

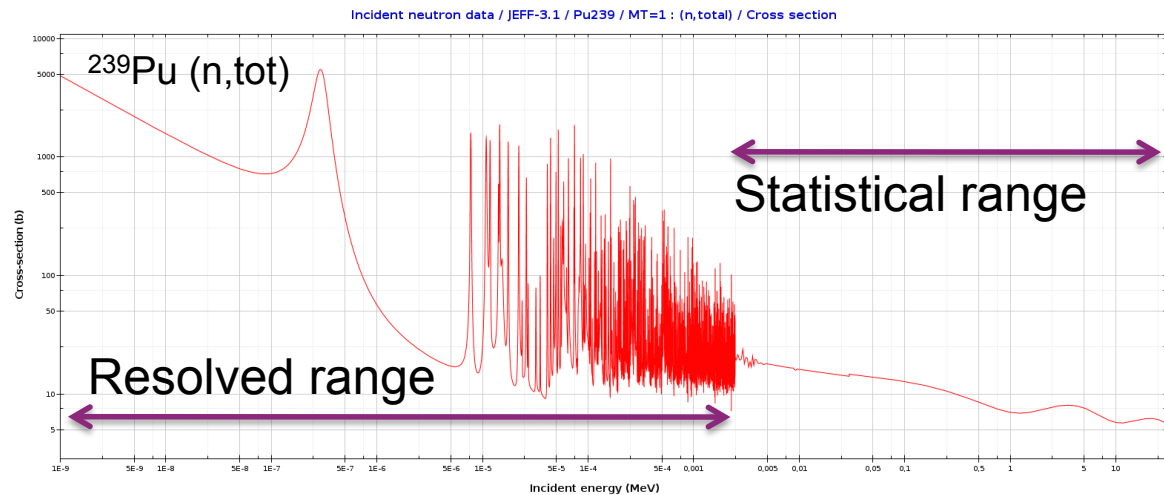
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TOWARDS IMPROVED METHODS FOR FISSION CROSS SECTION EVALUATION IN STATISTICAL ENERGY RANGE



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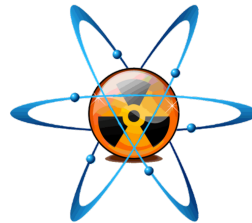
State of the art of evaluated fission cross sections and evaluation models



Undergoing developments in the **Conrad*** code and numerical validation



Physics underneath



Conclusion and outlook

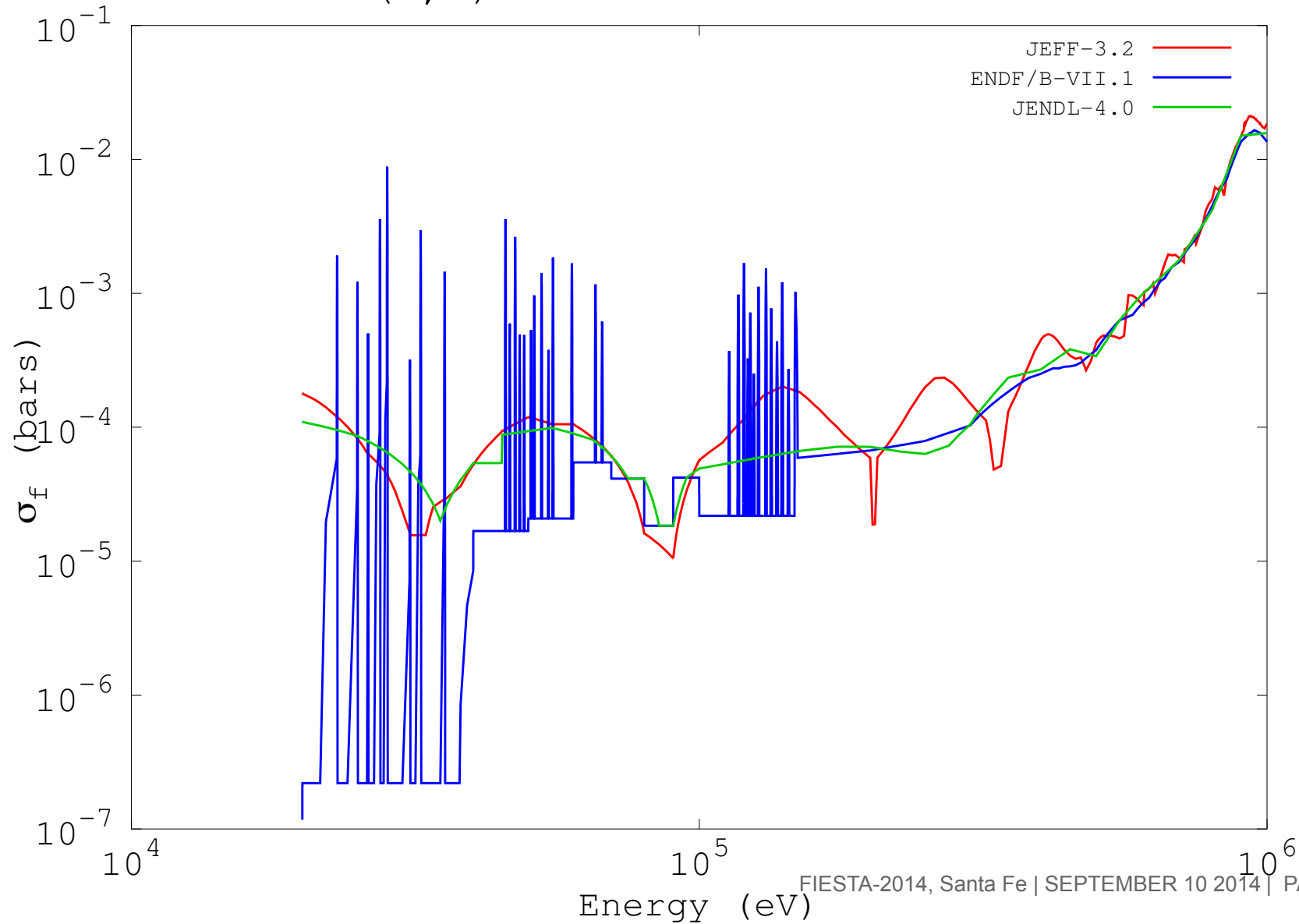




STATE OF THE ART OF EVALUATED FISSION CROSS SECTIONS AND EVALUATION MODELS

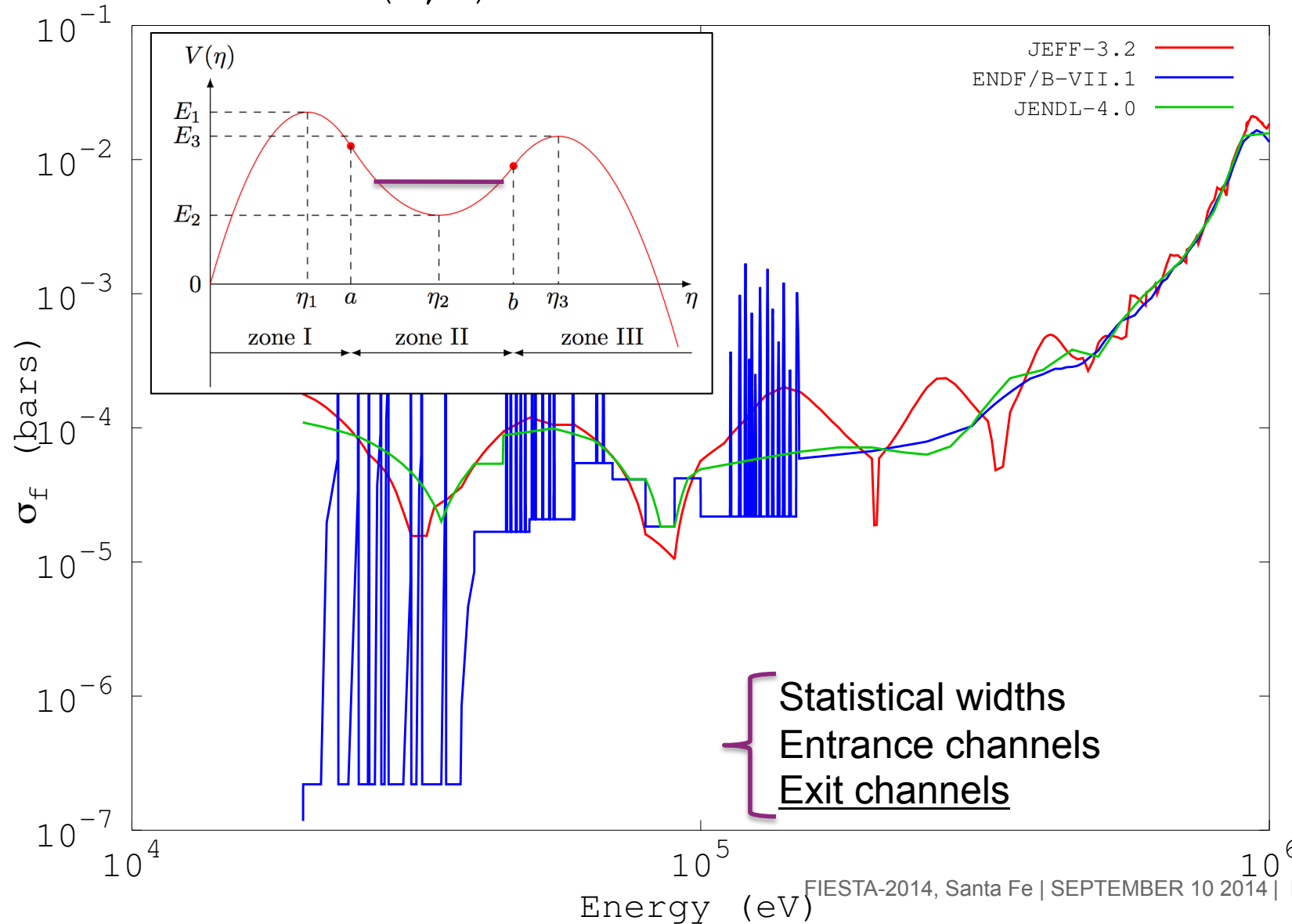


^{238}U (n, f) cross section evaluations





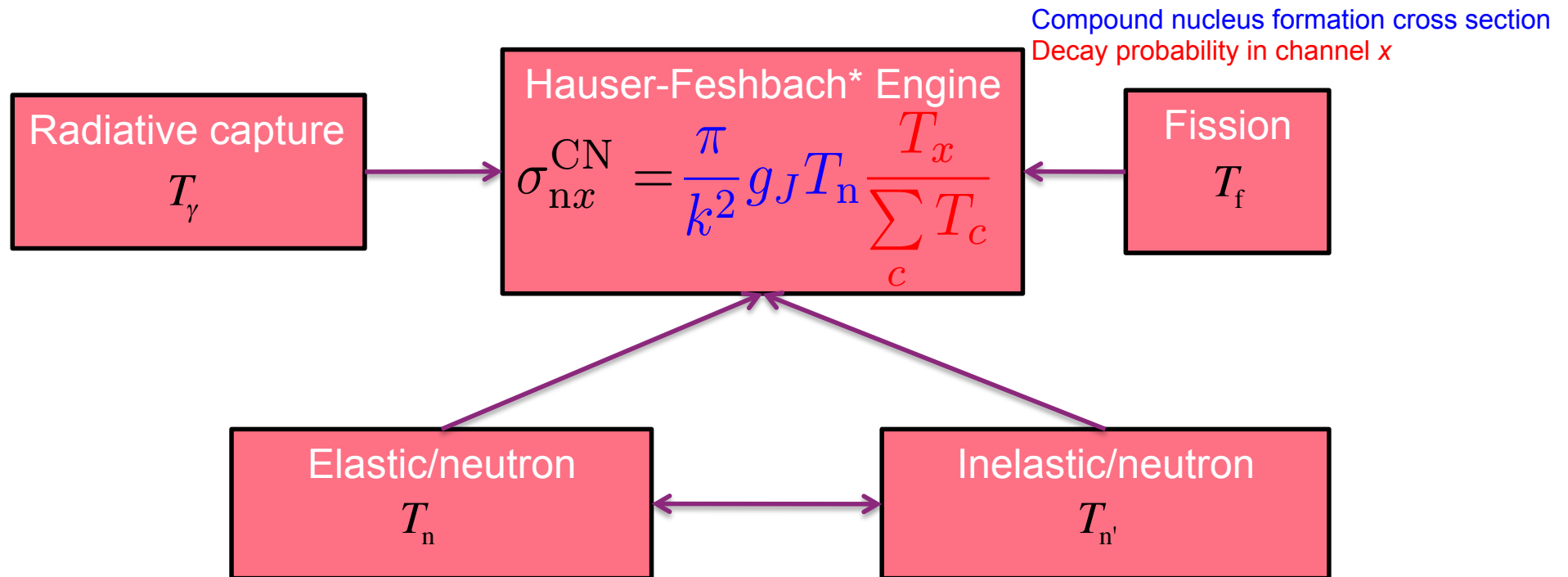
^{238}U (n, f) cross section evaluations





$$\sigma_{nx} = \sigma_{nx}^{\text{dir}} + \sigma_{nx}^{\text{CN}}$$

Each type of reaction is treated by a sub-model



The formalism intrinsically correlates the calculated cross sections

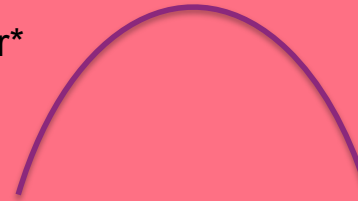


Fission T_f

Fission process modeled as particle passing through parabolic potential barrier

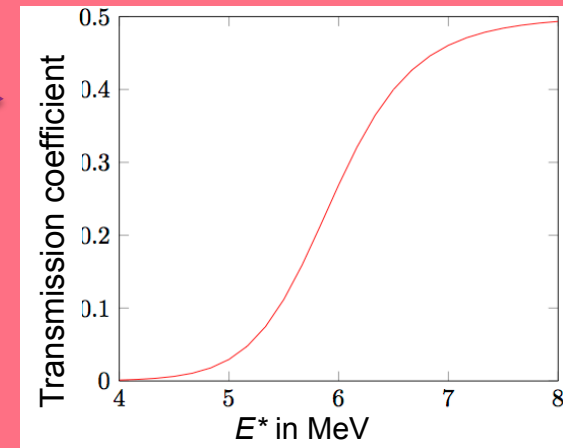
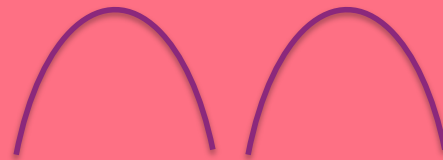
$$T_f = \frac{1}{1 + \exp\left(-2\pi \frac{E^* - V}{\hbar\omega}\right)}$$

Hill-Wheeler*
analytical
solution



Or passing through two barriers with statistical equilibrium

$$T_{\text{eff}} = \frac{T_A T_B}{T_A + T_B}$$



Possibility of local enhancements

$$T_{\text{eff}} = \frac{T_A T_B}{T_A + T_B} F_{AB}(E)$$

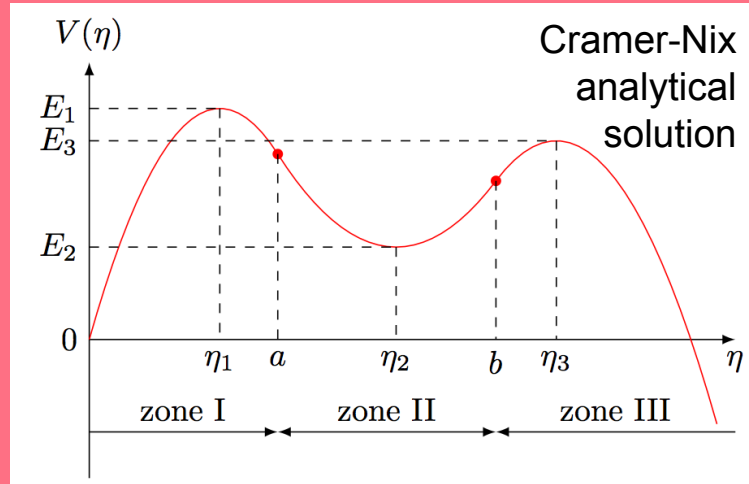
Intermediate state

$$F_{AB}(E) = 1 + \sum_{\text{class II}} \left[\frac{4}{T_A + T_B} + \left(\frac{E - E_{II}}{\Gamma_{II}/2} \right)^2 \left(1 - \frac{4}{T_A + T_B} \right) - 1 \right] \delta_{E \in [E_{II} \pm \Gamma_{II}/2]}$$

*D. L. Hill and J. A. Wheeler, Phys. Rev. 89, 1102 (1953)



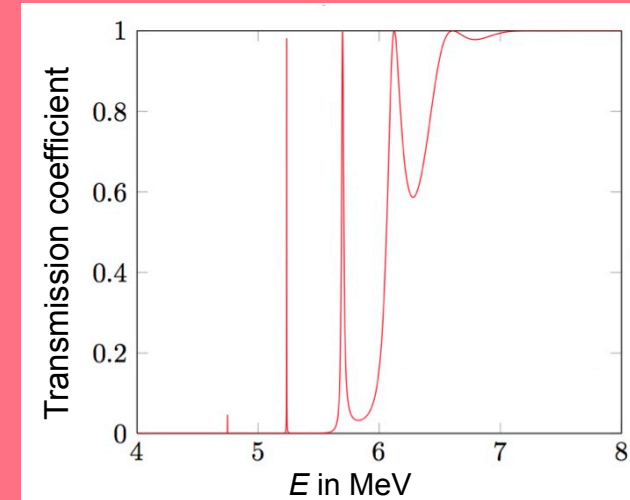
Fission T_f



Cramer-Nix potential*

$$V(\eta) = \begin{cases} E_1 - \frac{1}{2}\mu\omega_1^2(\eta - \eta_1)^2, & \eta \leq a \\ E_2 + \frac{1}{2}\mu\omega_2^2(\eta - \eta_2)^2, & a \leq \eta \leq b \\ E_3 - \frac{1}{2}\mu\omega_3^2(\eta - \eta_3)^2, & \eta \geq b \end{cases}$$

$$T_f = \sqrt{\frac{\omega_3}{\omega_1}} \left| \frac{T}{A} \right|^2$$



$$\frac{T}{A} = \frac{v'u'2i\sqrt{2/\pi}}{\det M} \quad \begin{aligned} u &= \sqrt{2\mu\omega_1/\hbar}(\eta - \eta_1) \\ v &= \sqrt{2\mu\omega_2/\hbar}(\eta - \eta_2) \\ w &= \sqrt{2\mu\omega_3/\hbar}(\eta - \eta_3) \end{aligned}$$

$$\alpha_1 = (E_1 - E)/\hbar\omega_1, \quad \alpha_2 = (E_2 - E)/\hbar\omega_2, \quad \alpha_3 = (E_3 - E)/\hbar\omega_3$$

$$M = \begin{bmatrix} E_a(\alpha_1, -u) & -V_a(\alpha_2, v) & -U_a(\alpha_2, v) & 0 \\ -u'E_a^{(-u)}(\alpha_1, -u) & -v'V_a^{(v)}(\alpha_2, v) & -v'U_a^{(v)}(\alpha_2, v) & 0 \\ 0 & V_b(\alpha_2, v) & U_b(\alpha_2, v) & -E_b(\alpha_3, w) \\ 0 & v'V_b^{(v)}(\alpha_2, v) & v'U_b^{(v)}(\alpha_2, v) & -w'E_b^{(w)}(\alpha_3, w) \end{bmatrix}$$

* J. D. Cramer and J. R. Nix, Phys. Rev. C 2, 1048 (1970)



UNDERGOING DEVELOPMENTS IN THE **Conrad** CODE AND NUMERICAL VALIDATION



The CEA/Cadarache code **Conrad** was created for such evaluation

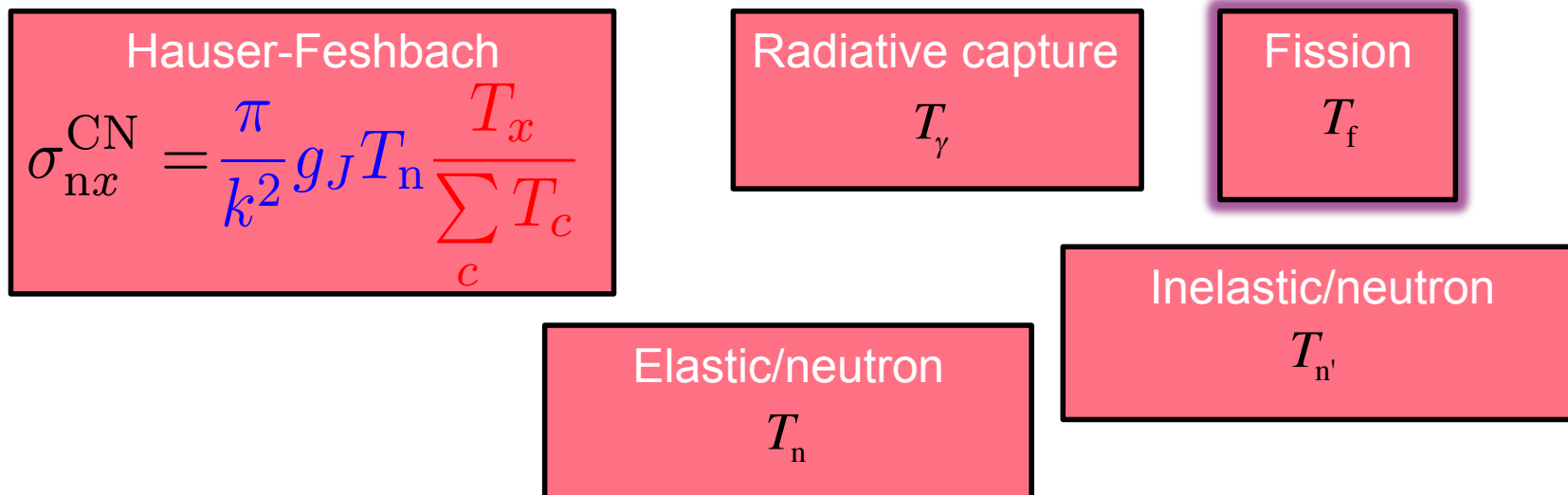
->originally developed for the resolved resonances range

- (Reich-Moore & multi-level Breit-Wigner)

->few statistical range capabilities

- Average R-matrix only

->New developments using the Talys* code as guideline





Neutron T_n $T_{n'}$ - Coupled channels calculation performed by the ECIS* code
- Average R-matrix

Gamma T_γ

$$T_\gamma^{J\Pi}(E^*) = 2\pi \sum_{X\ell} \sum_{I'=|J-\ell|}^{J+\ell} \int_0^{E^*} d\epsilon_\gamma \epsilon_\gamma^{2\ell+1} f_{X\ell}(\epsilon_\gamma, E^*) \rho(E^* - \epsilon_\gamma, I', \Pi(-1)^{\ell+\delta_{MX}})$$

Strength function for the
electromagnetic transition

Level density for the
nucleus after
emission of gamma

→ Several available models for both level densities and gamma strength functions thanks to library sharing with the FIFRELIN** code

$$\sigma_{nx} = \sigma_{nx}^{\text{dir}} + \sigma_{nx}^{\text{CN}}$$

For inelastic levels not taken into account in the coupled channels scheme:

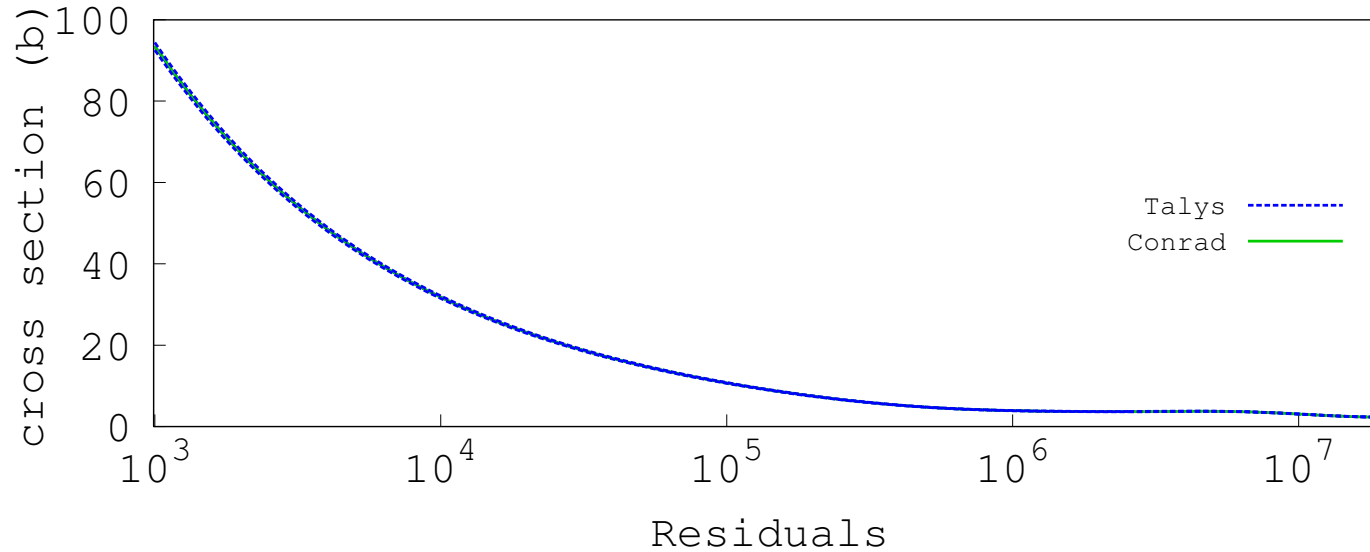
->DWBA calculation performed either by ECIS or internally by Conrad

*J. Raynal, Notes on ECIS94, CEA Saclay report No. CEA-N-2772 (1994) FIESTA-2014, Santa Fe | SEPTEMBER 10 2014 | PAGE 11 / 29

**O. Litaize, O. Serot, Phys. Rev. C 82, 054616 (2010)

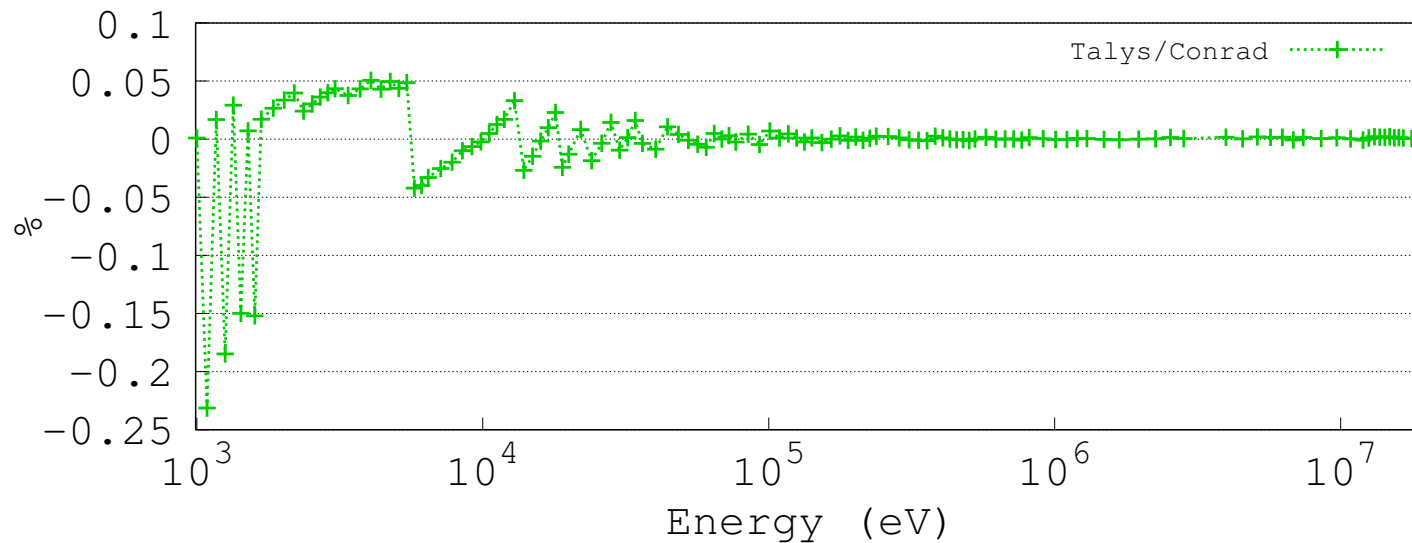


^{56}Fe (n,total) calculation



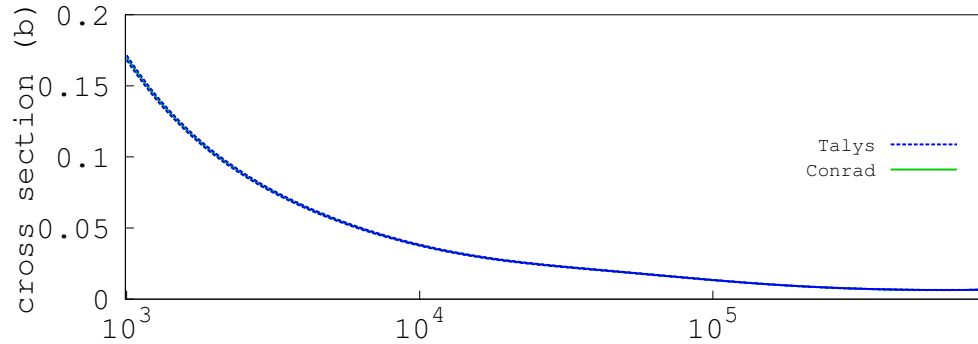
$$\sigma_{n,\text{tot}} \propto T_n$$

Validation of T_n

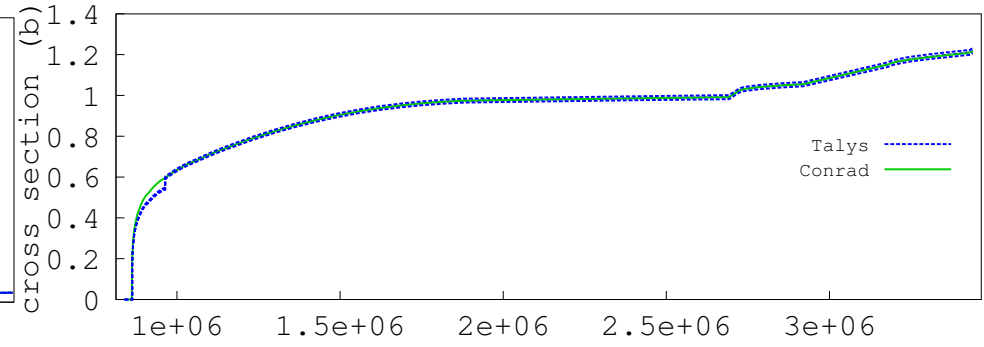




^{56}Fe (n, γ) calculation

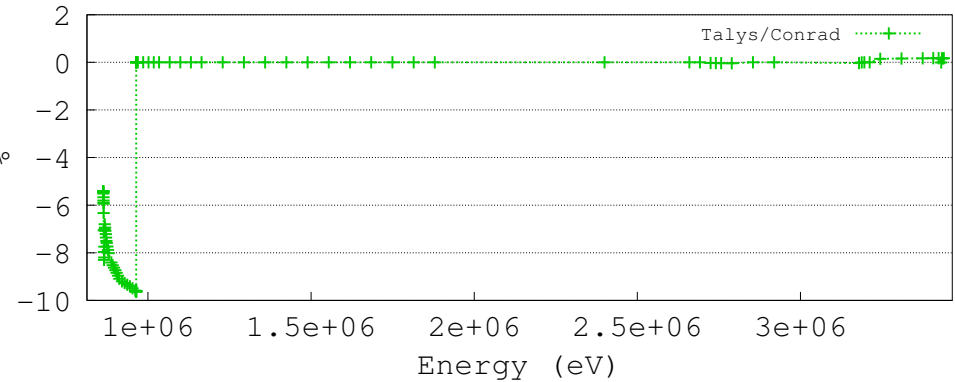
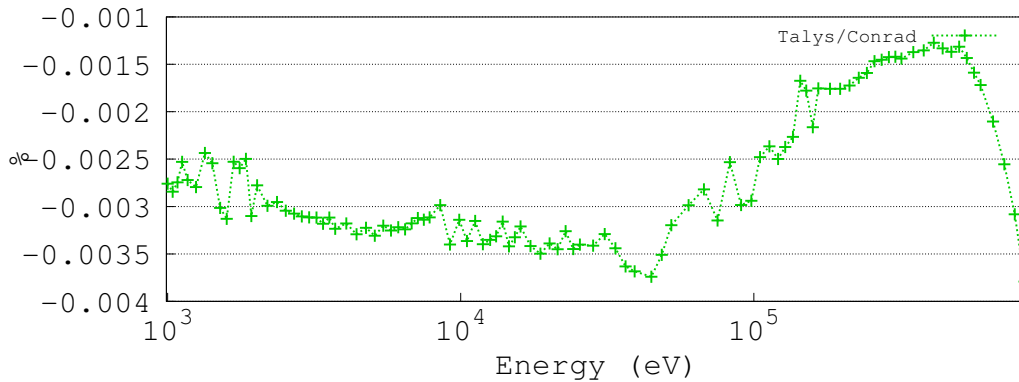


^{56}Fe (n,inel_{tot}) calculation



Residuals

Residuals



$$\sigma_{n\gamma} \propto \frac{T_n T_\gamma}{T_n + T_\gamma}$$

Validation of T_γ

$$\sigma_{nn'}^{\text{CN}} \propto \frac{T_n T_{n'}}{T_n + T_\gamma + T_{n'}}$$

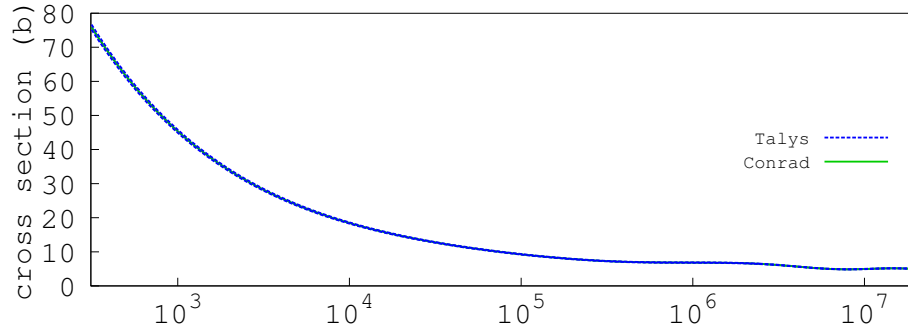
Validation of $T_{n'}$

$$\sigma_{nn'}^{\text{dir}} = \sigma_{nn'}^{\text{DWBA}}$$

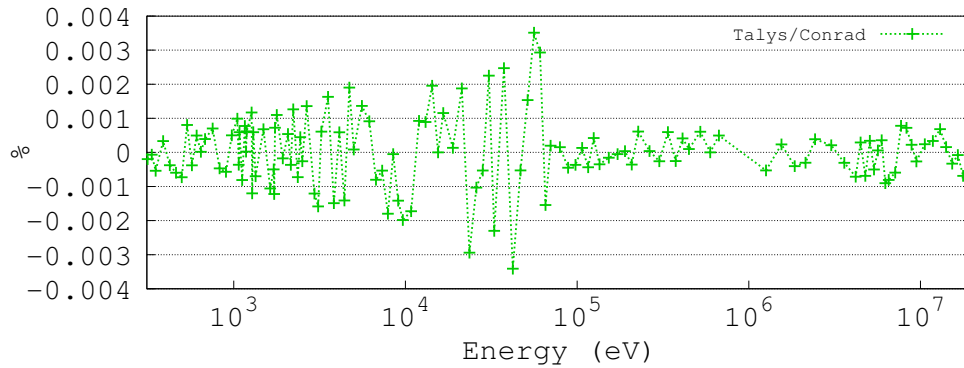
Validation of DWBA calculation



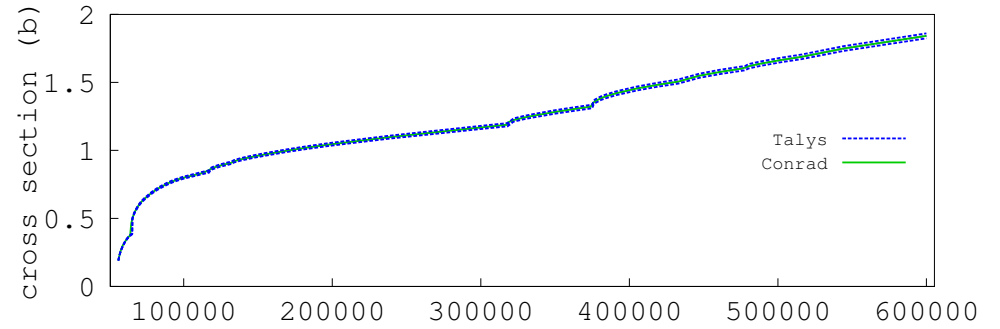
^{157}Gd (n,total) calculation



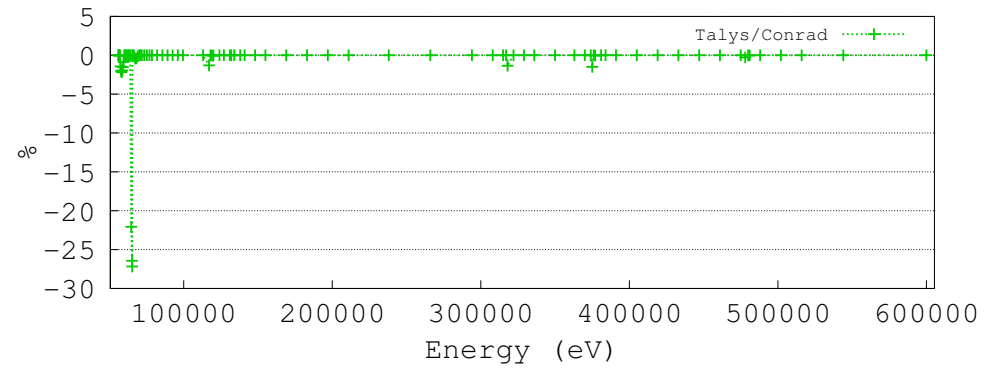
Residuals



^{157}Gd (n,inel_{tot}) calculation



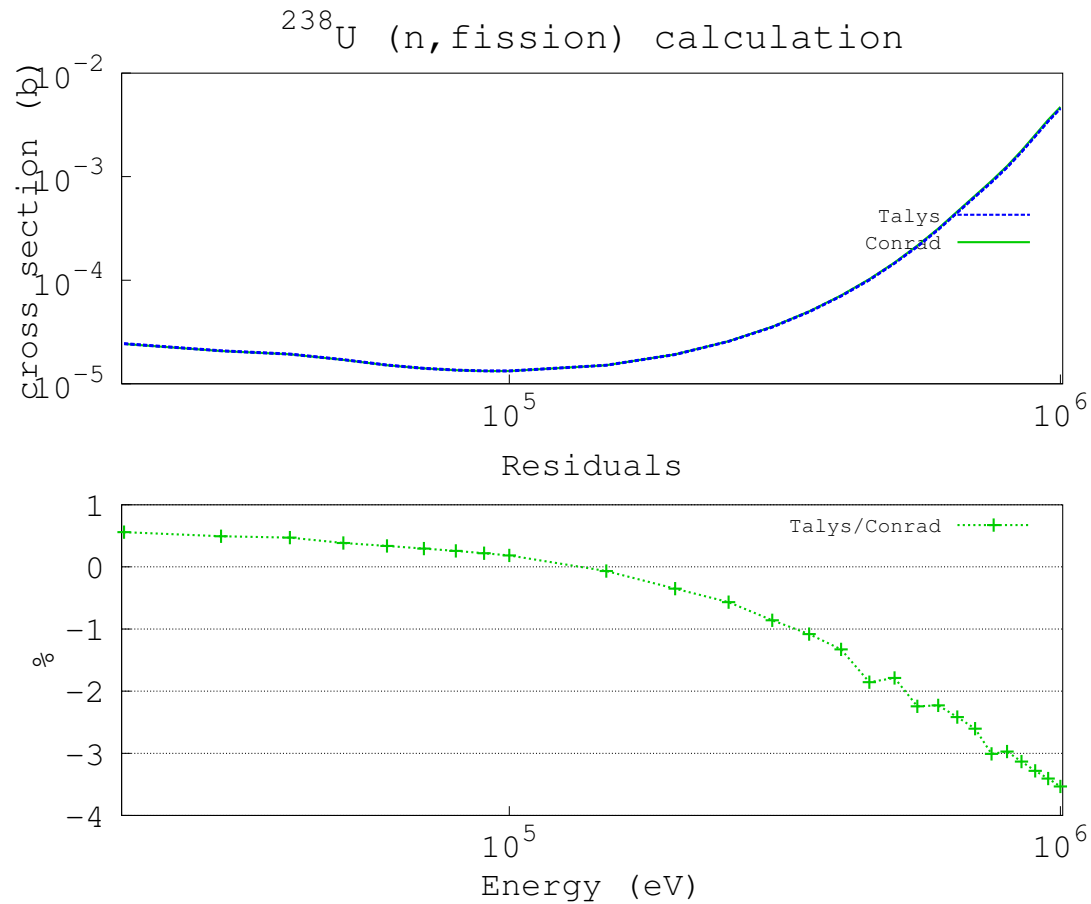
Residuals



Validation of the coupled channels calculated $T_{jl}^{J\Pi}$

$\sigma_{nn'}^{\text{dir}}$ is given by the coupled channels calculation

Other reaction are verified, what about fission?

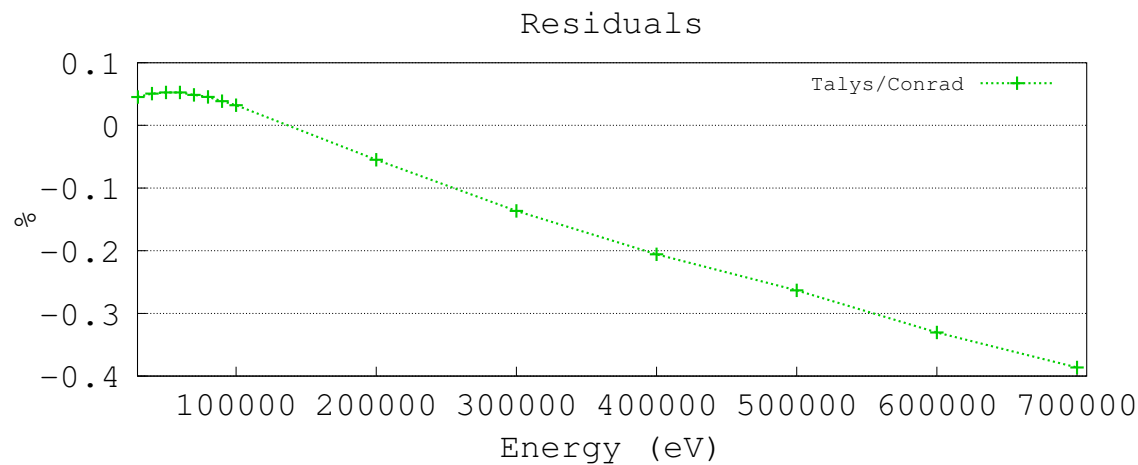
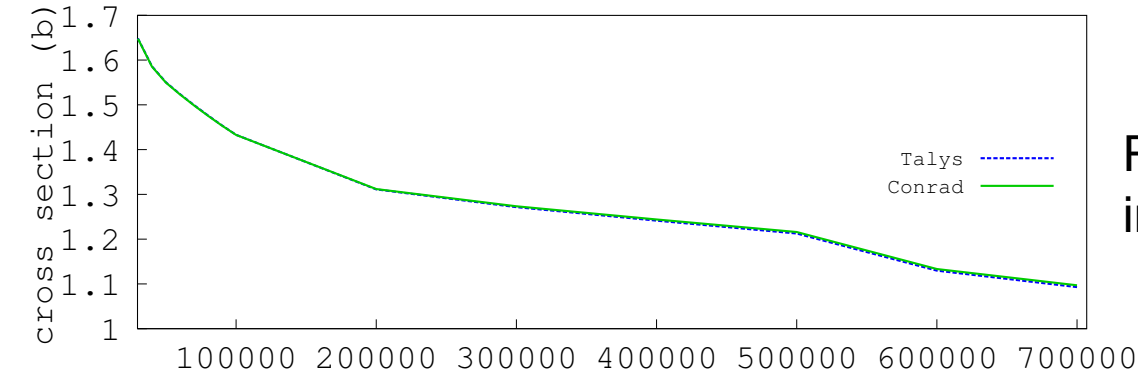


Remaining differences due to
integration method

$$T_{\text{barrier } i}(E^*) = \sum_{\text{disc. trans. states } t_i} T^{\text{Hill-Wheeler}}(E^* - \epsilon_{t_i}) + \int_{E_i^{\text{limit}}}^{E^*} \rho_i(J, \Pi, E^* - \epsilon) T^{\text{Hill-Wheeler}}(E^* - \epsilon) d\epsilon$$



^{239}Pu (n, fission) calculation



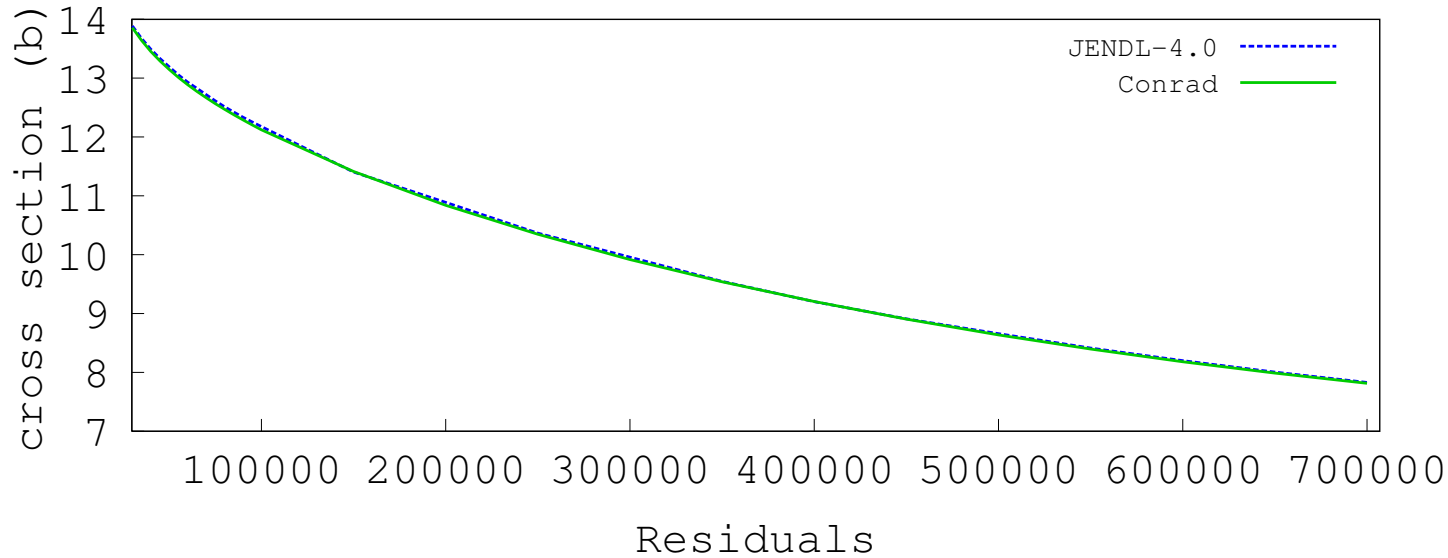
Remaining differences due to
integration method



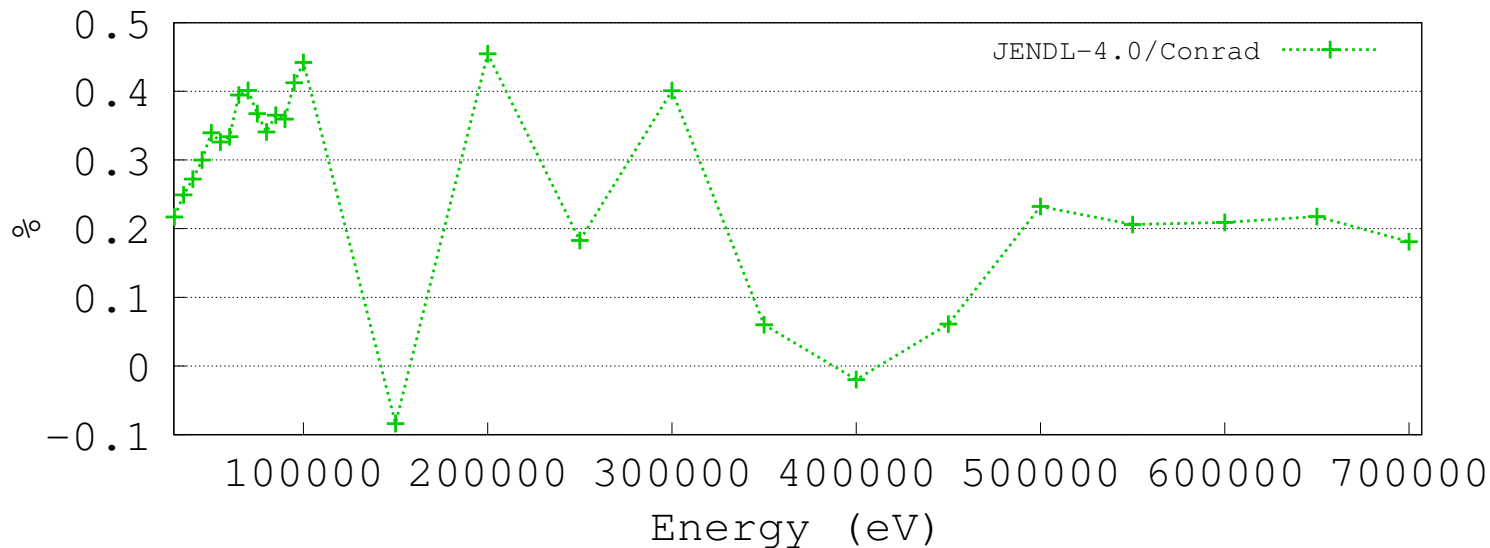
$$T_{\text{barrier } i}(E^*) = \sum_{\text{disc. trans. states } t_i} T^{\text{Hill-Wheeler}}(E^* - \epsilon_{t_i}) + \int_{E_i^{\text{limit}}}^{E^*} \rho_i(J, \Pi, E^* - \epsilon) T^{\text{Hill-Wheeler}}(E^* - \epsilon) d\epsilon$$



Comparison JENDL-4.0/Conrad for ^{239}Pu (n,total)

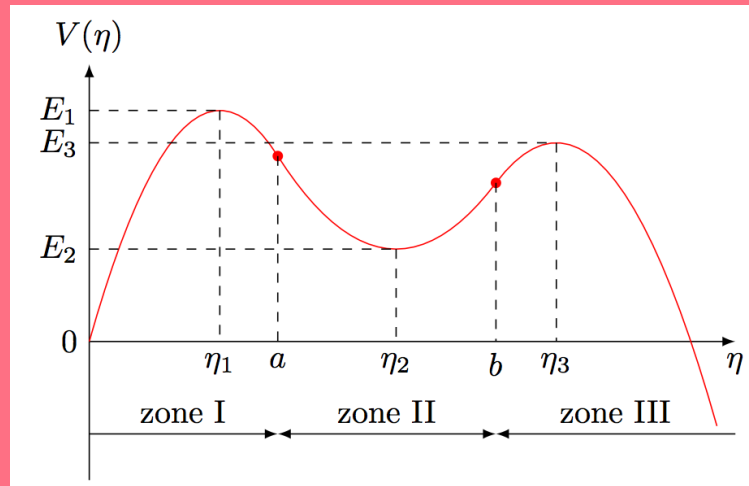


Numerical validation is achieved.

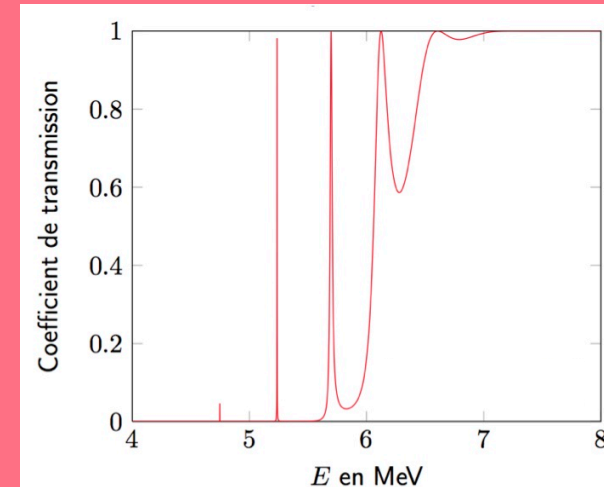


Comparison with evaluations is possible when parameters are known.

Fission T_f



$$V(\eta) = \begin{cases} E_1 - \frac{1}{2}\mu\omega_1^2(\eta - \eta_1)^2, & \eta \leq a \\ E_2 + \frac{1}{2}\mu\omega_2^2(\eta - \eta_2)^2, & a \leq \eta \leq b \\ E_3 - \frac{1}{2}\mu\omega_3^2(\eta - \eta_3)^2, & \eta \geq b \end{cases}$$



6 adjustable parameters per “global transition states”
(continuous potential)



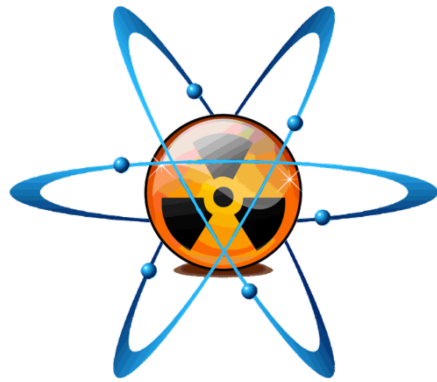
Too many for adjustments

6 adjustable parameters identical for all “global transition states”



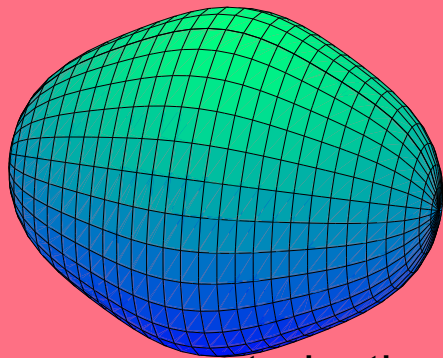
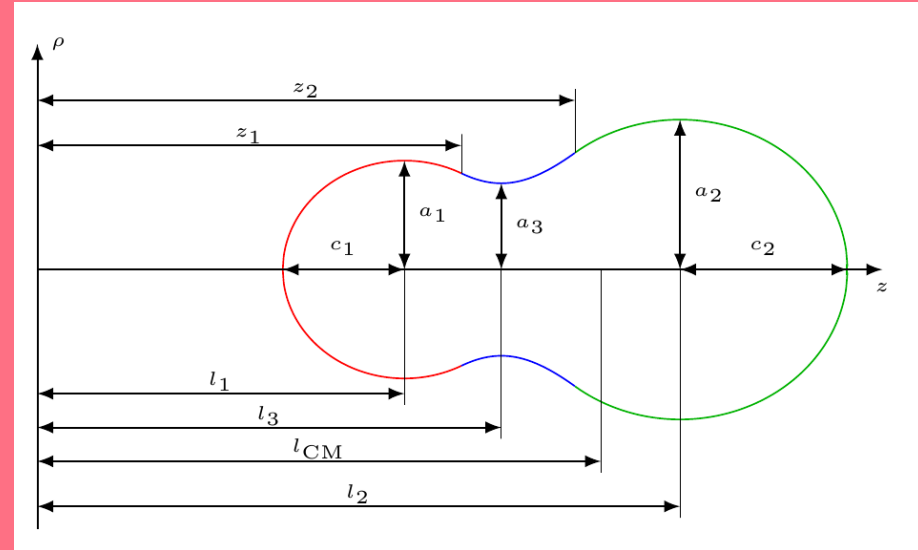
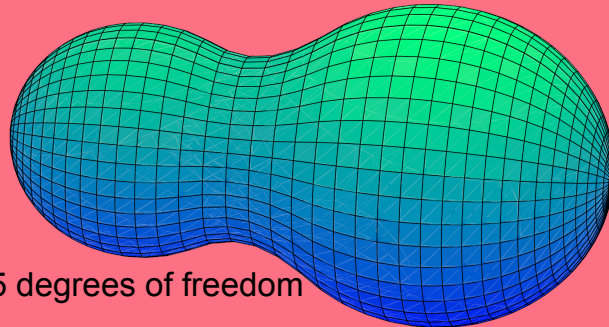
Section enhancements occur at similar energies and accumulate (too much)

Need for distinct – not one-by-one adjustable – parameters for each « global transition states »



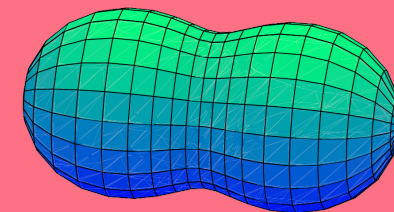
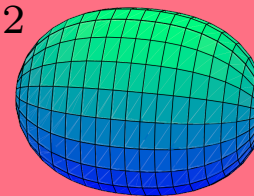
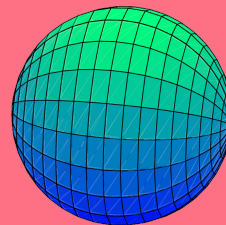
PHYSICS UNDERNEATH

Fission T_f Three quadratic surfaces



Legendre polynomials expansion

$$r(\theta) = R_0 \left[1 + \sum_{i=2}^{\infty} \beta_i P_i(\cos \theta) \right] / \lambda(\beta_2, \beta_3, \dots)$$

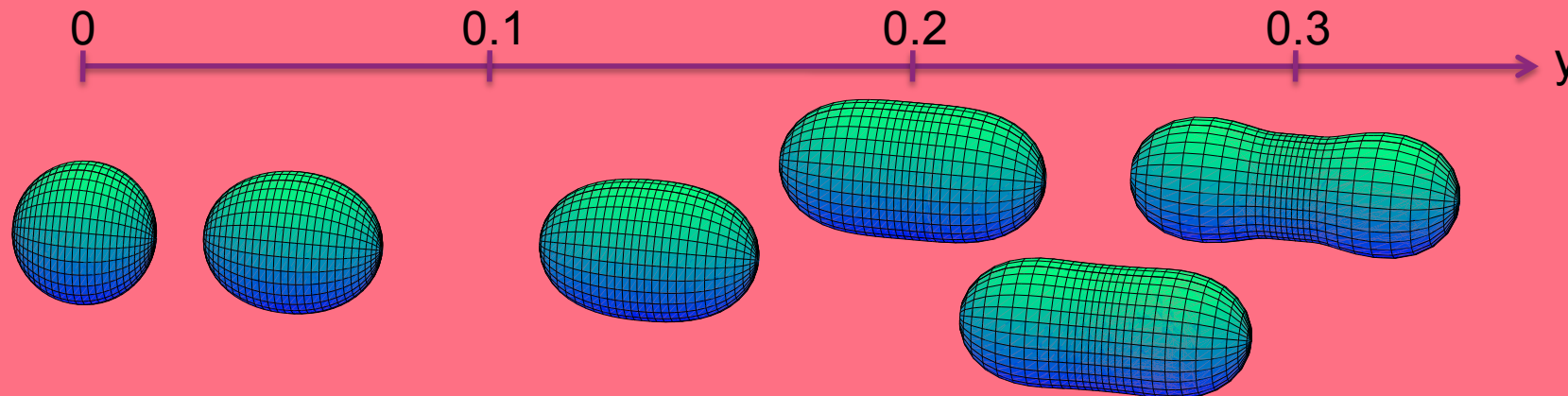


1) Provide $\rho(z)$ or $r(\theta)$

Fission T_f

$$r(\cos \theta) = a_0(y) \left[1 + \sum_{i=1}^4 a_{2i}(y) P_{2i}(\cos \theta) \right]$$

Numerical
expression

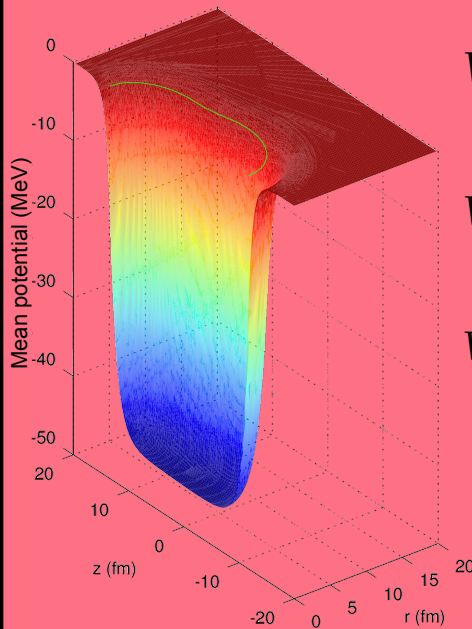


1 parameter = a “reasonable” shape description

Fission T_f $V(\text{shape}) = V_{\text{macro}}(\text{shape}) + \Delta V_{\text{micro}}(\text{shape})$

$$E_C(\text{shape}) = \frac{1}{2} \int_V d^3\vec{r}_1 \int_V d^3\vec{r}_2 \frac{\rho(\vec{r}_1)\rho(\vec{r}_2)}{\|\vec{r}_1 - \vec{r}_2\|} \quad E_S(\text{shape}) = \gamma \int_S d^2S$$

2) $V_{\text{macro}}(\text{shape})^* = E_C(\text{shape}) + E_S(\text{shape})$

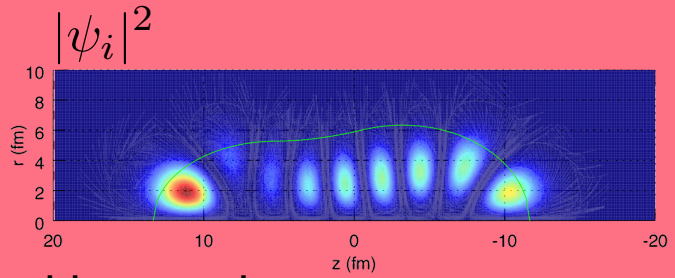


$$V_1(\vec{r}) = -\frac{V_0}{4\pi a_{\text{pot}}^3} \int_V \frac{e^{-\|\vec{r}-\vec{r}'\|/a_{\text{pot}}}}{\|\vec{r}-\vec{r}'\|/a_{\text{pot}}} d^3\vec{r}'$$

$$V_C(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{e^2 Z}{V} \int_V \frac{d^3\vec{r}'}{\|\vec{r}-\vec{r}'\|}$$

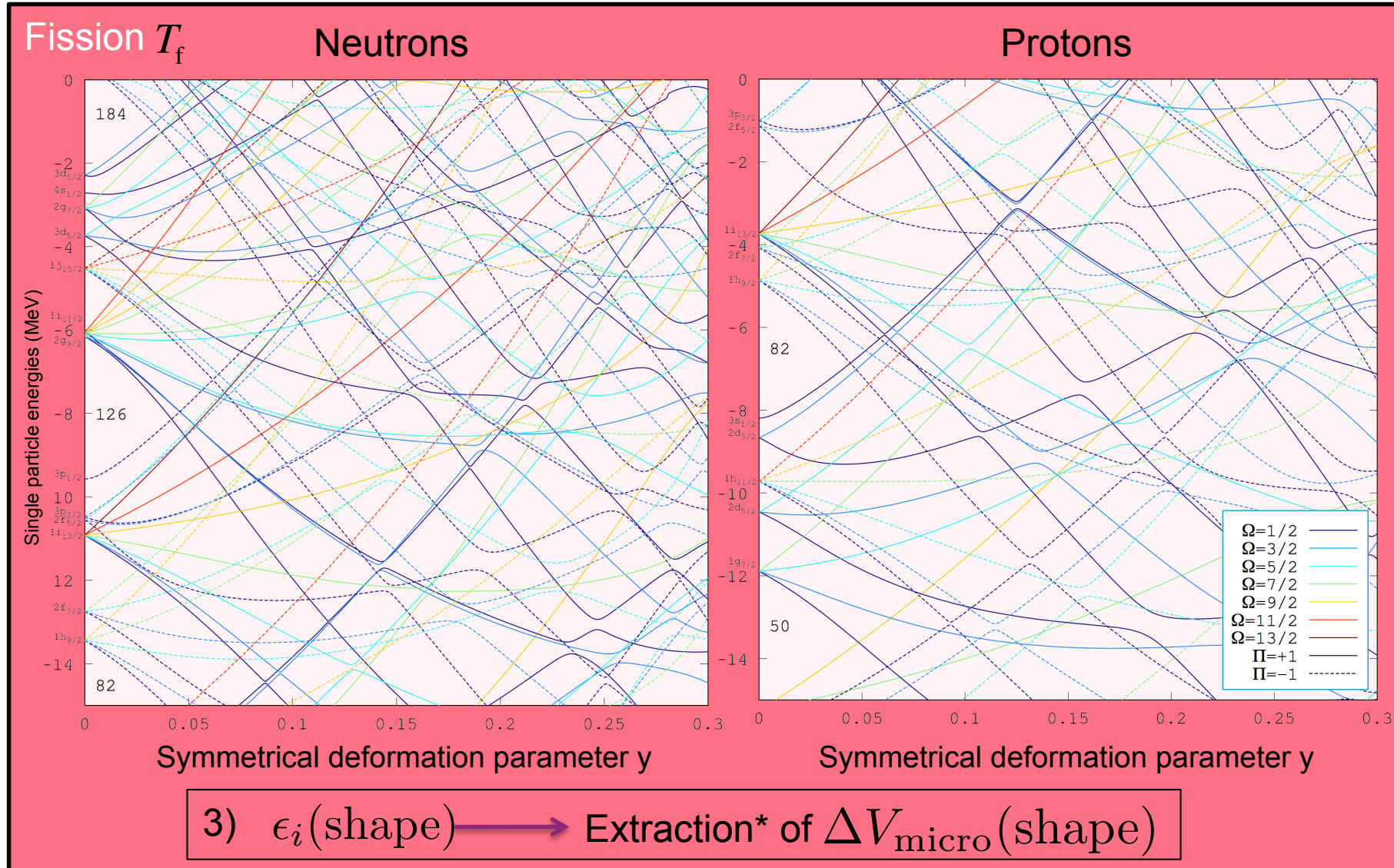
$$V_{\text{s.o.}}(\vec{r}) = -\lambda \left(\frac{\hbar^2}{2mc} \right)^2 \frac{\vec{\sigma} \cdot \vec{\nabla} V_1 \times \vec{p}}{\hbar}$$

Independent particle equation $[\hat{T} + \hat{V}_1 + \hat{V}_C + \hat{V}_{\text{s.o.}}] \psi_i = \epsilon_i \psi_i$



Resolution using deformed harmonic oscillator base functions $**|n_r, n_z, \Lambda, \Sigma\rangle$

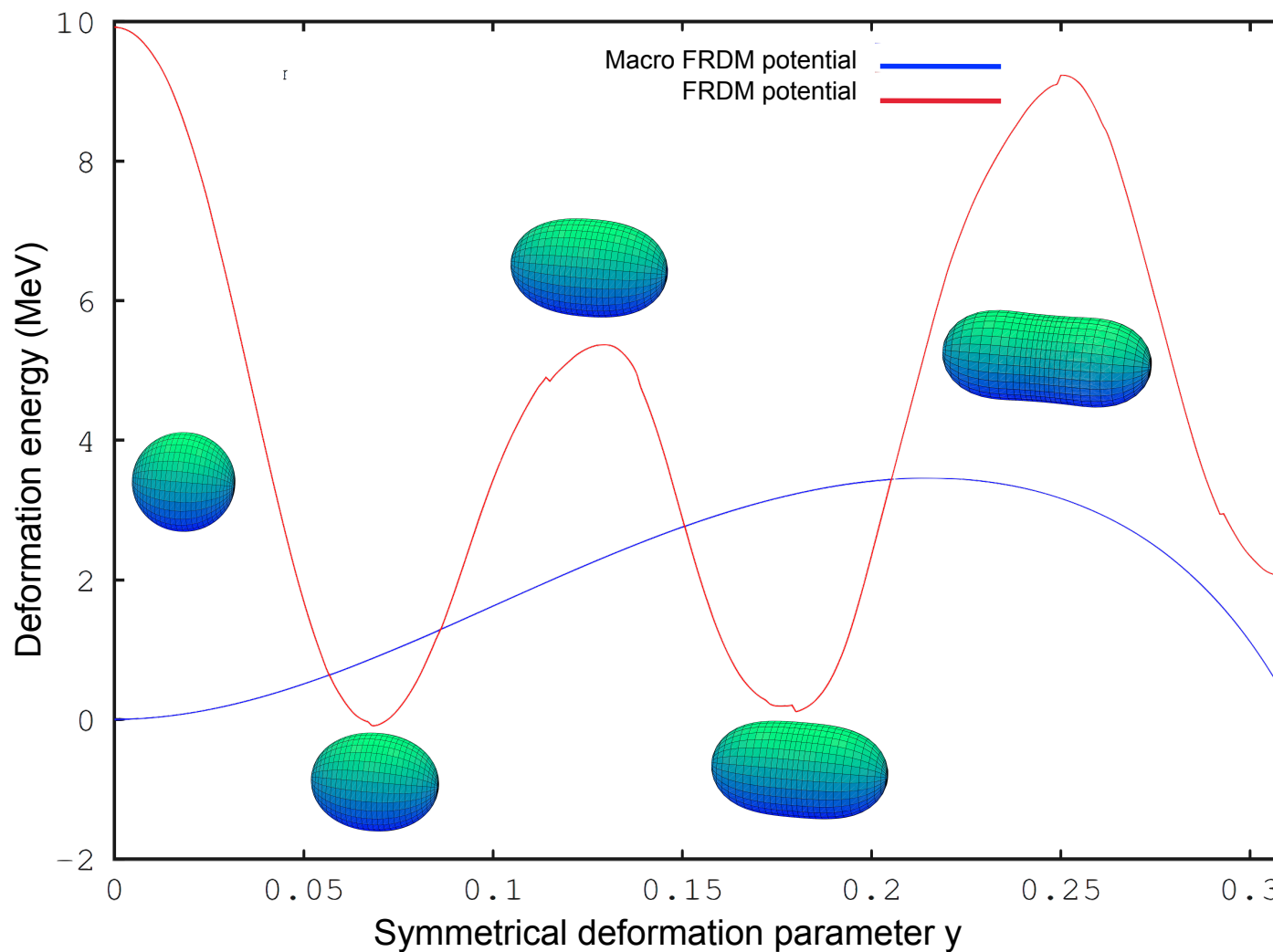
•P. Moller et al, At. Data Nucl. Data Tables 39 225 (1988)
 P. Moller et al, At. Data Nucl. Data Tables 39 213 (1988)
 ** J. Damgaard et al., Nucl. Phys. A 135, 432 (1969)



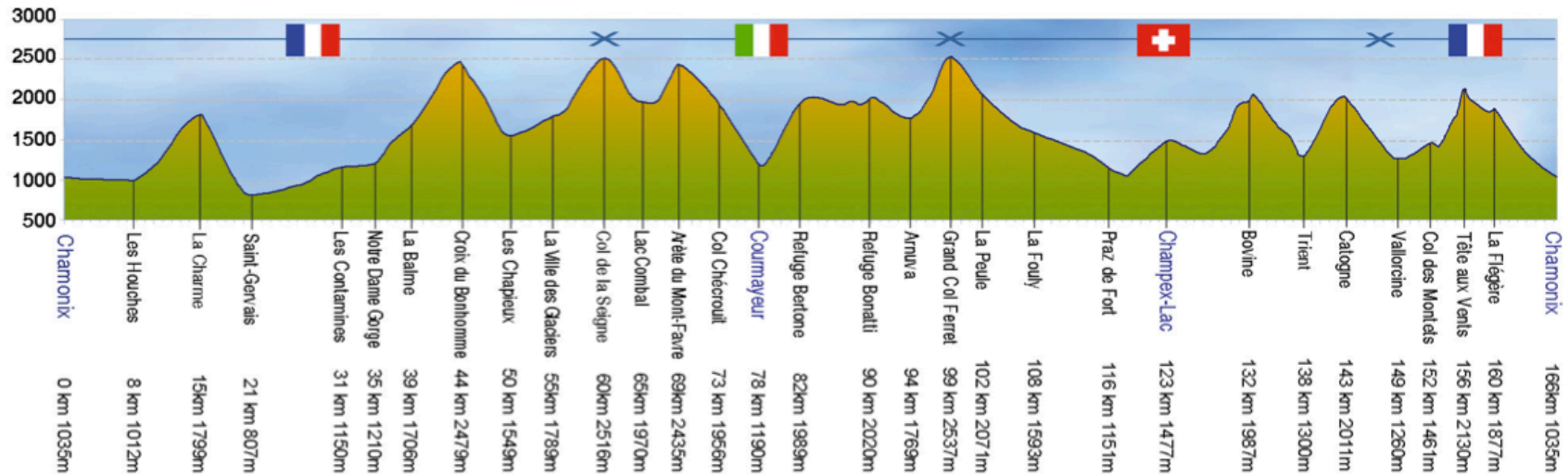
* M. Bolsterli et al., Phys. Rev. C 5, 1050 (1972)

Fission T_f

Reasonable values* for symmetric fission

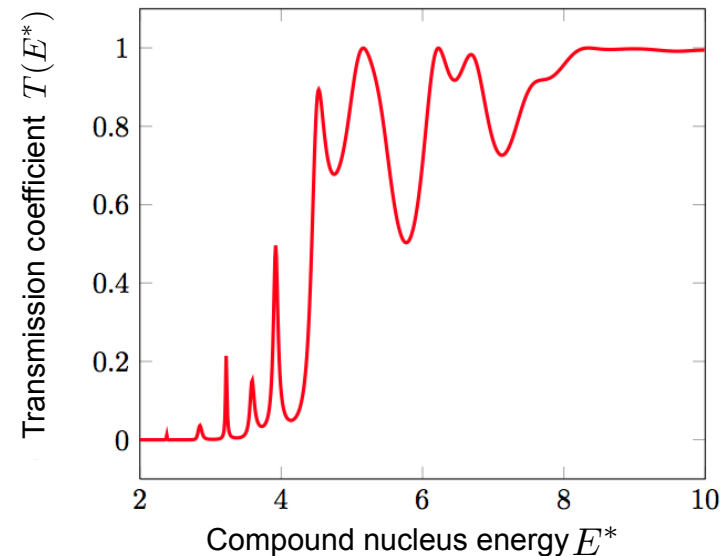


* M. Bolsterli et al., Phys. Rev. C 5, 1050 (1972)



Topologic profile of the Tour du Mont-Blanc (150km)

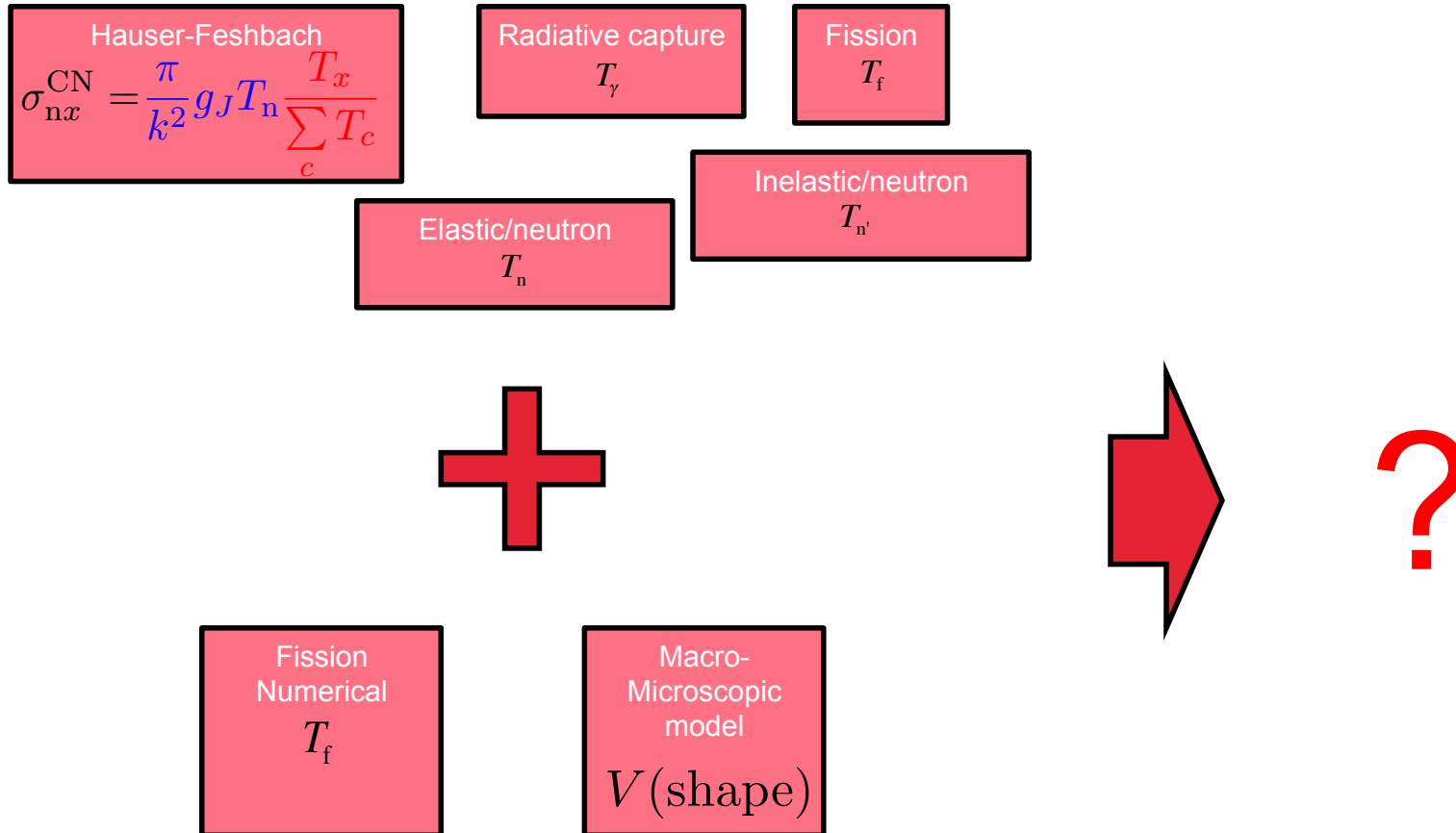
Implementation of the Numerov numerical method* for the calculation of fission transmission coefficients



*H. Durate, B. Morillon, P. Romain, CEA/(DAM/DEN/DSM) seminar (2013)



CONCLUSION AND OUTLOOK



Actual evaluated fission cross section calculated on some selected nuclei



Determination of a single dimension path for fission (mass asymmetry exploration)

Axially asymmetric shapes (gamma deformations)

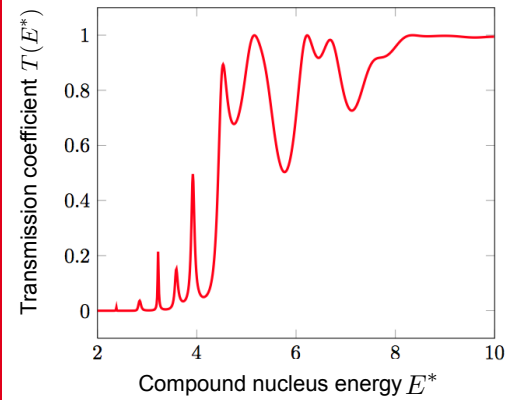
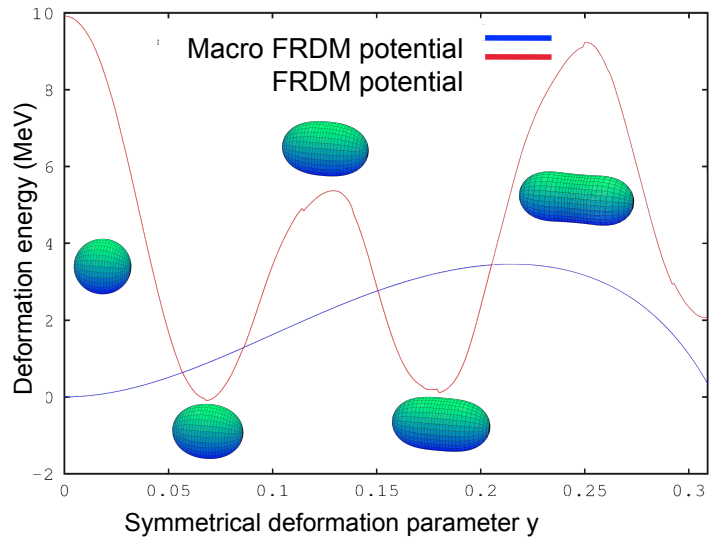
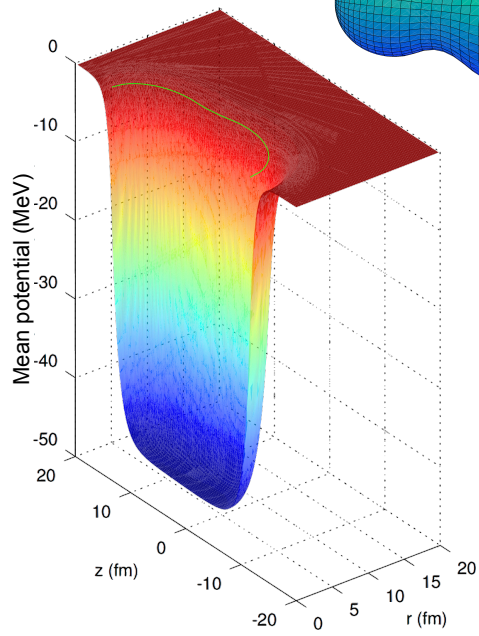
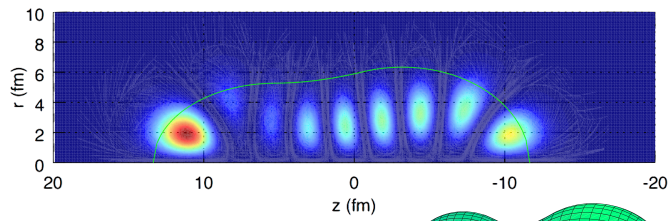
Dependence of the inertial parameter on deformation

Degree of freedom in the width fluctuation factor for fission*

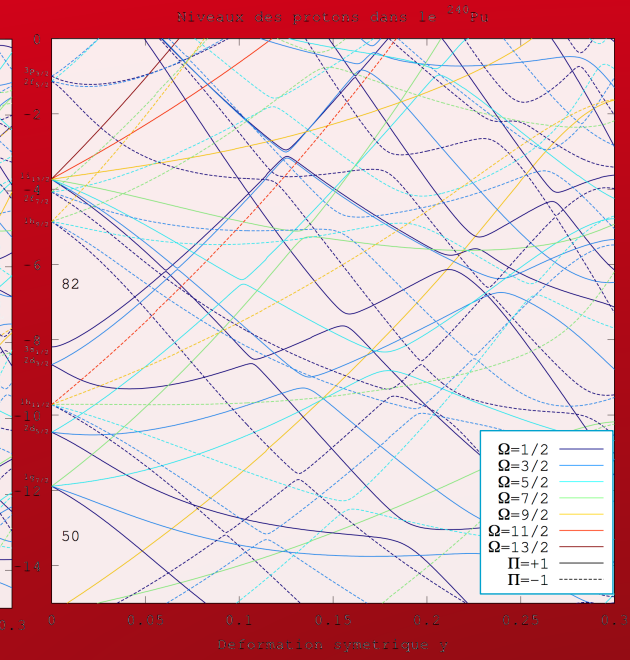
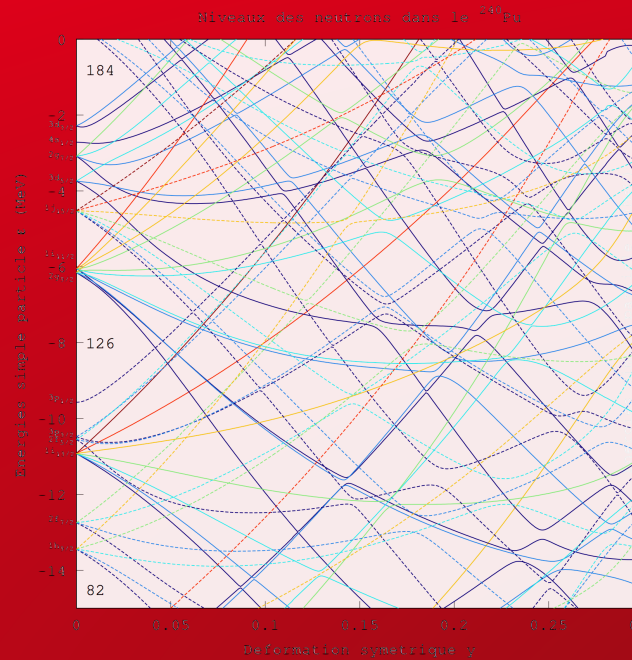
Comparison with the AVXSF** code

* O. Bouland, J.E. Lynn and P. Talou, Phys. Rev. C 88, 054612 (2013)

**J. E. Lynn, Harwell Report AERE-R 7468 (1974)



Thank you for your attention



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