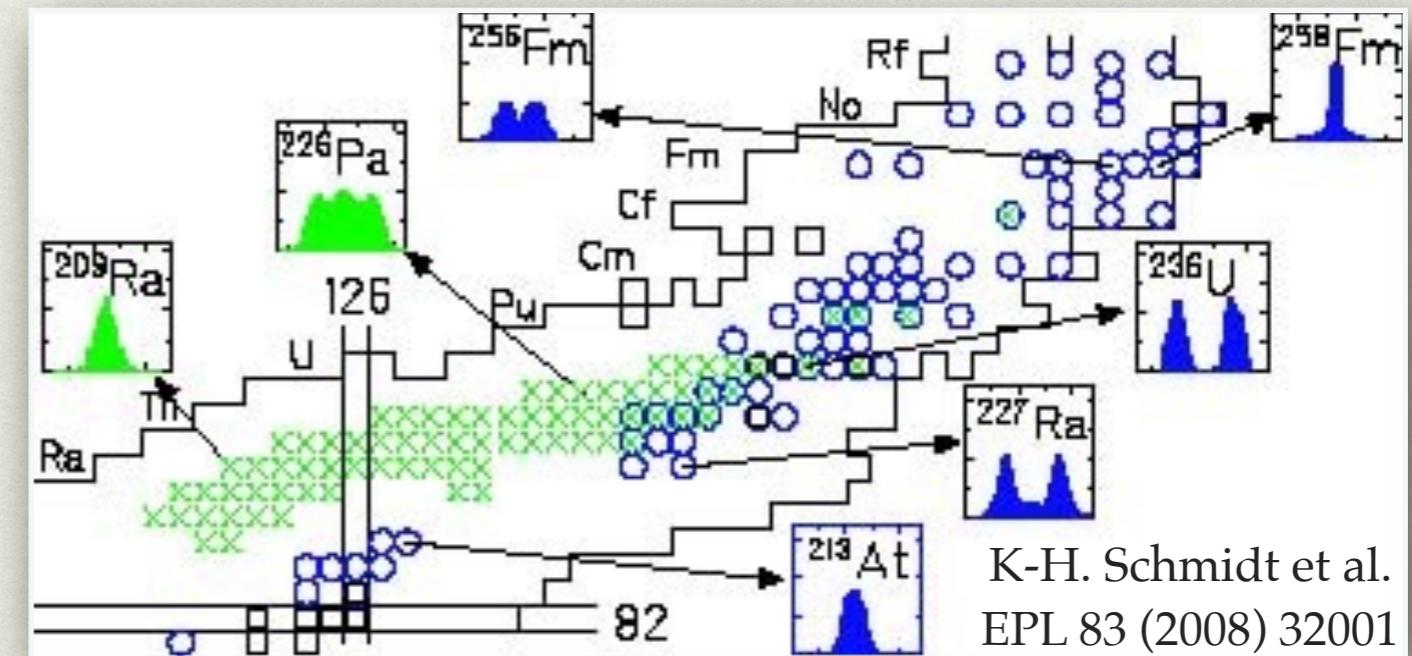
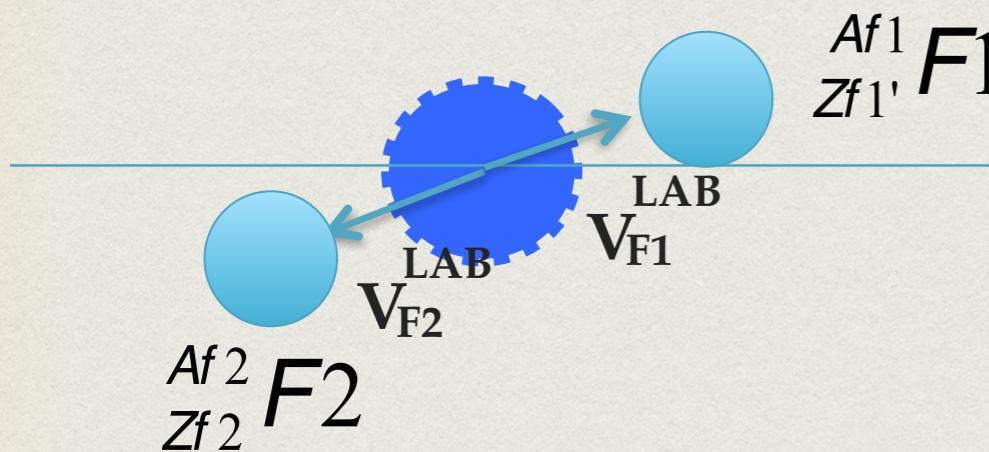
Fragment Distributions of Transfer- and Fusion-Induced Fission in Inverse Kinematics. The Impact of the Excitation Energy.

D. Ramos^{1,2}, M. Caamaño¹, F. Farget³, C. Rodríguez-Tajes³, L. Audouin², J. Benlliure¹, E. Casarejos⁴, E. Clement³, D. Cortina¹, O. Delaune³, X. Derkx^{5,*}, A. Dijon³, D. Doré⁶, B. Fernández-Domínguez¹, G. de France³, A. Heinz⁷, B. Jacquot³, A. Navin³, C. Paradela¹, M. Rejmund³, T. Roger³, M.-D. Salsac⁶, C. Schmitt³

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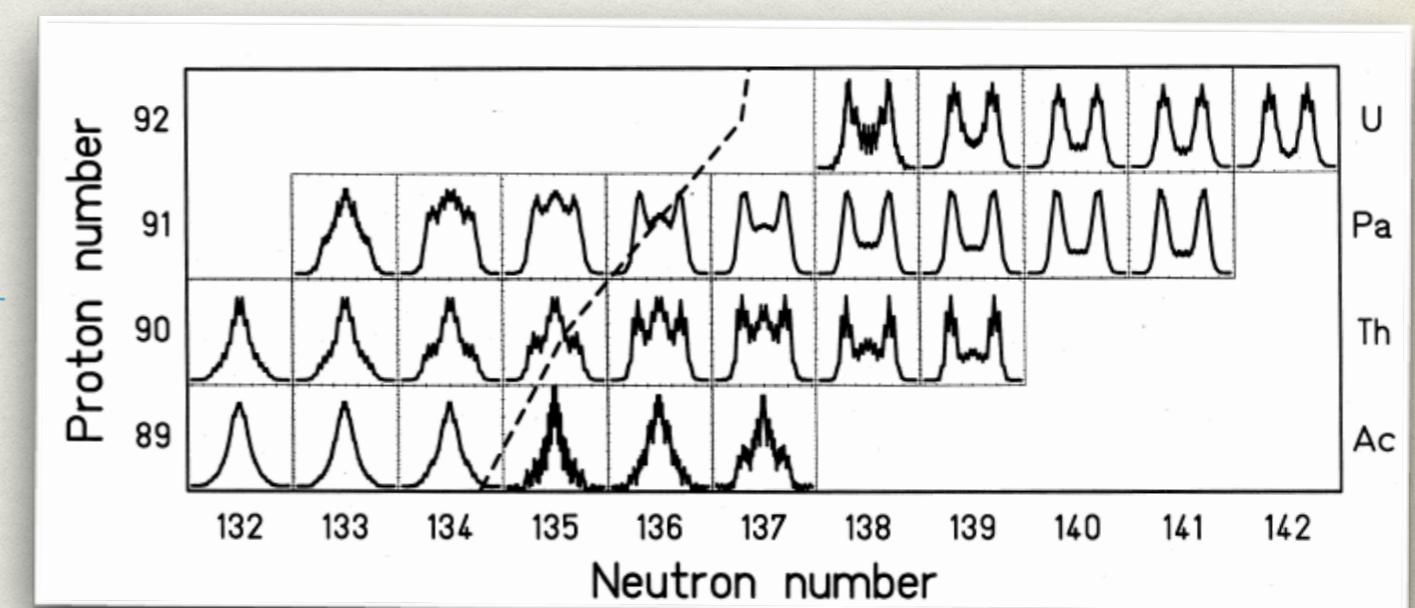
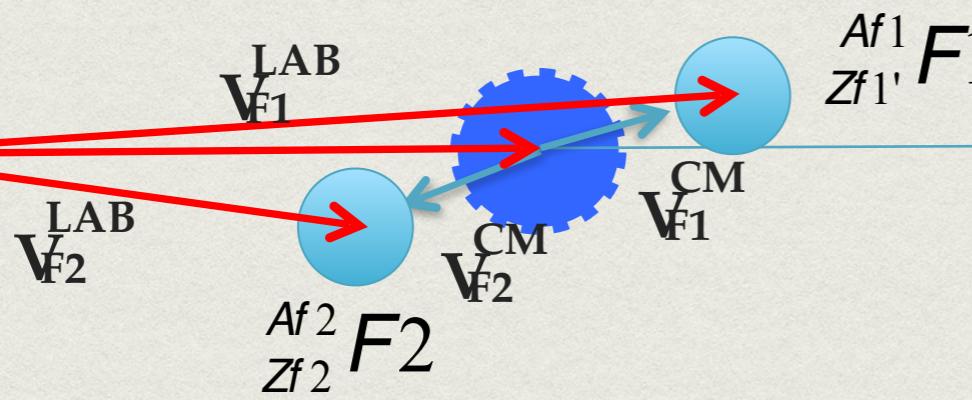
Direct kinematics

- Fragments with low energy
- Mass distribution measurements



Inverse kinematics

- Fragments with high energy
- Direct Z distribution measurements



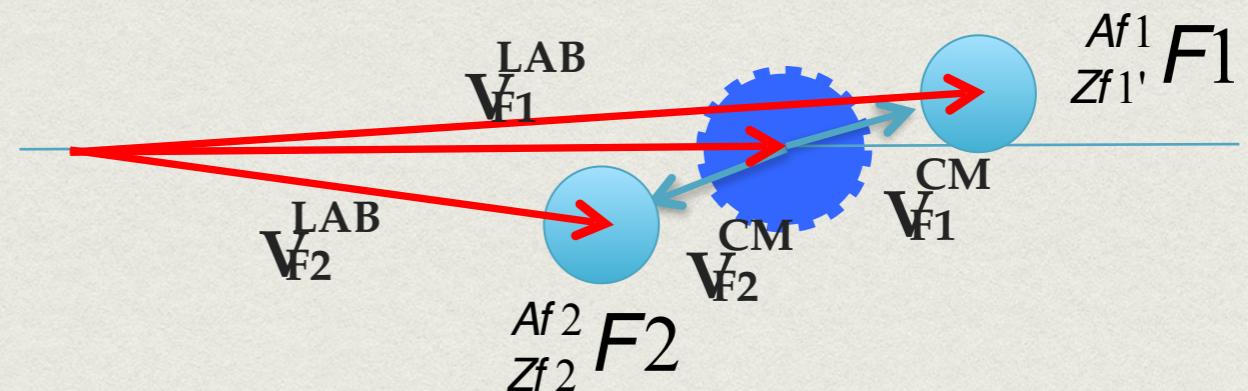
K-H. Schmidt et al. NPA 665 (2000) 221

INVERSE KINEMATICS +
MAGNETIC SPECTROMETER



FULL ISOTOPIC FRAGMENT
DISTRIBUTION

$^{238}\text{U} + ^{12}\text{C}$ @ 6.14 AMeV



Fissioning Systems

^{242}Cf	^{243}Cf	^{244}Cf	^{245}Cf	^{246}Cf	^{247}Cf	^{248}Cf	^{249}Cf	^{250}Cf	^{251}Cf	^{252}Cf
^{241}Bk	^{242}Bk	^{243}Bk	^{244}Bk	^{245}Bk	^{246}Bk	^{247}Bk	^{248}Bk	^{249}Bk	^{250}Bk	^{251}Bk
^{240}Cm	^{241}Cm	^{242}Cm	^{243}Cm	^{244}Cm	^{245}Cm	^{246}Cm	^{247}Cm	^{248}Cm	^{249}Cm	^{250}Cm
^{239}Am	^{240}Am	^{241}Am	^{242}Am	^{243}Am	^{244}Am	^{245}Am	^{246}Am	^{247}Am	^{248}Am	^{249}Am
^{238}Pu	^{239}Pu	^{240}Pu	^{241}Pu	^{242}Pu	^{243}Pu	^{244}Pu	^{245}Pu	^{246}Pu	^{247}Pu	
^{237}Np	^{238}Np	^{239}Np	^{240}Np	^{241}Np	^{242}Np	^{243}Np	^{244}Np			
^{236}U	^{237}U	^{238}U	^{239}U	^{240}U	^{241}U	^{242}U				

Transfer-Fission:

10 n-rich actinides produced with a distribution of E_x below 30 MeV

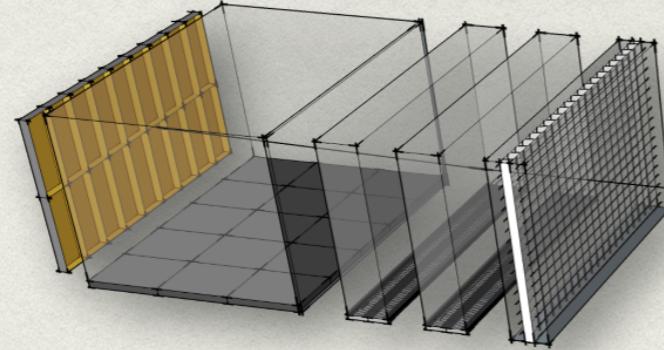
Fusion-Fission:

production of ^{250}Cf with $E_x = 46$ MeV
10 times more likely than any transfer channel

EXPERIMENTAL SETUP

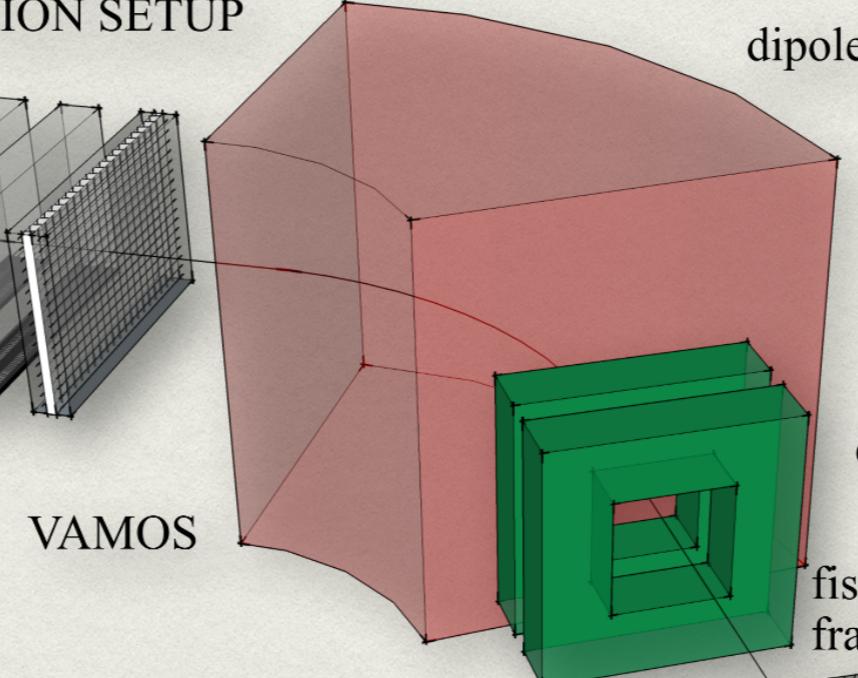
-Fissioning system identification
-Excitation energy reconstruction

FOCAL PLANE
DETECTION SETUP

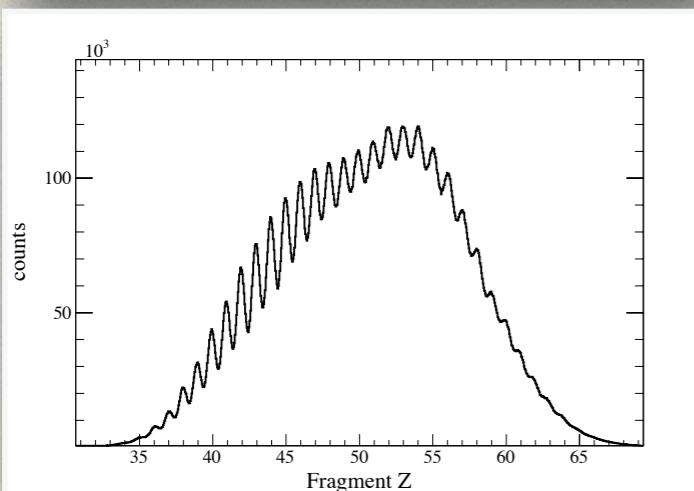
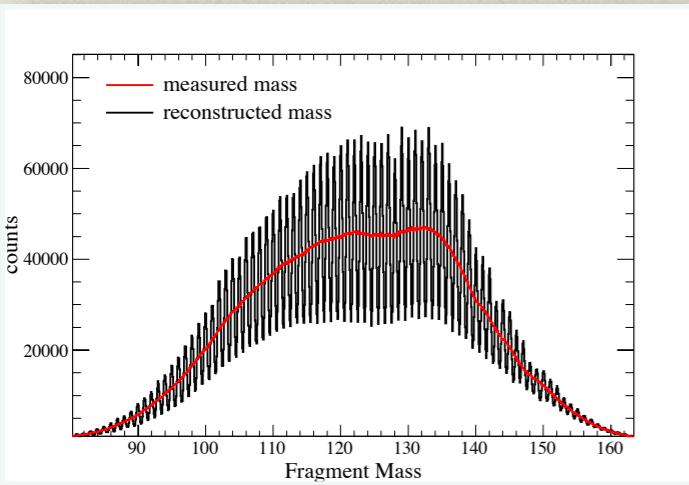


-Fission fragment identification
-Fission kinematics reconstruction

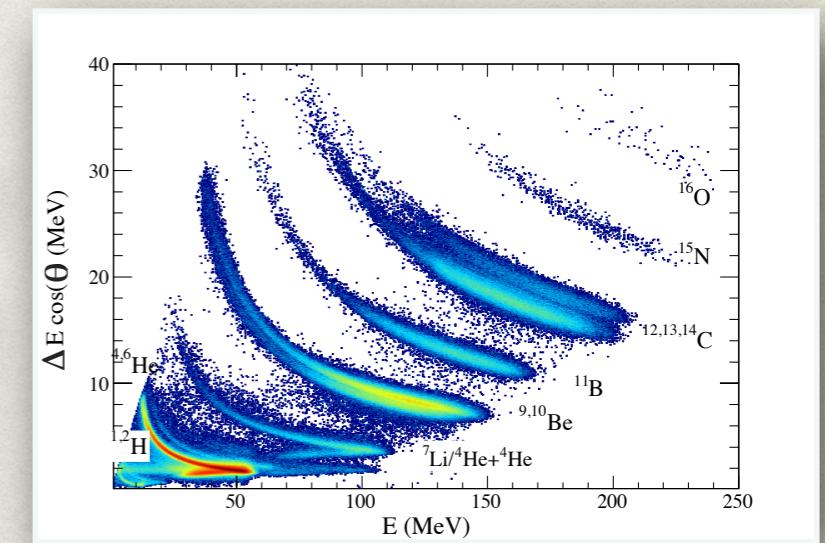
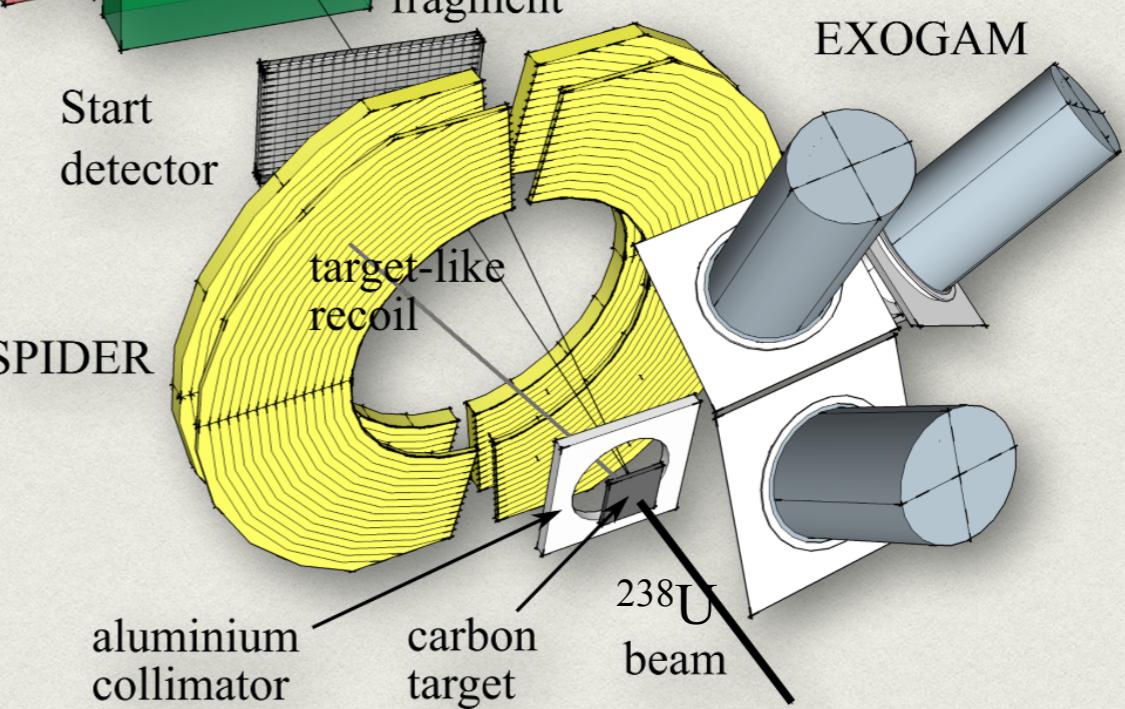
VAMOS



quadrupoles

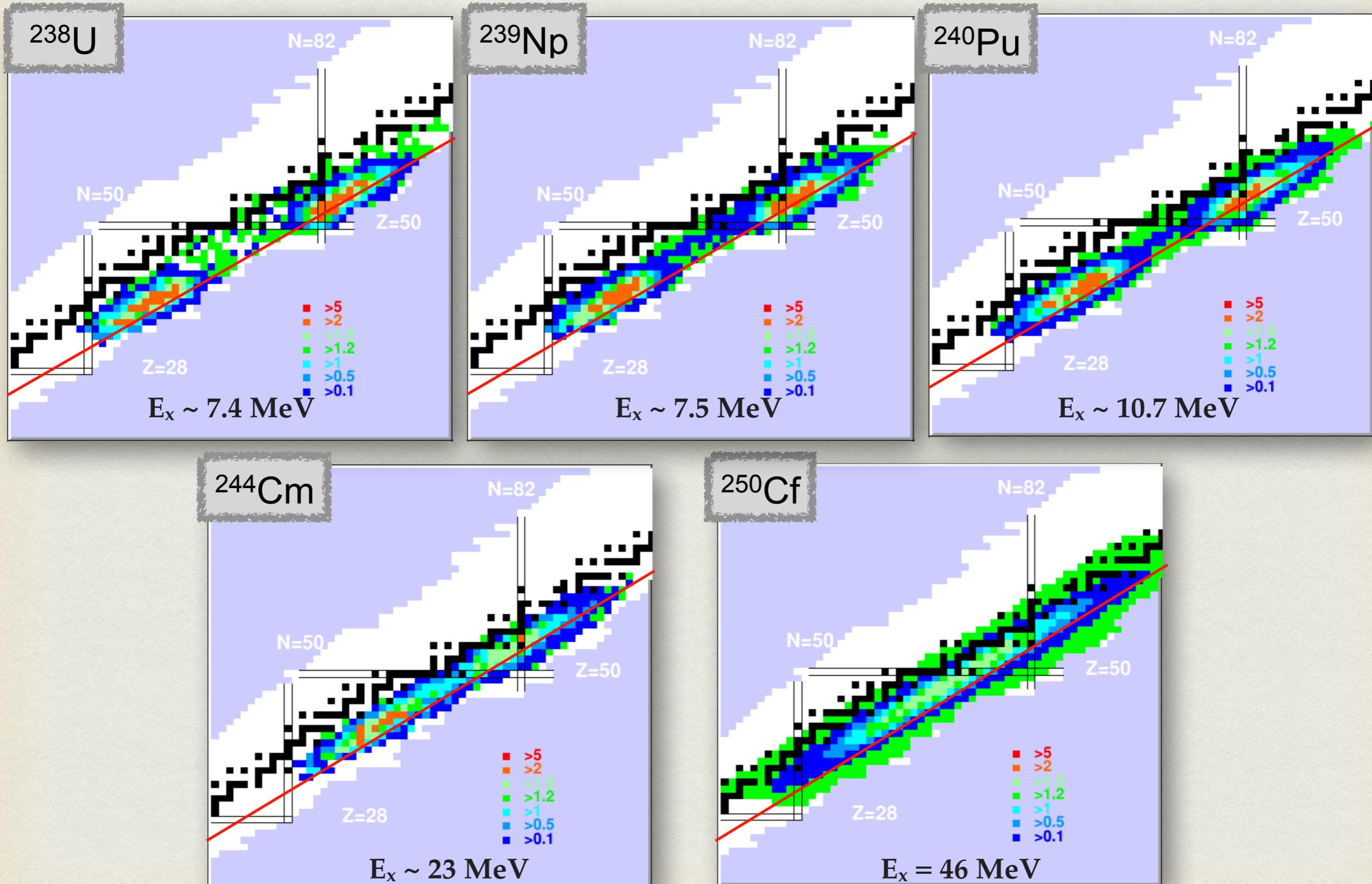


SPIDER

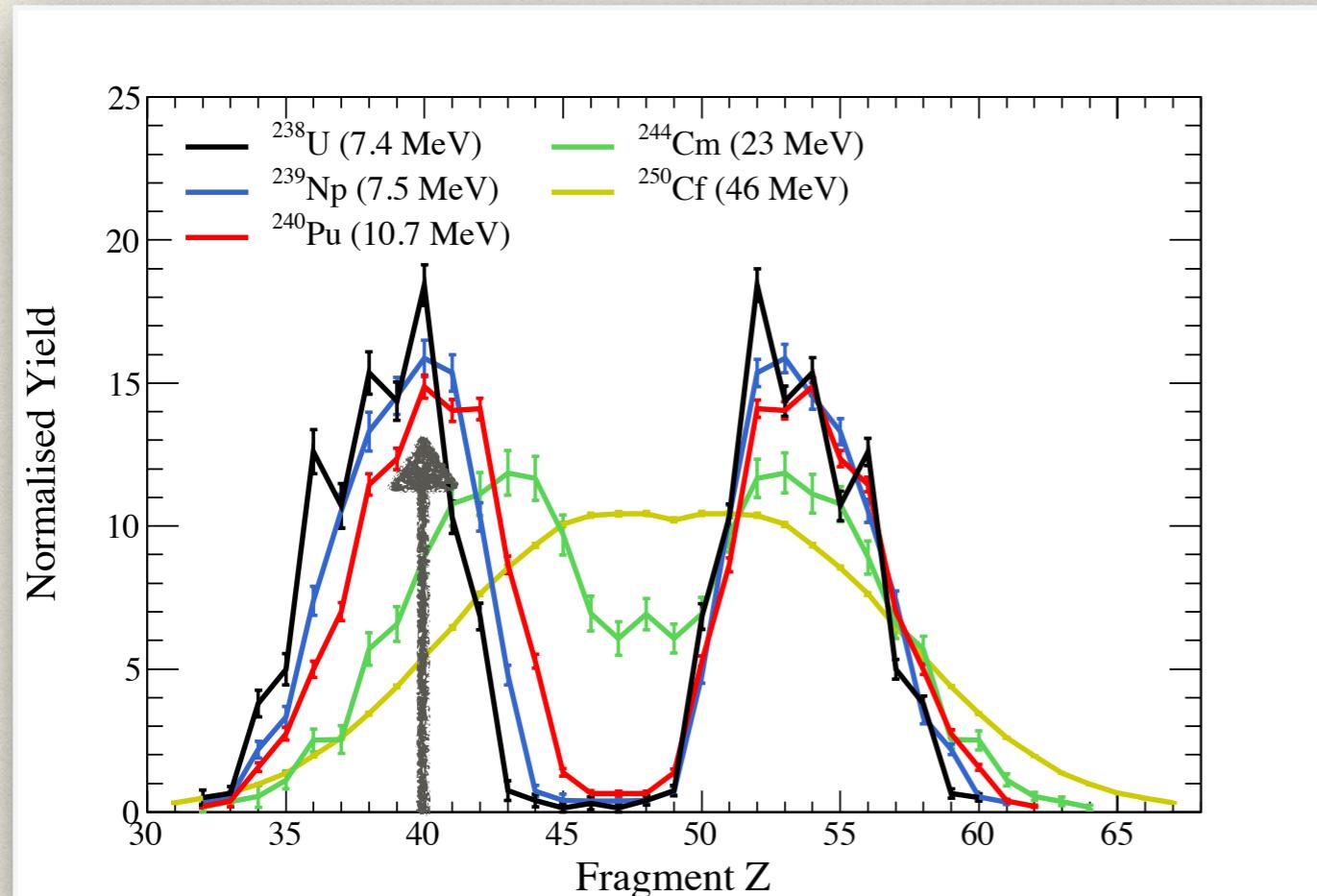


MORE DETAILS IN POSTER

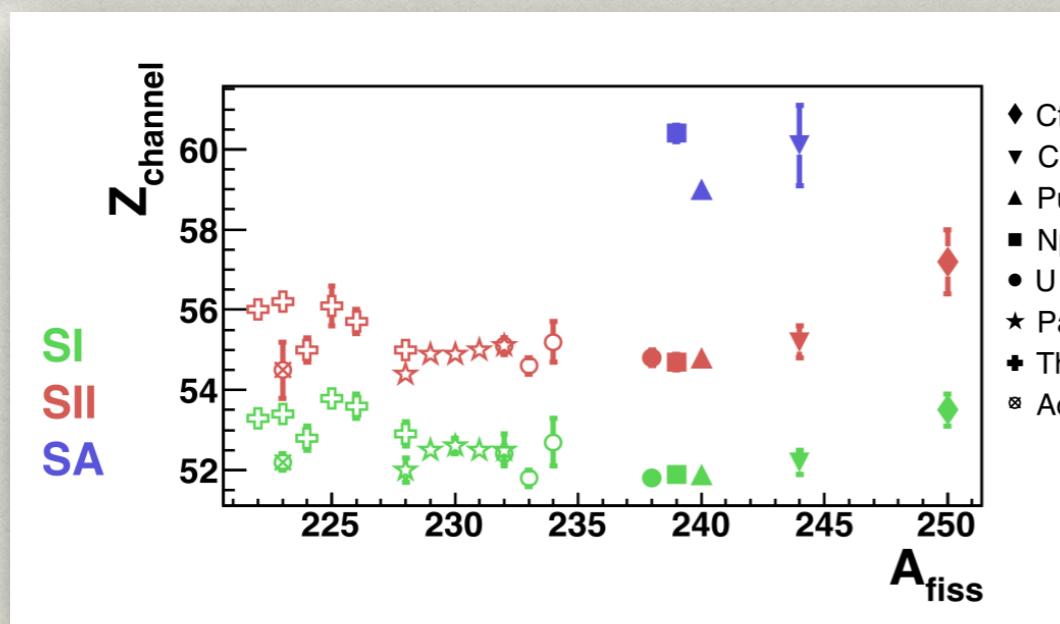
ISOTOPIC FISSION YIELDS



ELEMENTAL FISSION YIELDS



- Stabilization of the heavy fragment
- Light fragment shifted with a constant maximum $Z=40$ at low E_x
- Strong even-odd staggering at low E_x in even- Z systems



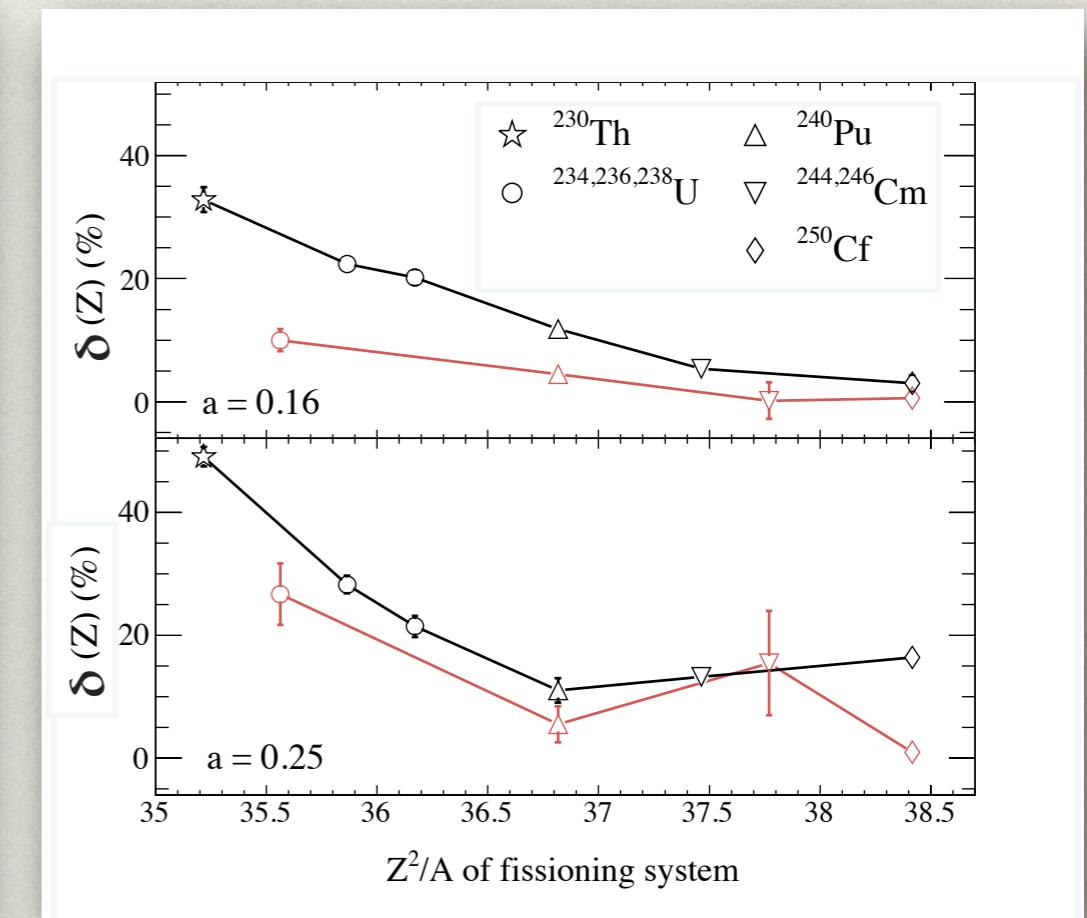
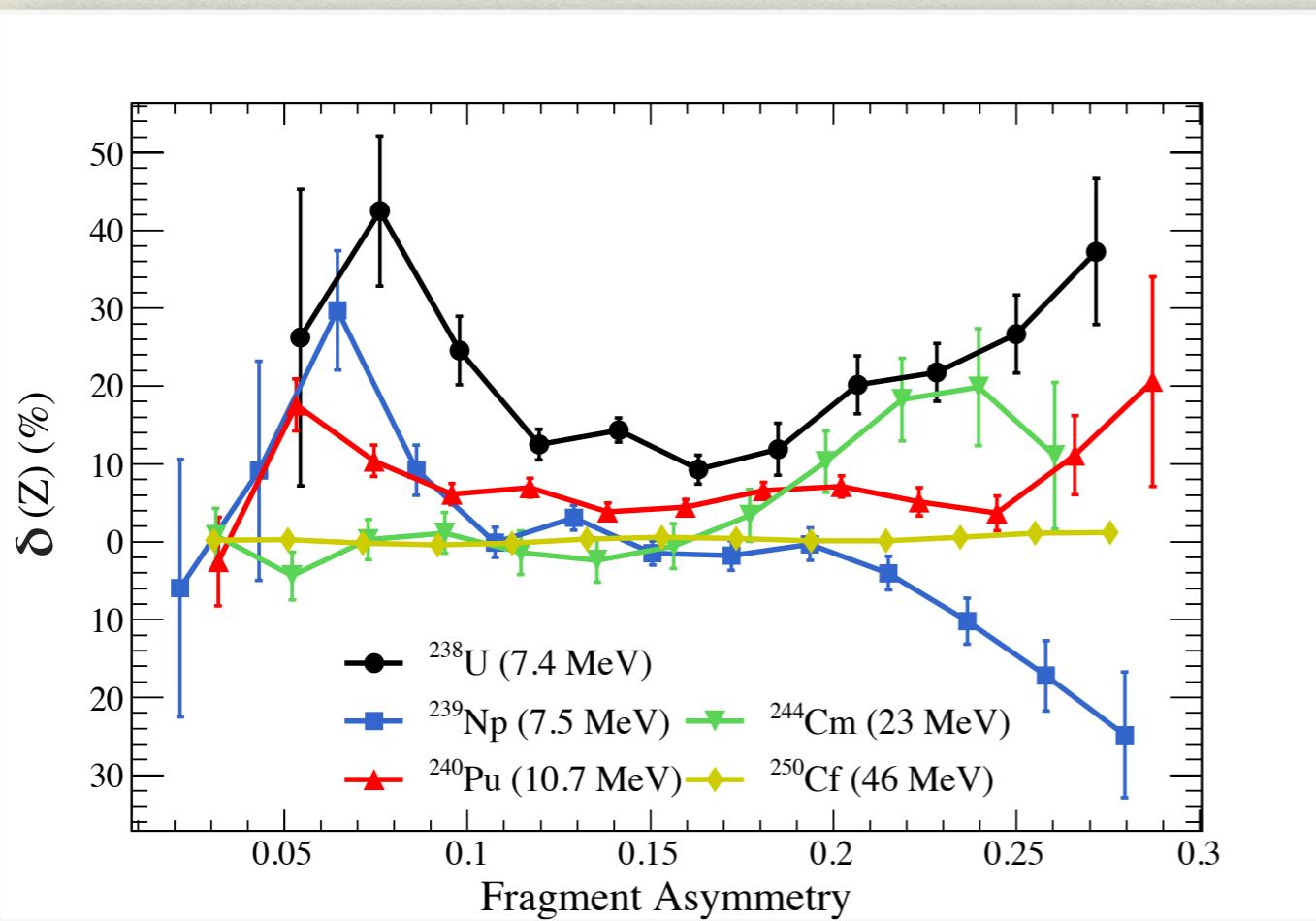
- Asymmetric Fission Modes positions:
(U. Brosa et al. Phys. Reports 197(1990)167)

Standard I : Spherical heavy fragment,
compact scission configuration $Z \sim 52$

Standard II : Deformed heavy fragment $Z \sim 55$

Z EVEN-ODD STAGGERING

$$\delta\left(Z_i + \frac{3}{2}\right) = \frac{(-1)^{Z_i}}{8} \cdot \left[\ln Y(Z_{i+3}) - \ln Y(Z_i) - 3 \left(\ln Y(Z_{i+2}) - \ln Y(Z_{i+1}) \right) \right]$$

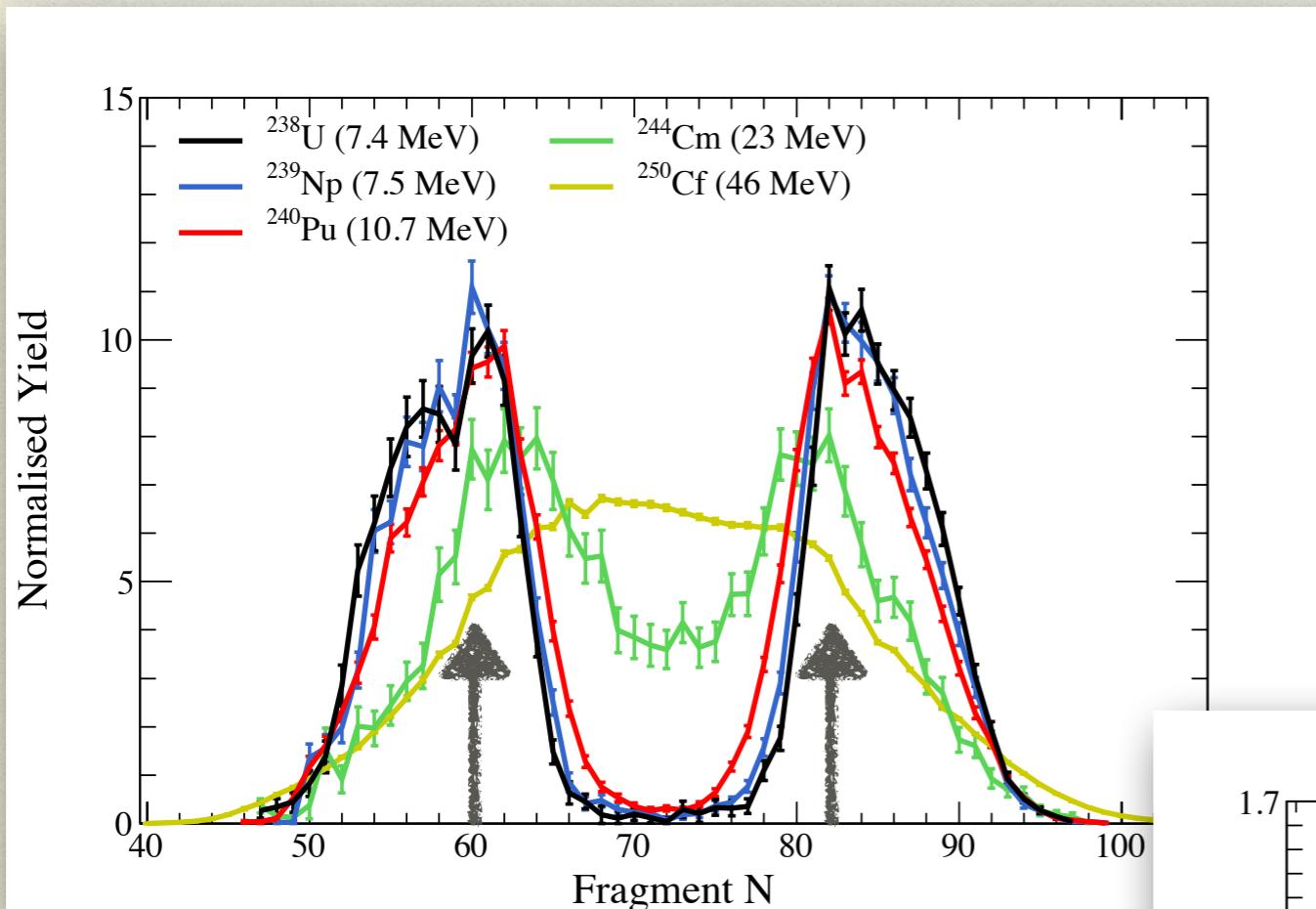


- The even-odd effect increases with the fragment asymmetry with a local maximum around $Z \sim 50$
- Systems at higher E_x show lower even-odd effect
- Odd-Z nuclei ^{239}Np show opposite behavior

Even-Z nuclei (red) compared with thermal-neutron induced fission (black) at different asymmetries

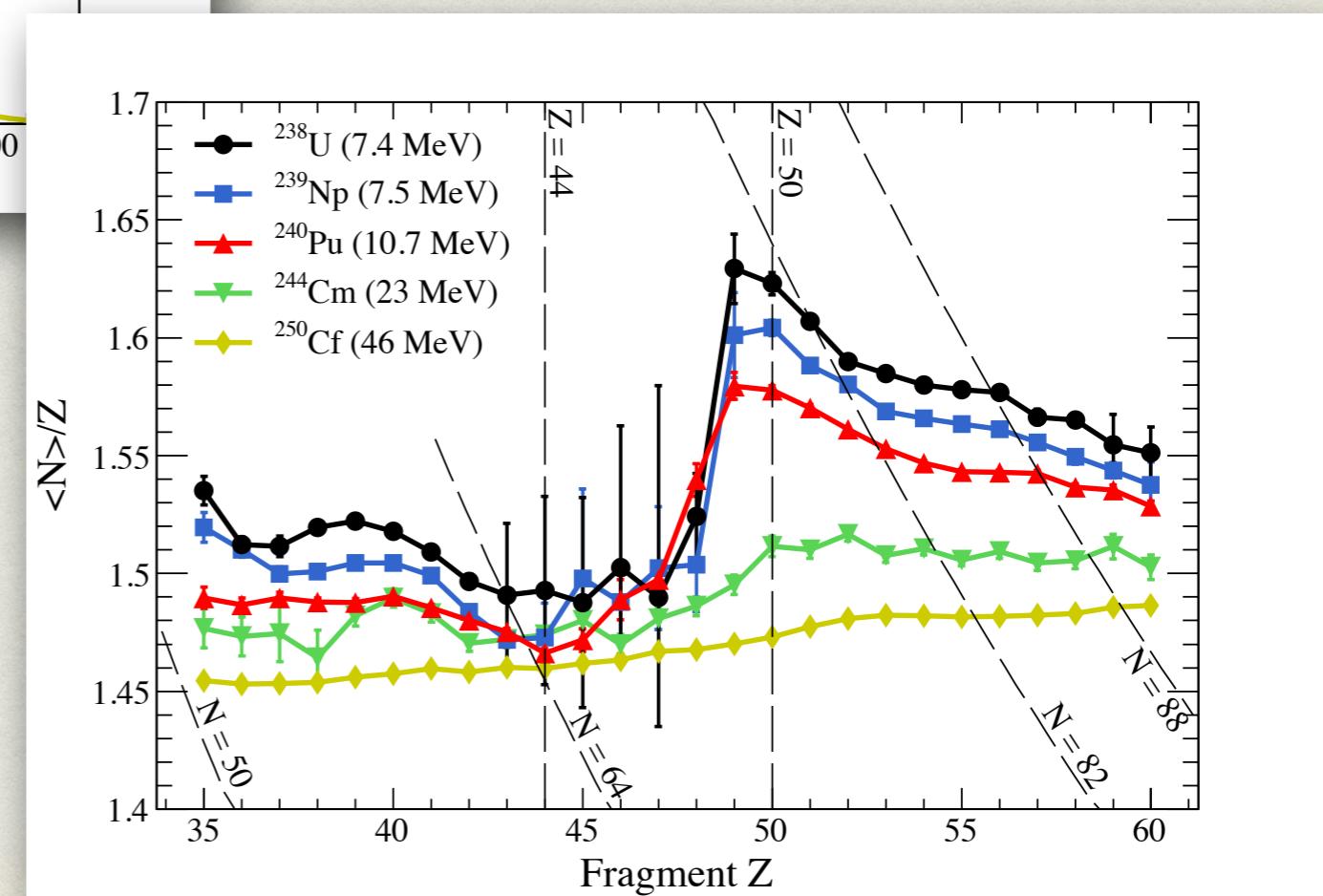
- For central asymmetry, the impact of Z^2/A dominates over E_x difference

NEUNTRON DISTRIBUTIONS AND NEUTRON EXCESS

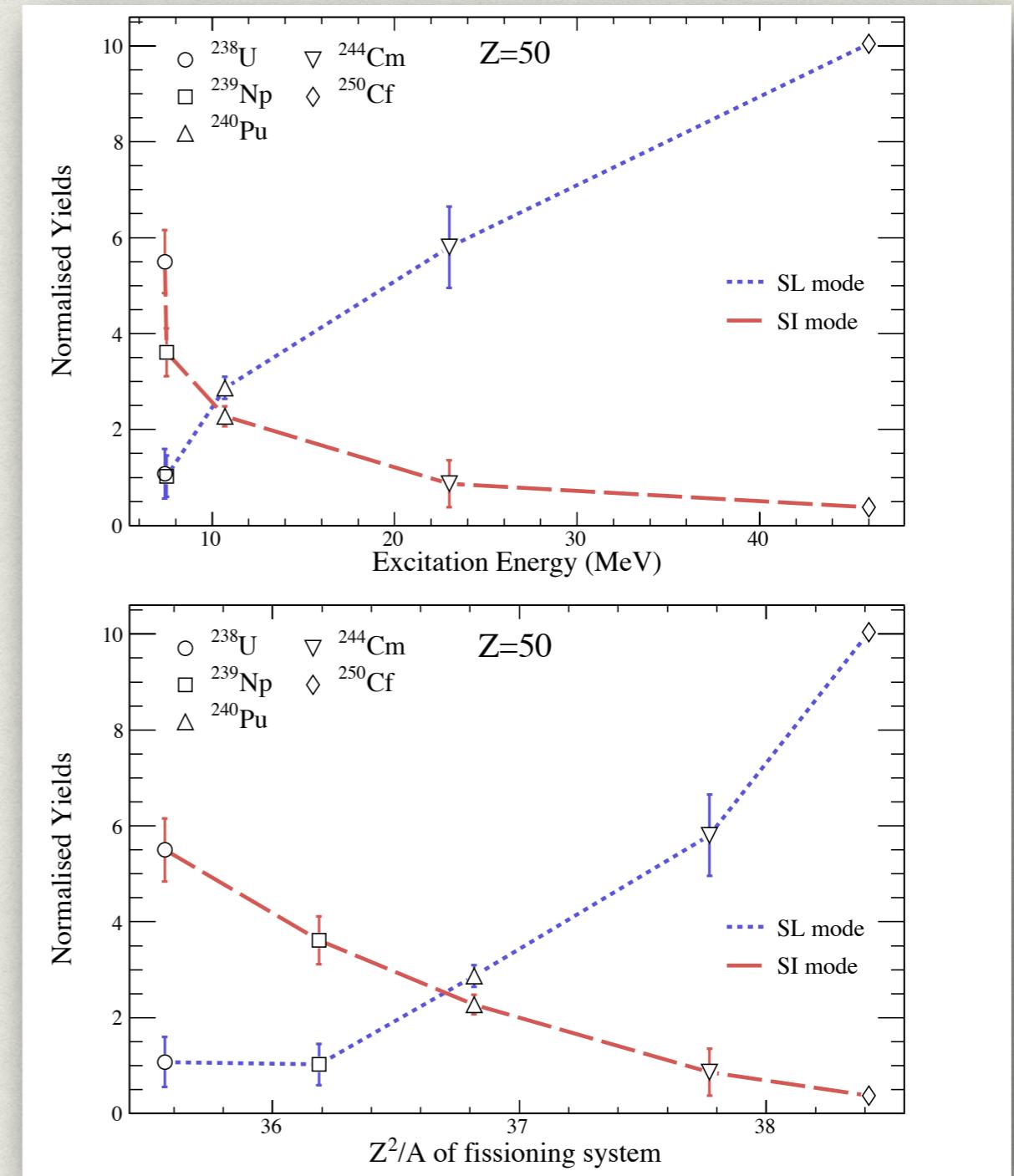
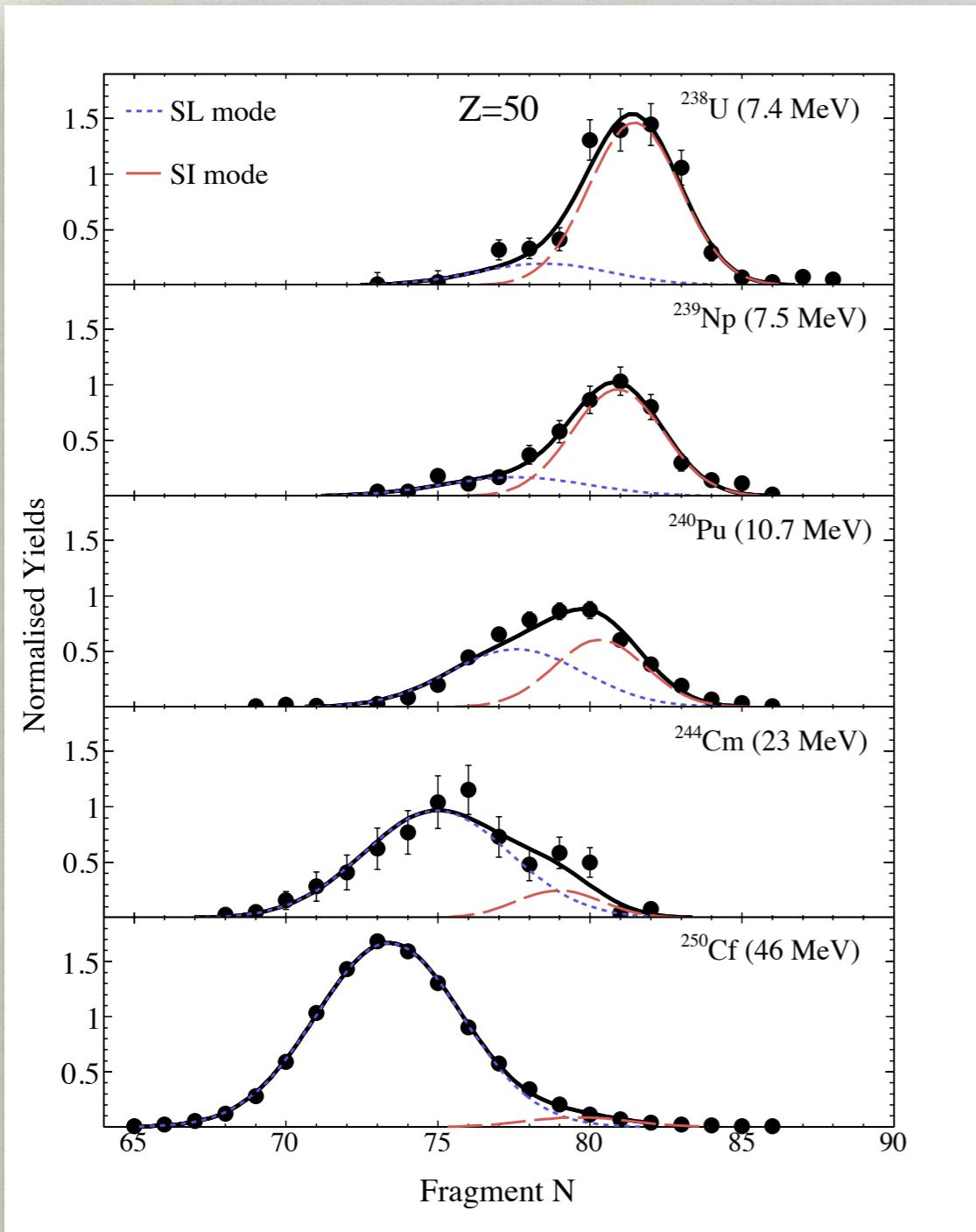


- Charge polarization with the heavy fragment more n-rich than the light fragment
- High $\langle N \rangle / Z$ at $Z \sim 50$ at low E_x , decreasing in systems at higher E_x
- Tendency to the double magic nucleus ^{132}Sn
- Local maximum at $Z \sim 40$

- Post neutron evaporation
- Maximum in the heavy fragment at $N=82$
- Unexpected maximum in the light fragment $N \sim 60$
- The no-flat plateau of ^{250}Cf reflects the effect of the increasing neutron evaporation



ISOTOPIC DISTRIBUTION IN Z=50

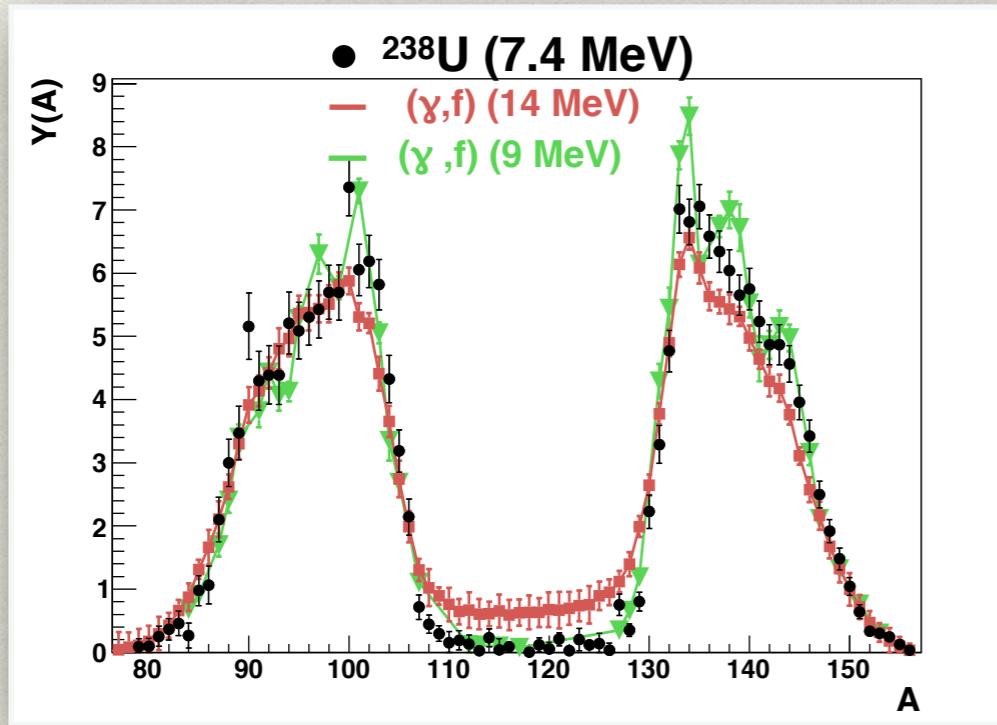


- The isotopic distribution of $Z=50$ is mainly described by two fission modes: SL, related to the symmetric fission and SI, related to spherical closed shells ($N=82$, $Z=50$)
- The E_x feeds the **SL mode** while the **SI mode shows an interplay between the E_x and the fissility parameter**

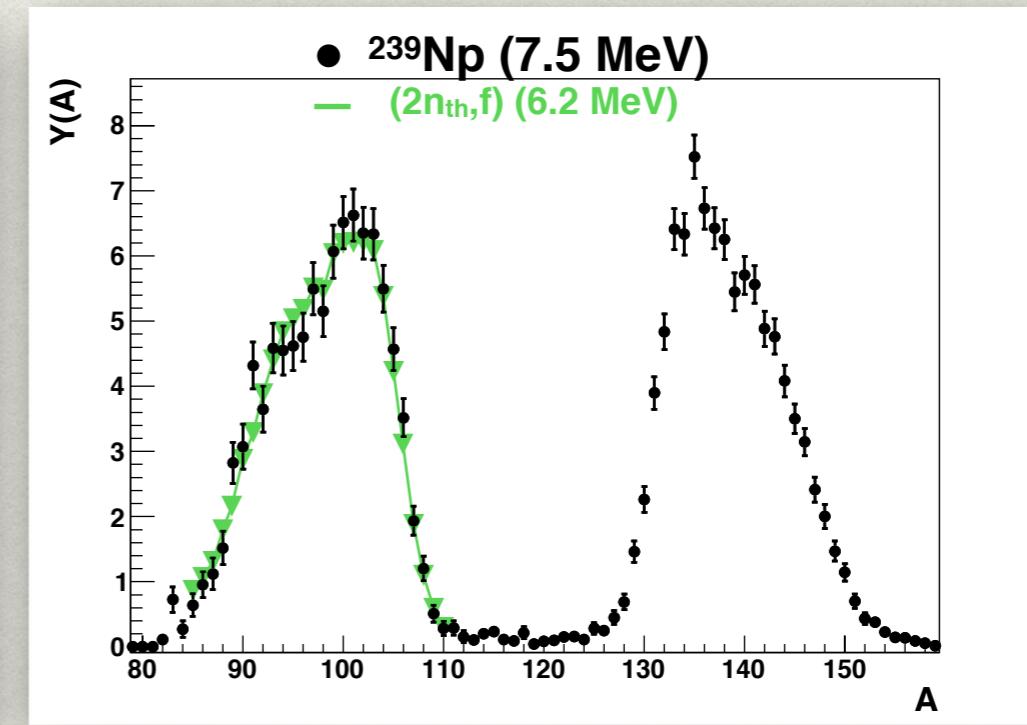
CONCLUSIONS

- The use of inverse kinematics with the VAMOS spectrometer allowed the characterization of fissioning system in terms of Z, A, and E_x , and the isotopic identification of their full fragment distribution.
- The measured elemental and neutron yield distributions show a dominant effect of structure at low E_x , while at higher E_x , the macroscopic features become more relevant. Maxima of production are observed in Z=40 and N=60,82.
- The local even-odd effect is found to increase around Z~50 splits, consistent with a larger binding energy, and at high asymmetry.
- The correlation between N and Z also reveals the effect of structure: a charge polarization with a saw-tooth behavior appears with a maximum governed by ^{132}Sn . Again, increasing E_x reduces this structure effect, but mainly on the heavy fragment.
- The isotopic distributions of Z=50 suggest that the SL mode is mainly governed by E_x while the SI shows also the impact of the fissility.

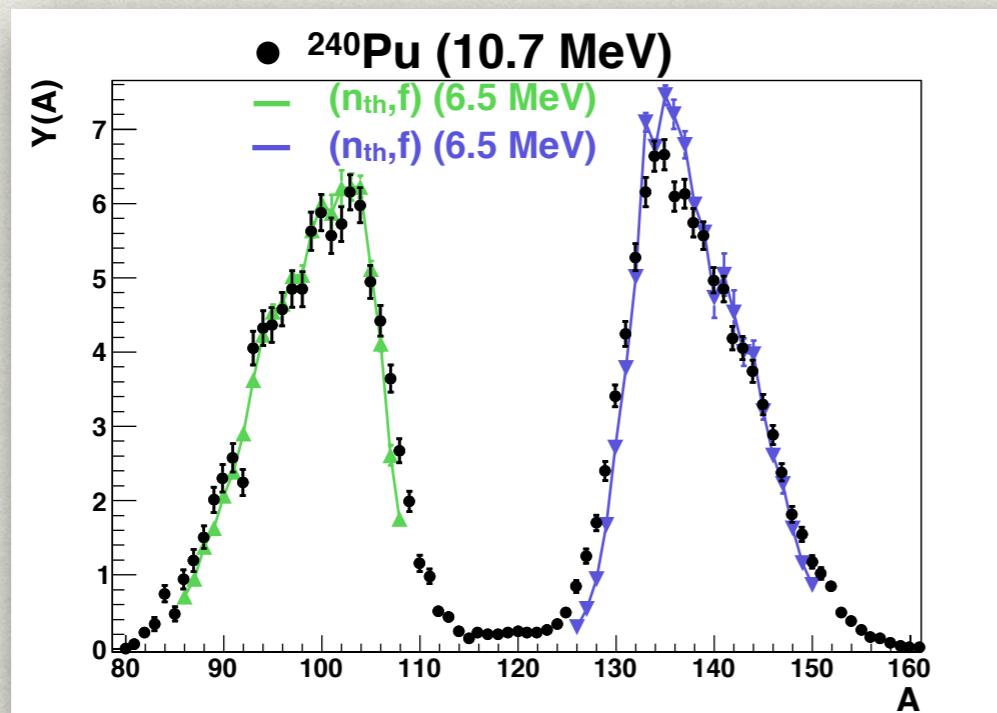
MASS FISSION YIELDS



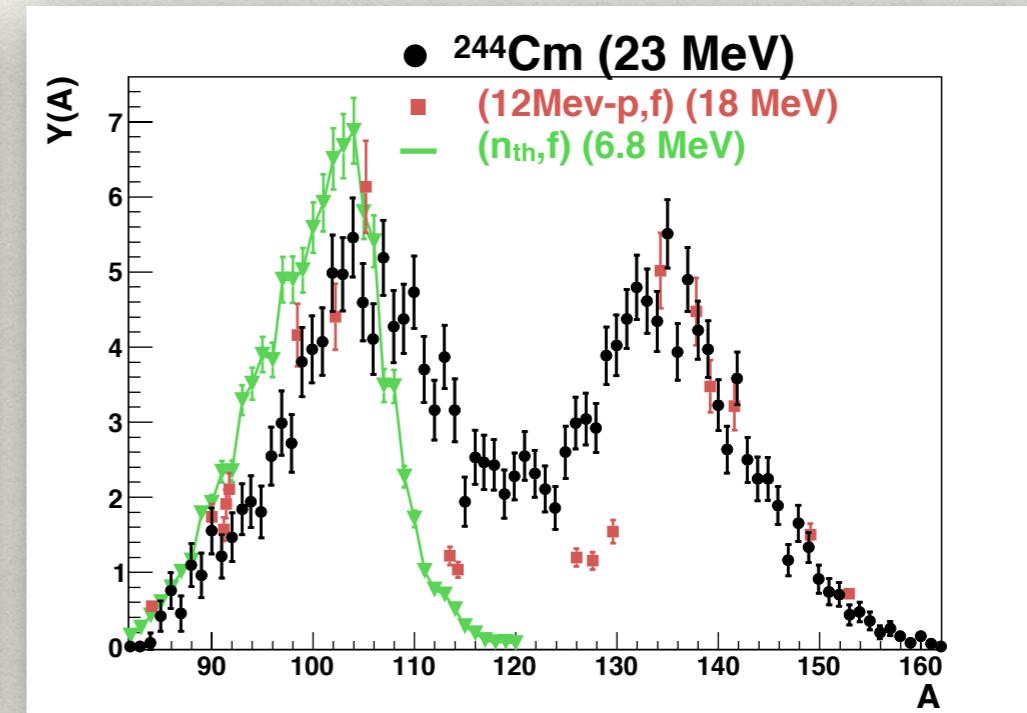
E. Pellereau et al. PRC 95 (2017) 054603
 H. Naik et al. EPJA 49 (2013) 94



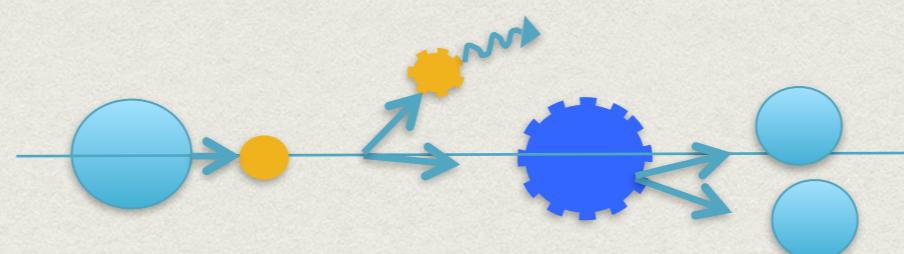
G. Martinez et al. NPA 515 (1990) 433



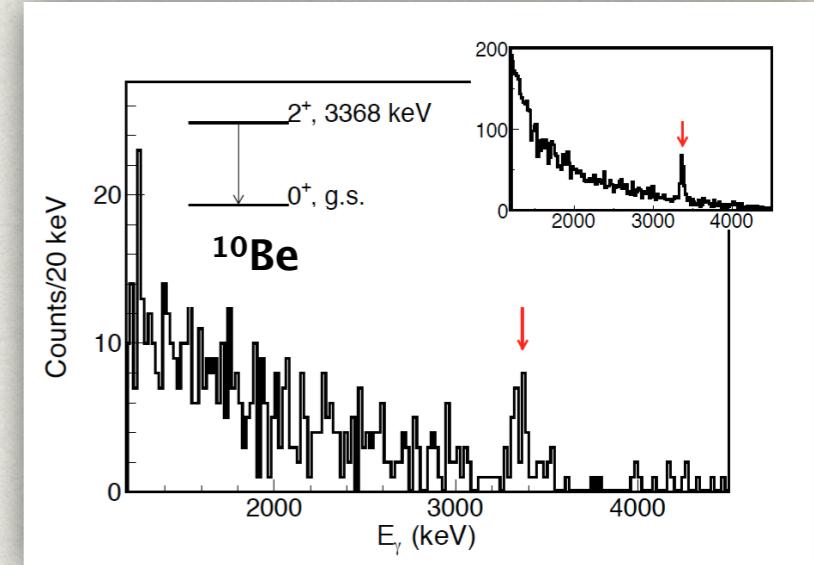
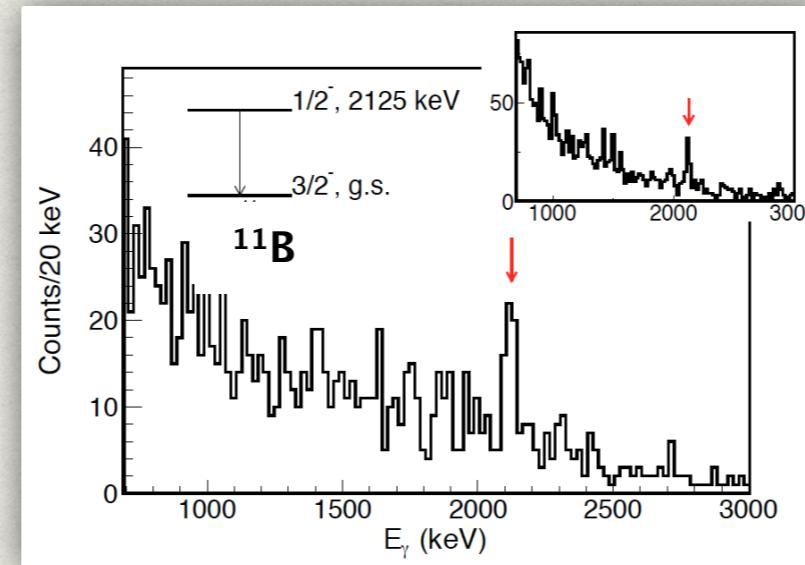
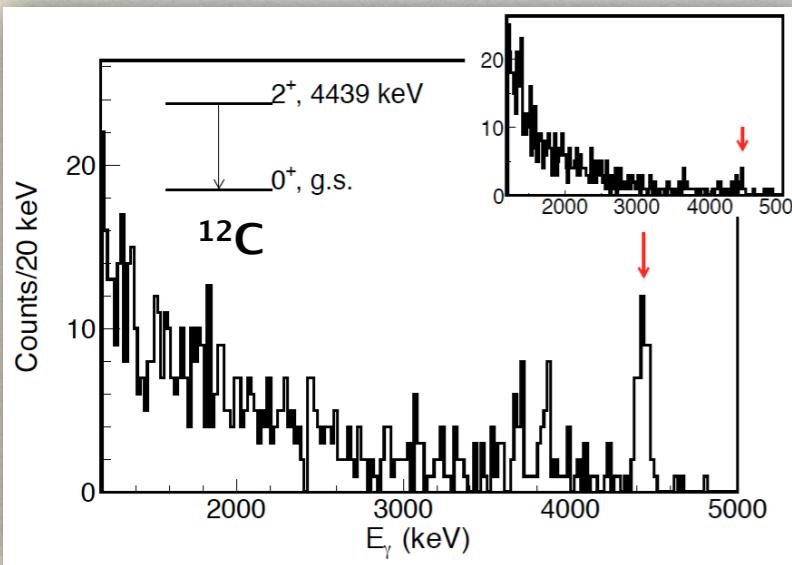
C. Schmitt et al. NPA 430 (1984) 21
 A. Bail et al. PRC 84 (2011) 034605



T. Ohtsuki et al. PRC 40 (1989) 2144
 I. Tsekhanovich et al. PRC 70 (2004) 044610



EXOGAM detector allow us to evaluate the excitation probability of the target-like nuclei



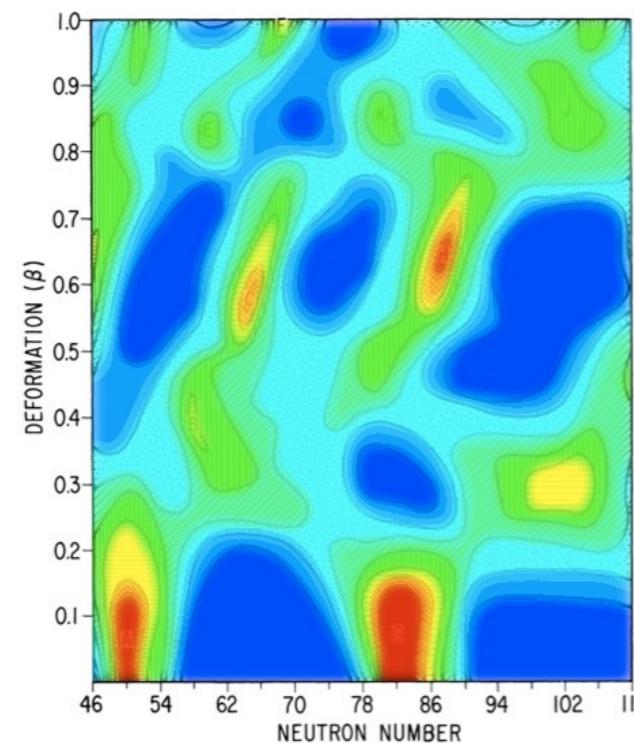
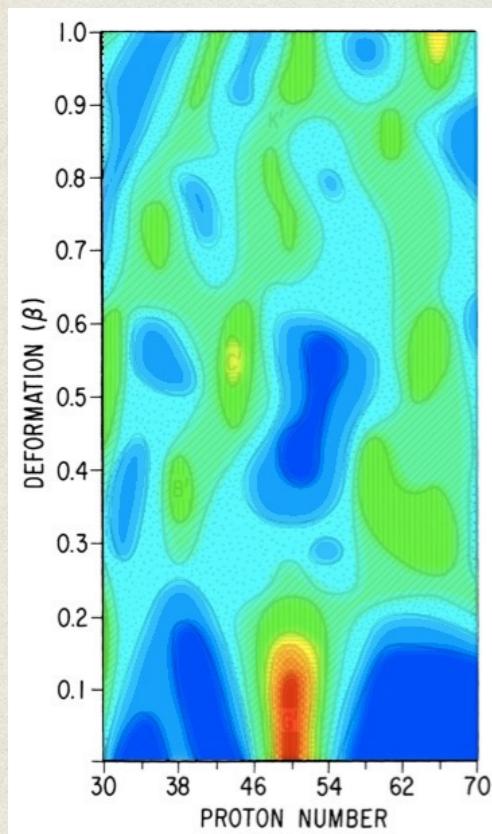
γ -rays measurements show excited states in ^{12}C , ^{11}B and ^{10}Be in coincidence with fission with $P_\gamma = 0.12\text{-}0.14$

Wilkins Scission Point

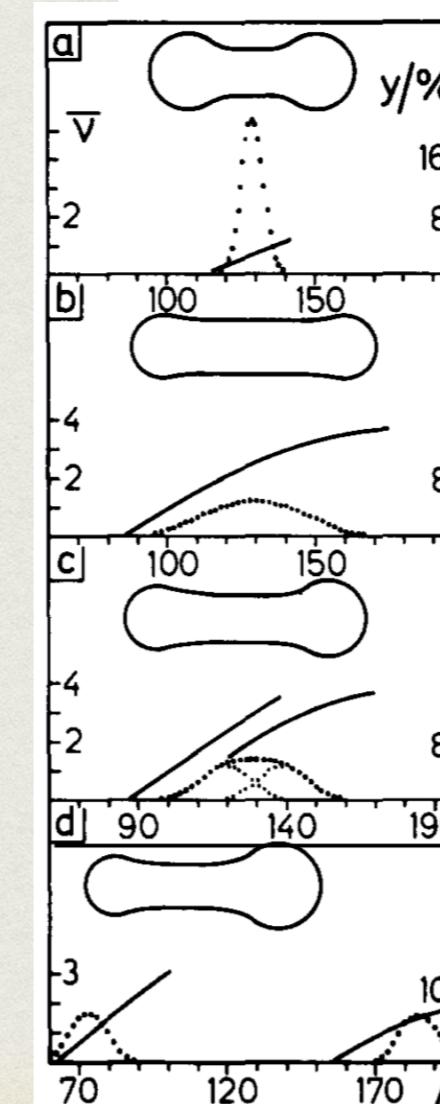
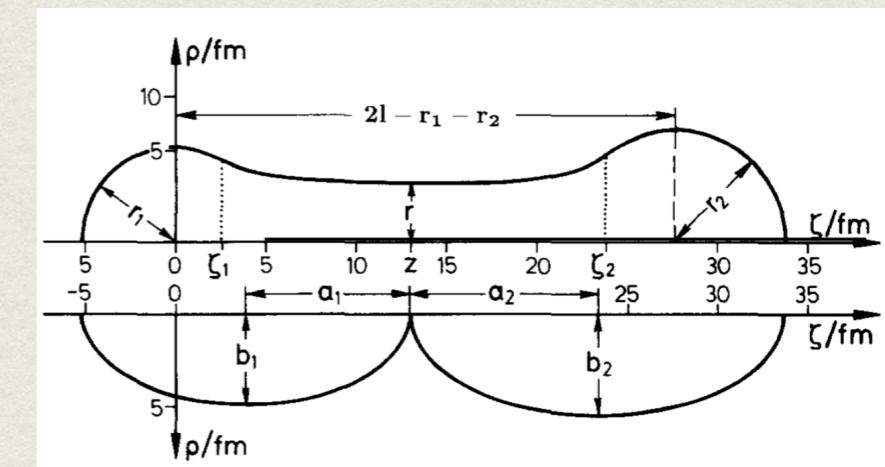
$$R = R_0 \left(1 + \frac{2}{3} \beta\right)$$

$$r = R_0 \left(1 - \frac{1}{3} \beta\right)$$

$$\begin{aligned} V &= V_1^{LDM} + V_2^{LDM} \\ &+ S_1^N + S_2^N + S_1^Z + S_2^Z \\ &+ P_1^N + P_2^N + P_1^Z + P_2^Z \\ &+ V_{12}^{Coulomb} + V_{12}^{Nuclear} \end{aligned}$$



Brosa Random neck rupture



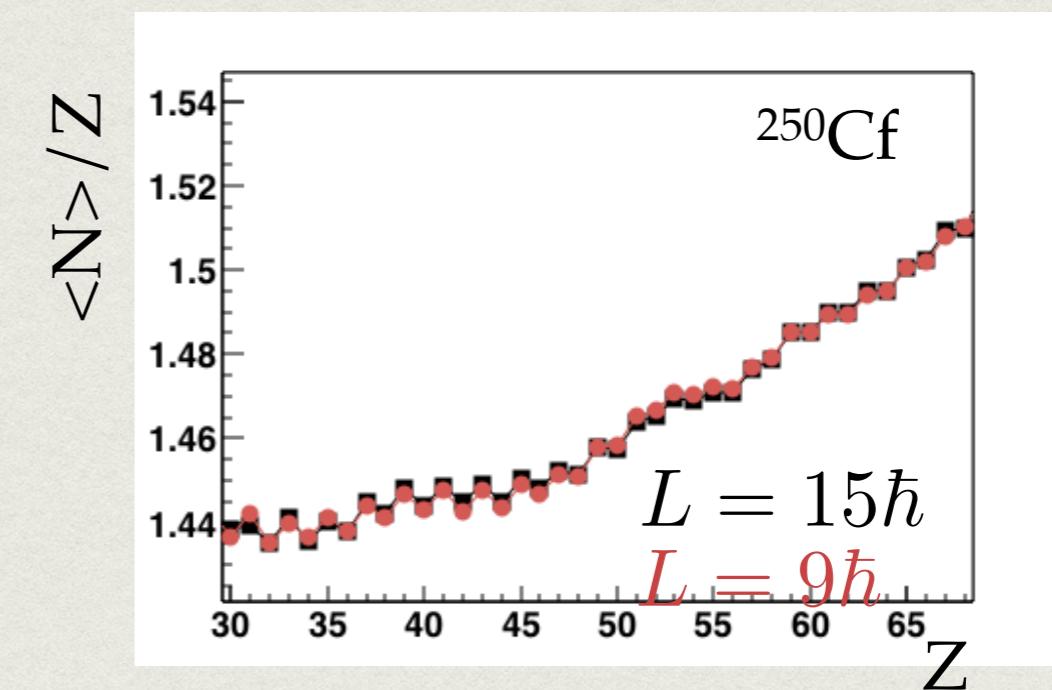
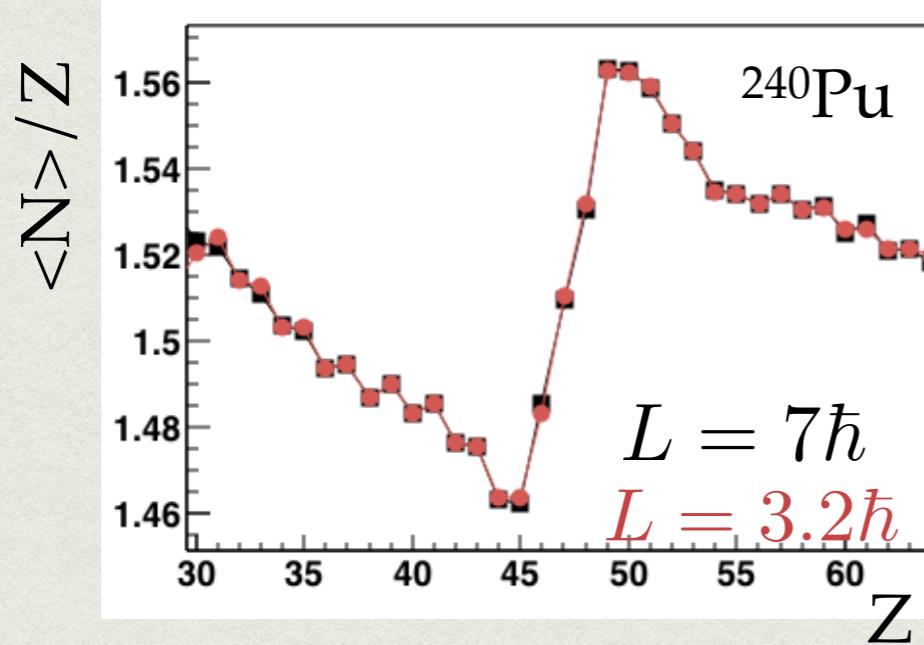
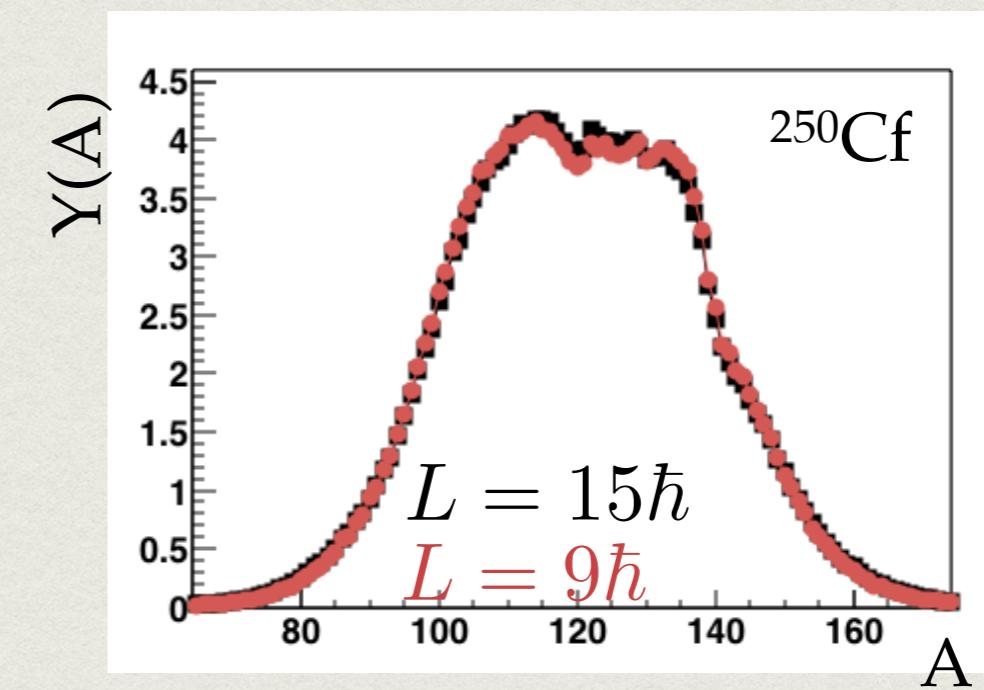
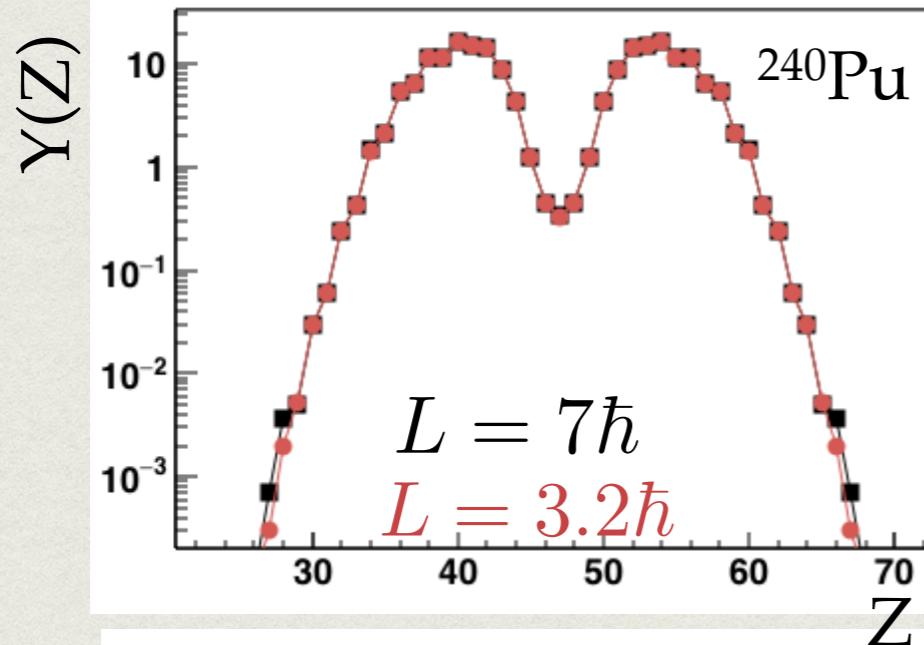
Super short channel

Super long channel

Standard I channel

Super asymmetric channel

GEF calculation



$$\langle \nu \rangle = 3.47$$

$$\langle \delta \rangle = 9.51$$

$$\langle \nu \rangle = 3.46$$

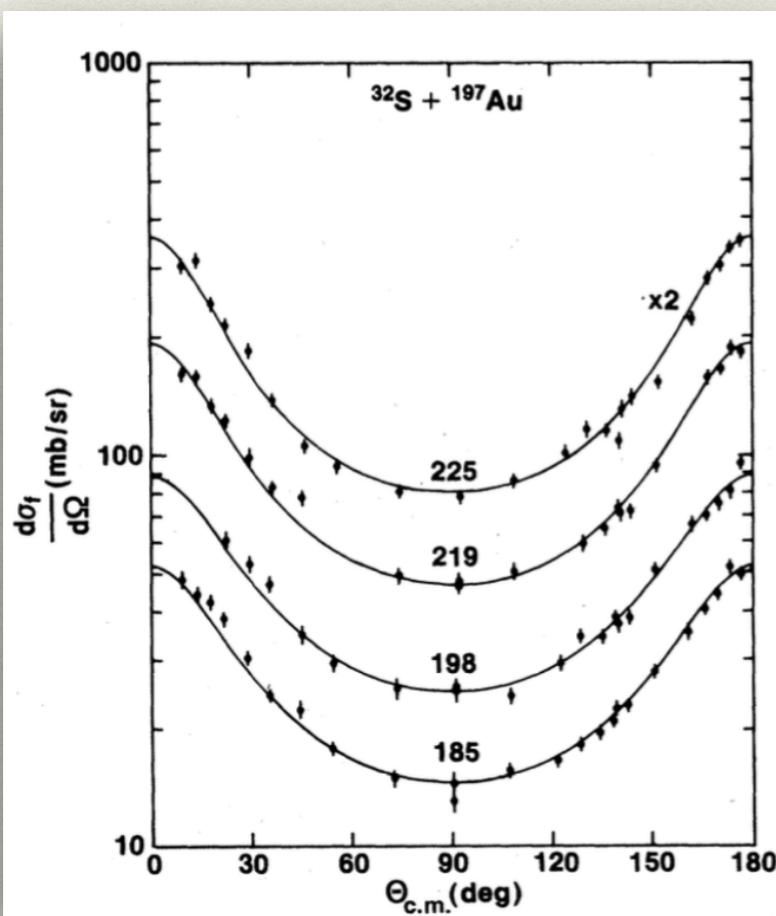
$$\langle \delta \rangle = 9.66$$

$$\langle \nu \rangle = 8.79$$

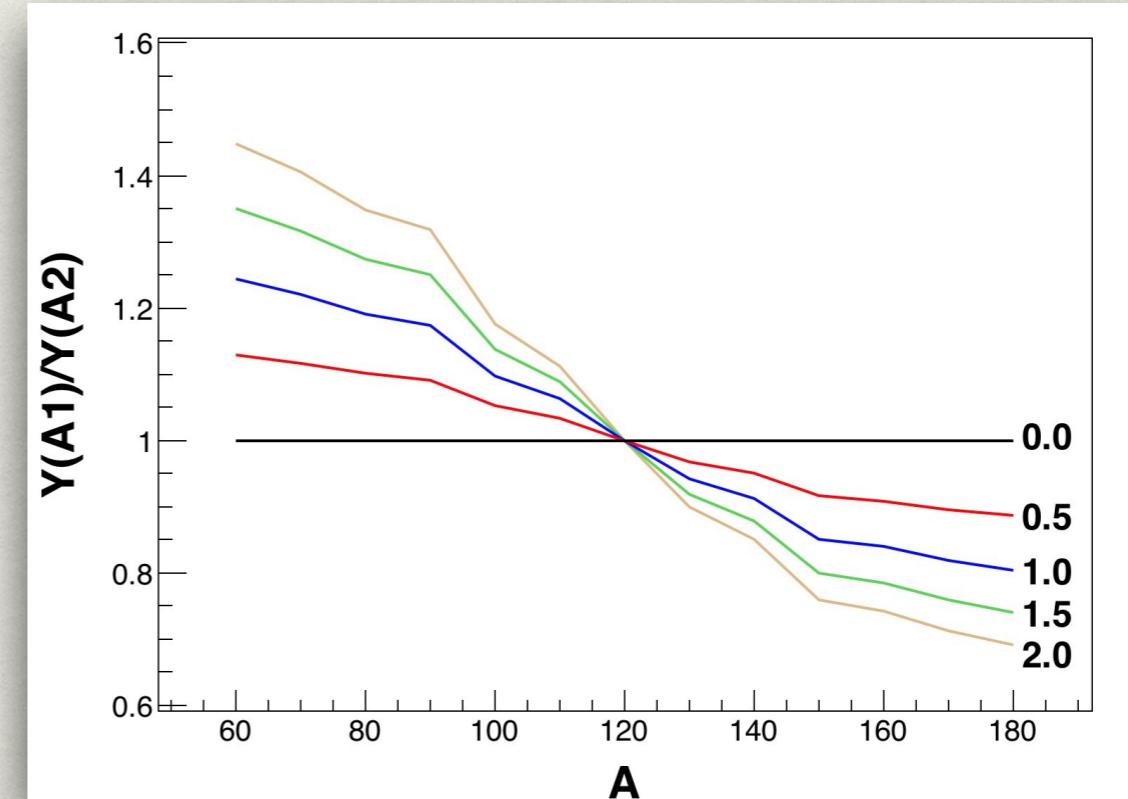
$$\langle \delta \rangle = 3.91$$

$$\langle \nu \rangle = 8.79$$

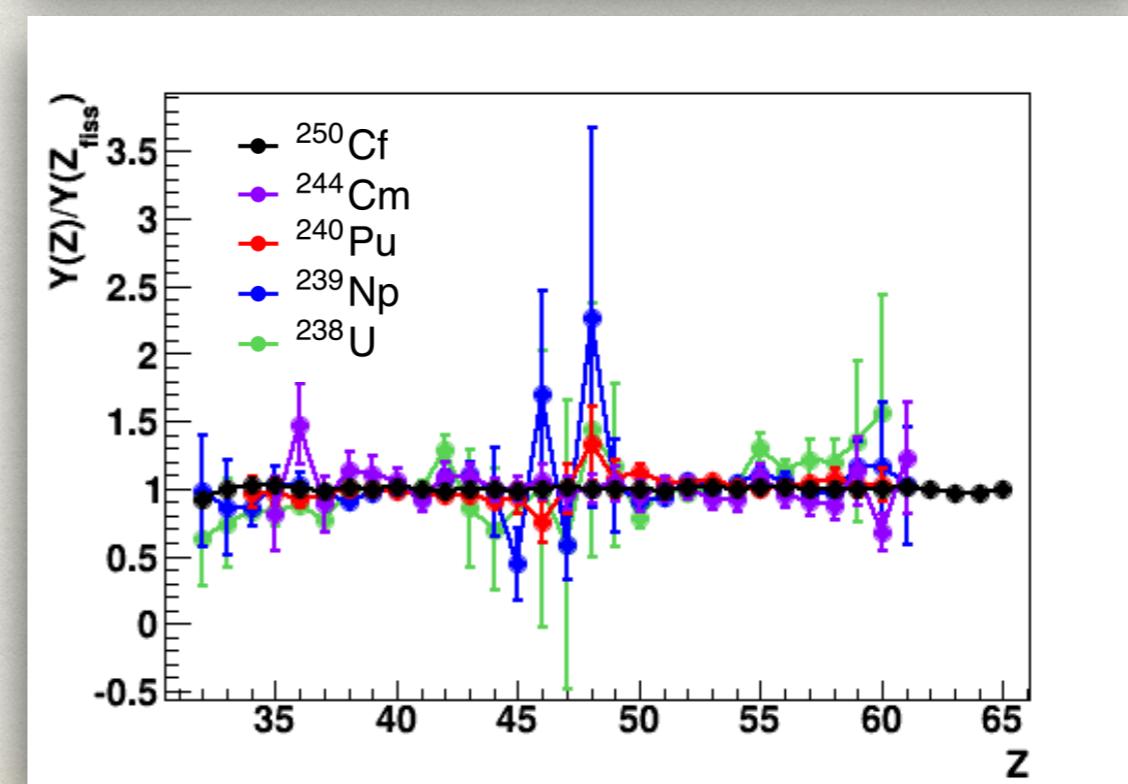
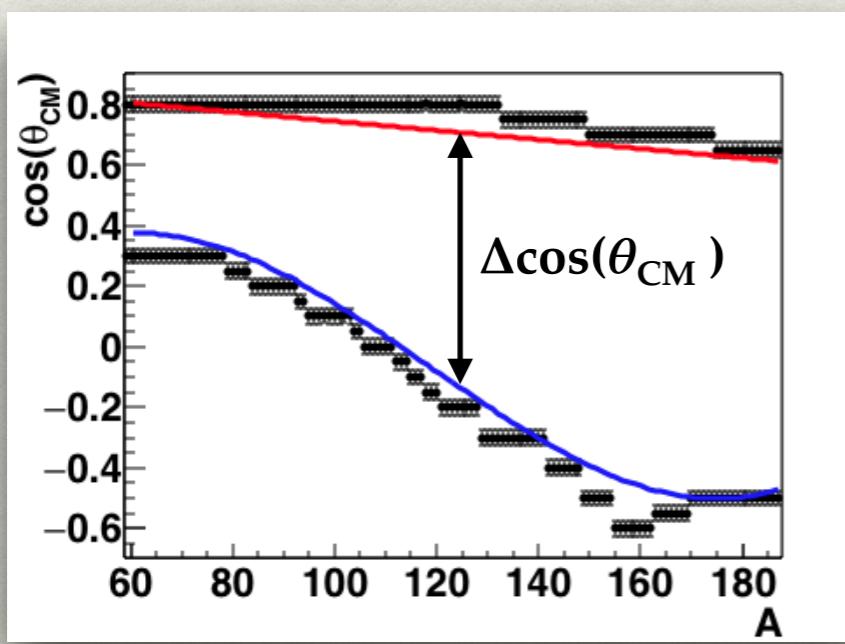
$$\langle \delta \rangle = 3.99$$

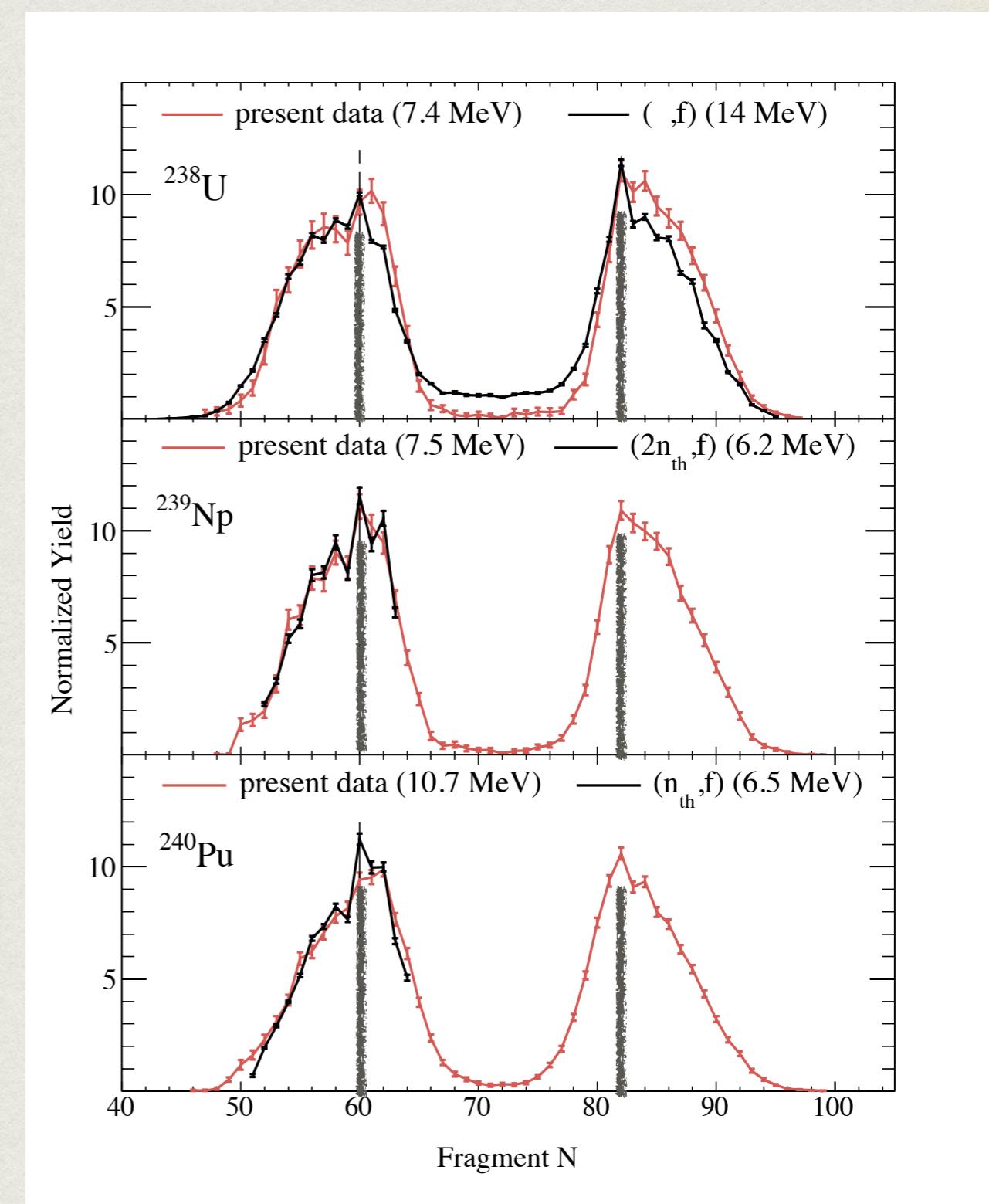
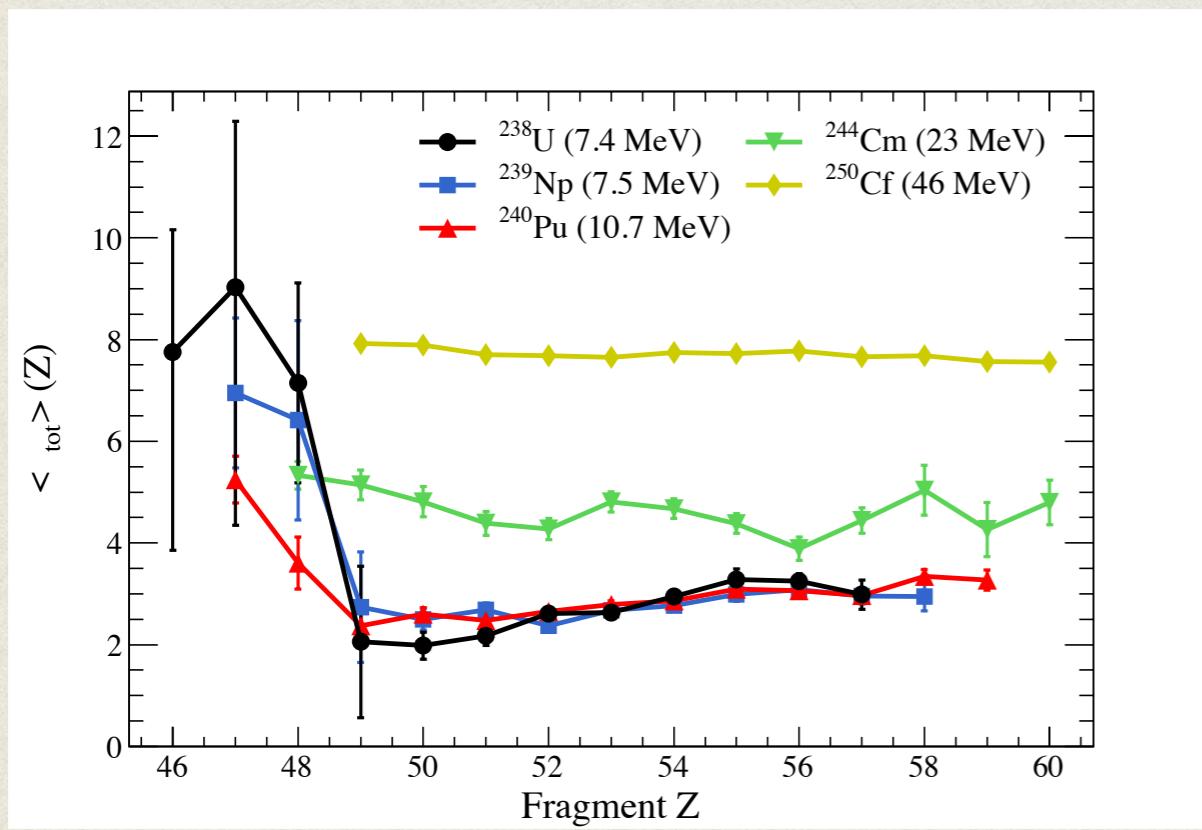


$$W(\theta_{CM}) = W(90^\circ) \left[1 + \alpha \cos^2(\theta_{CM}) \right]$$

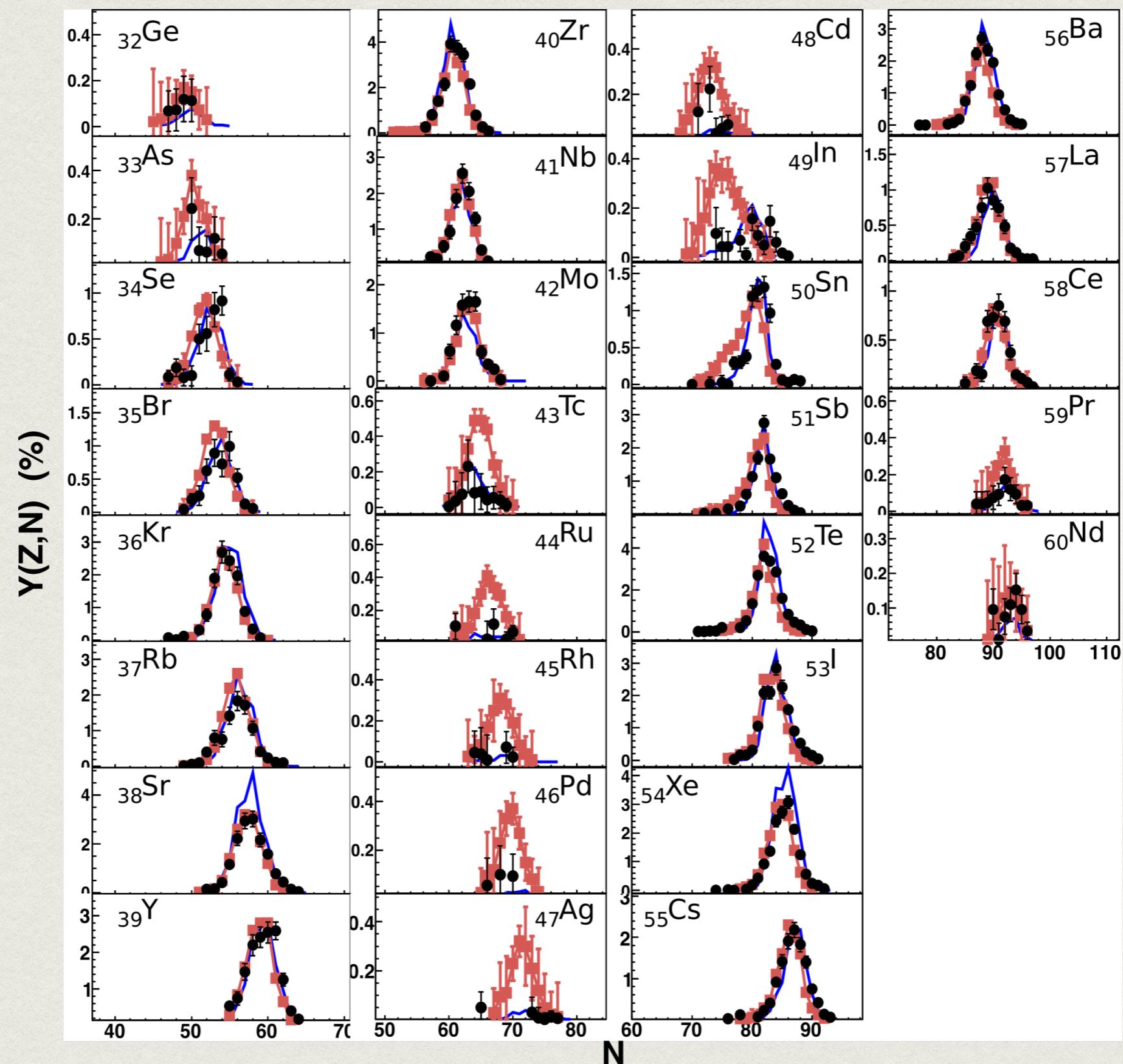


B. B. Back et al. PRC 32 (1985) 195

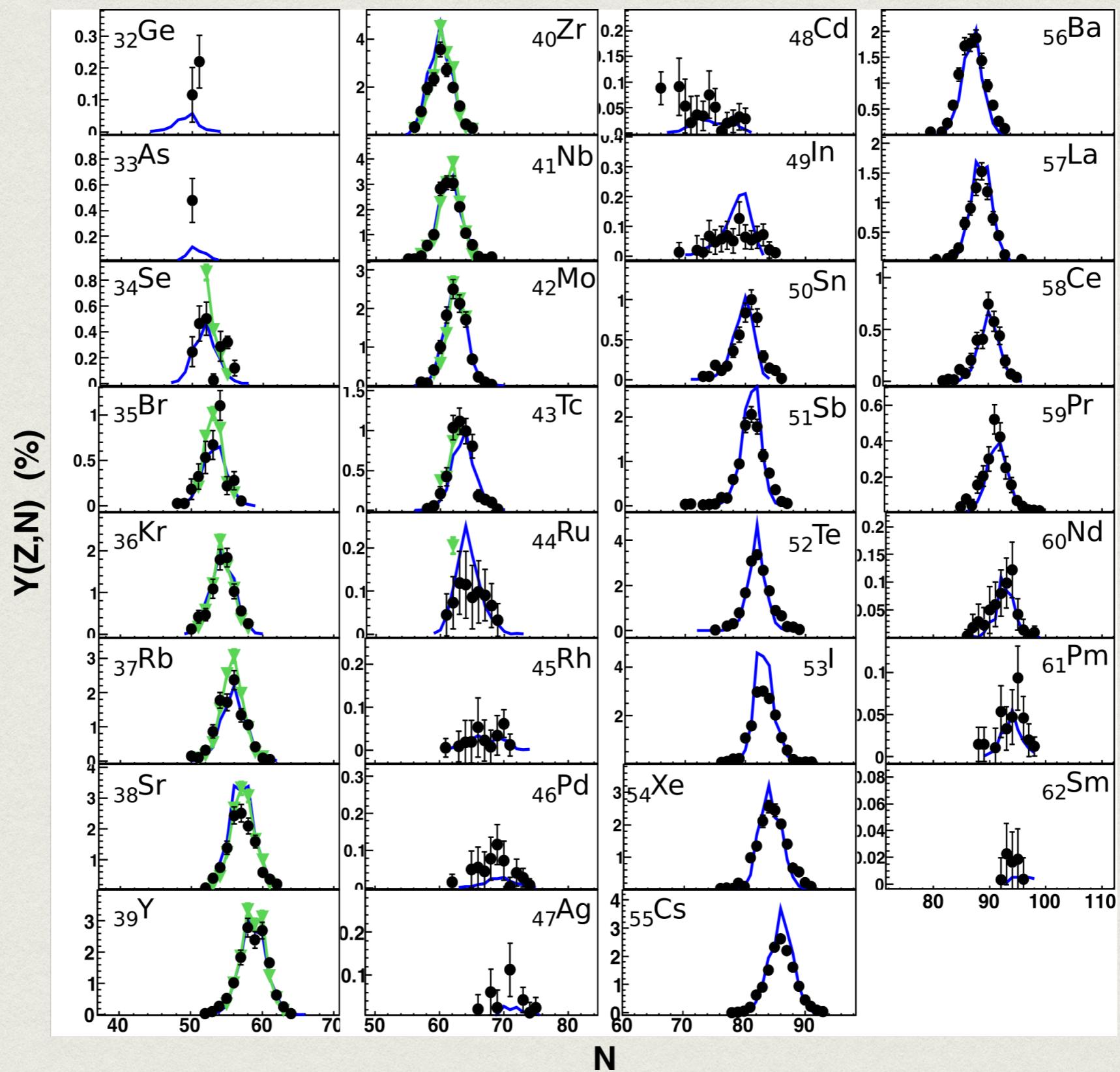




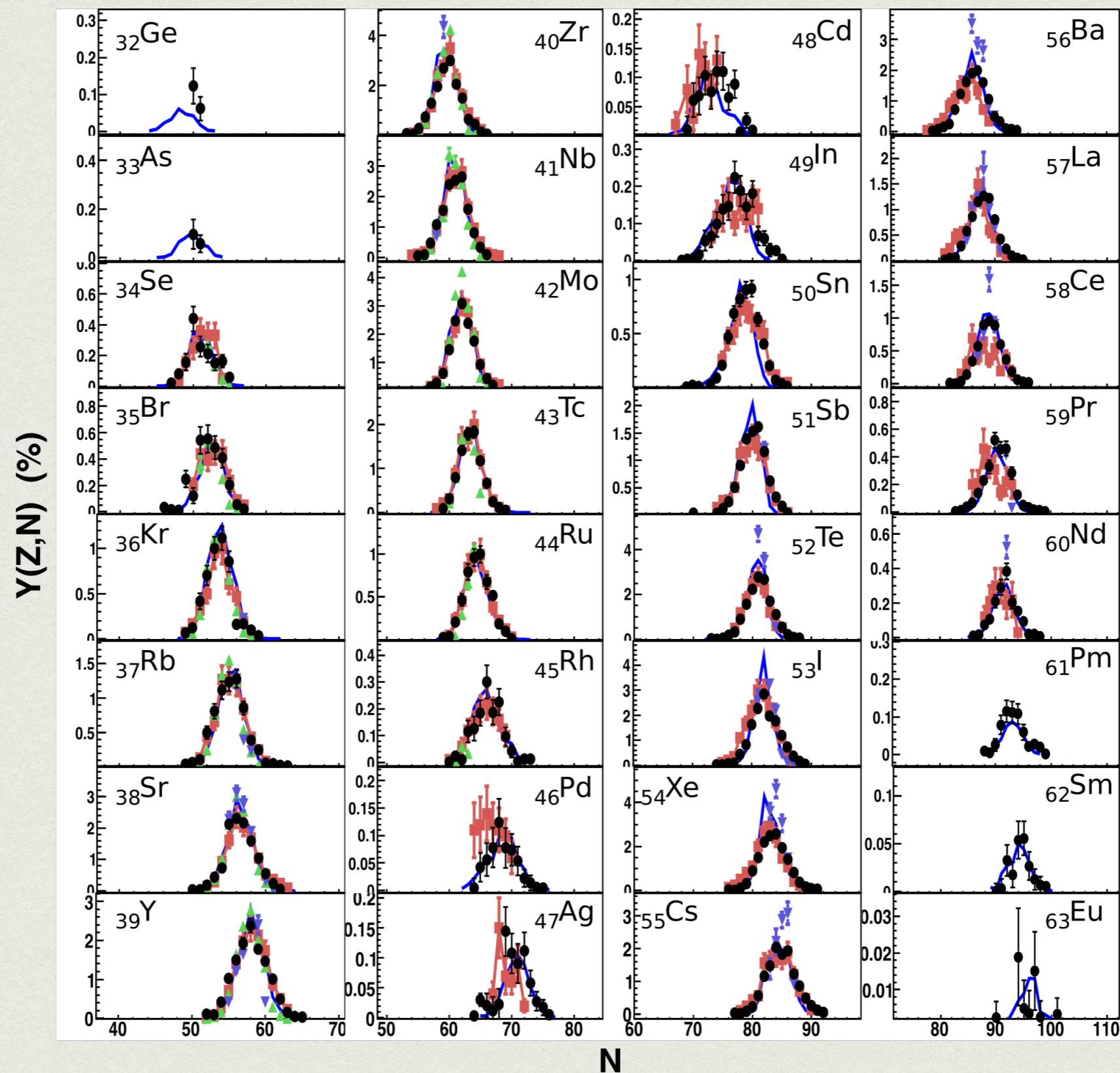
^{238}U



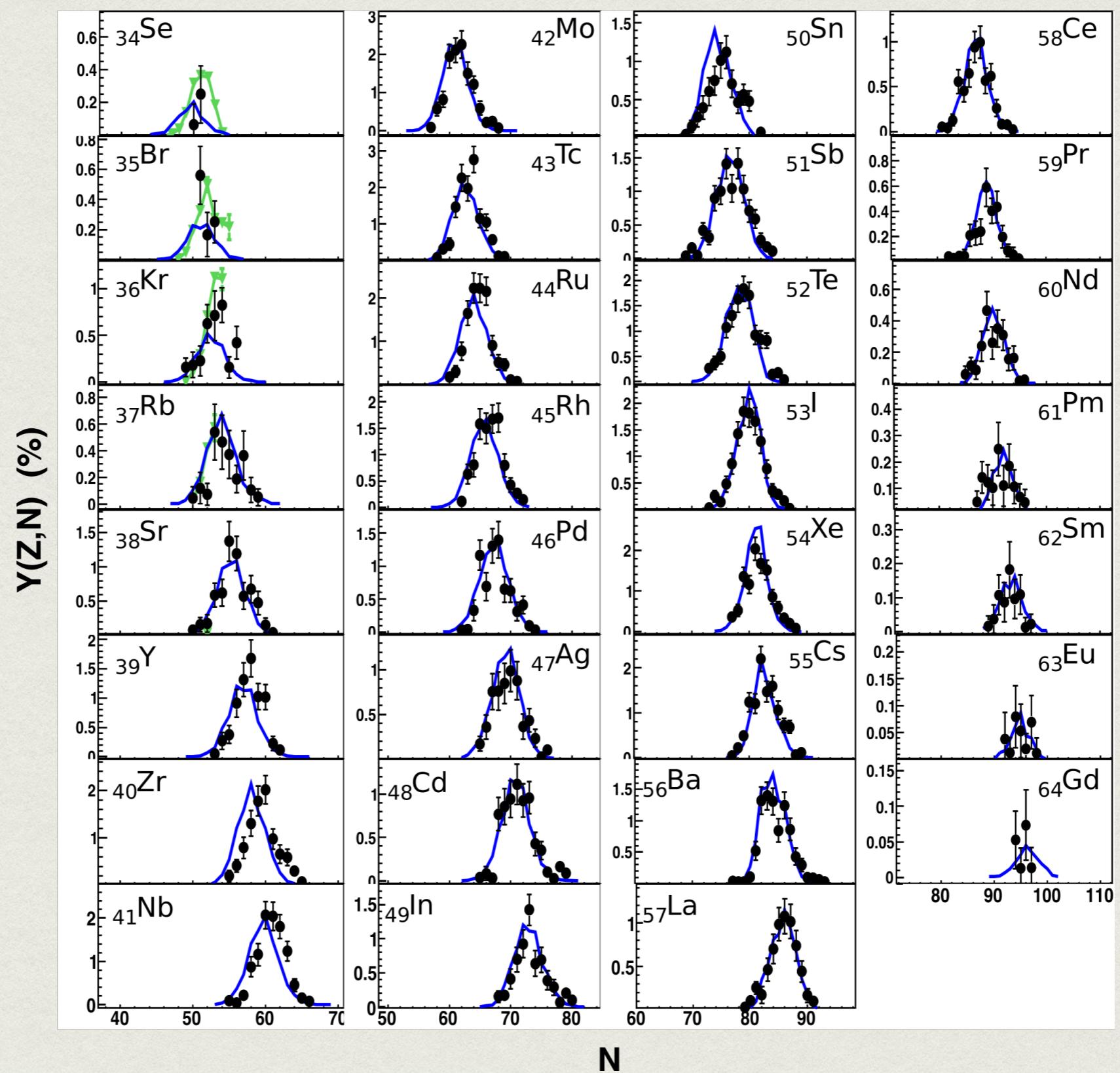
^{239}Np



^{240}Pu



^{244}Cm



^{250}Cf

